

Studies of b quark decays using experiment plus lattice QCD

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Particle Physics Seminar, University of Birmingham, 7 February 2018

Outline

- Quark flavour & Lattice QCD
- DiRAC facility
- Example: $|V_{cb}|$ from $B \rightarrow D^* l \nu$

Quark Flavour
&
Lattice QCD

Motivation

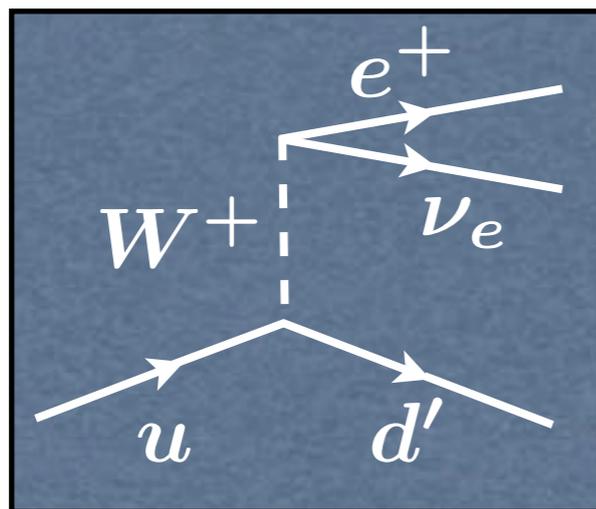
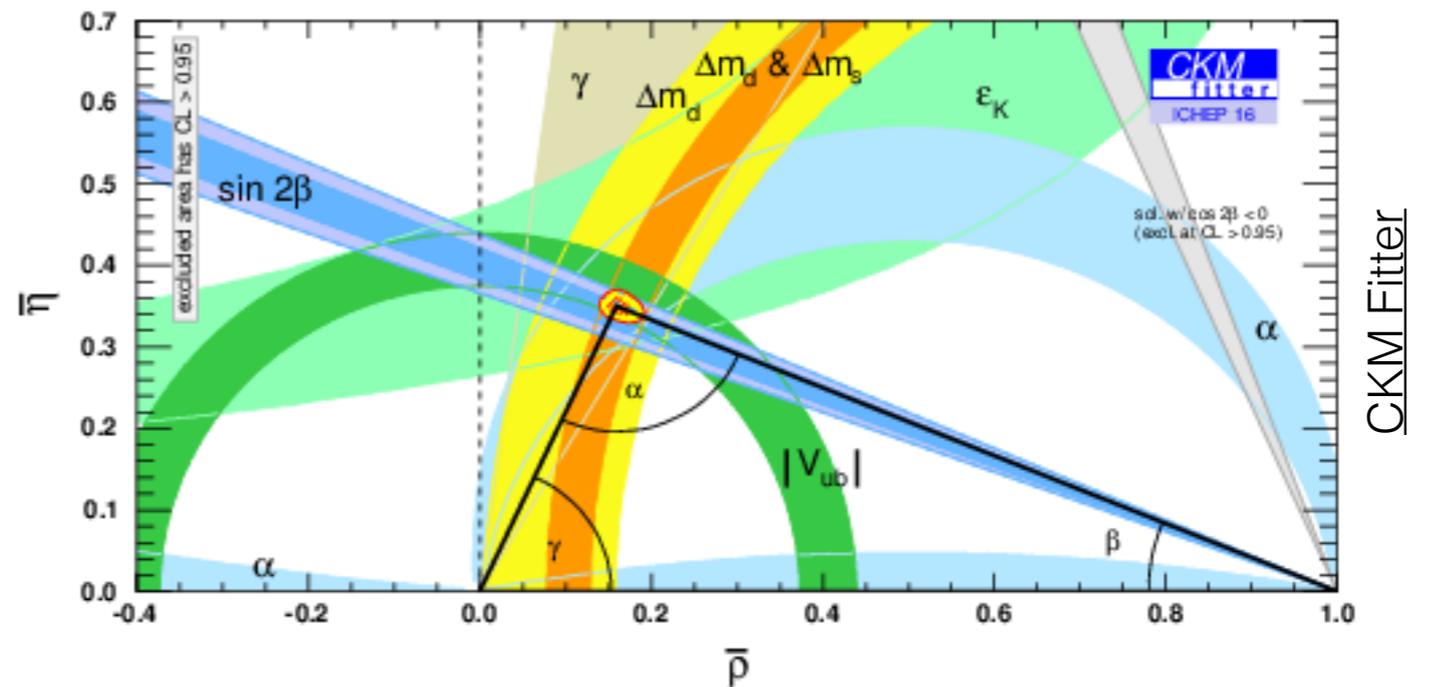
- Precision predictions & measurements of quark flavour interactions
- Is the Standard Model description of EWSB complete?
- If not, quark flavour measurements constrain models of new physics
- Experimental measurements of hadron decays: increasing precision, new modes
- Precision QCD calculations required in order to make inferences about quark interactions

Quark flavour physics

CKM matrix

$$\begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} =$$

$$\begin{pmatrix} 1 - \lambda^2/2 & \lambda & A\lambda^3(\rho - i\eta) \\ -\lambda & 1 - \lambda^2/2 & A\lambda^2 \\ A\lambda^3(1 - \rho - i\eta) & -A\lambda^2 & 1 \end{pmatrix} + \mathcal{O}(\lambda^4)$$



tree

$$K \rightarrow \pi l \nu \quad B \rightarrow \pi l \nu$$

$$D \rightarrow \pi l \nu \quad D \rightarrow K l \nu \quad B_{(s)} \rightarrow D_{(s)}^{(*)} l \nu$$

$$B_{(s)}^0 - \bar{B}_{(s)}^0 \quad B_c \rightarrow J/\psi l \nu$$

CKM matrix from Higgs couplings

LH SU(2) doublets $Q_L^i = \begin{pmatrix} u'^i \\ d'^i \end{pmatrix}_L$ RH SU(2) singlets u_R^i d_R^i

Interact with gauge bosons in covariant derivative

$$\mathcal{L}_{\text{quark}} = \bar{Q}_L^i i\not{D} Q_L^i + \bar{u}_R^i i\not{D} u_R^i + \bar{d}_R^i i\not{D} d_R^i$$

Gives rise to weak current $J_{\text{weak}}^{\mu,+} = \bar{u}'^i_L \gamma^\mu d'^i_L$

The coupling to the Higgs field is not apparently diagonal in generation

$$\mathcal{L}_{\text{quark},\phi} = -\sqrt{2} \left[\lambda_d^{ij} \bar{Q}_L^i \phi d_R^j + \lambda_u^{ij} \bar{Q}_{La}^i \epsilon_{ab} \phi_b^\dagger u_R^j + \text{h.c.} \right]$$

Fields may be transformed to mass basis

$$\mathcal{L}_{\text{quark},\phi}|_{vev} = -\sum_i (m_d^i \bar{d}_L^i d_R^i + m_u^i \bar{u}_L^i u_R^i + \text{h.c.})$$

Showing the weak current allows mixing between generations

$$J_{\text{weak}}^{\mu,+} = \bar{u}_L^i \gamma^\mu V_{CKM}^{ij} d_L^j$$

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Role of Lattice QCD

Models

Experiment

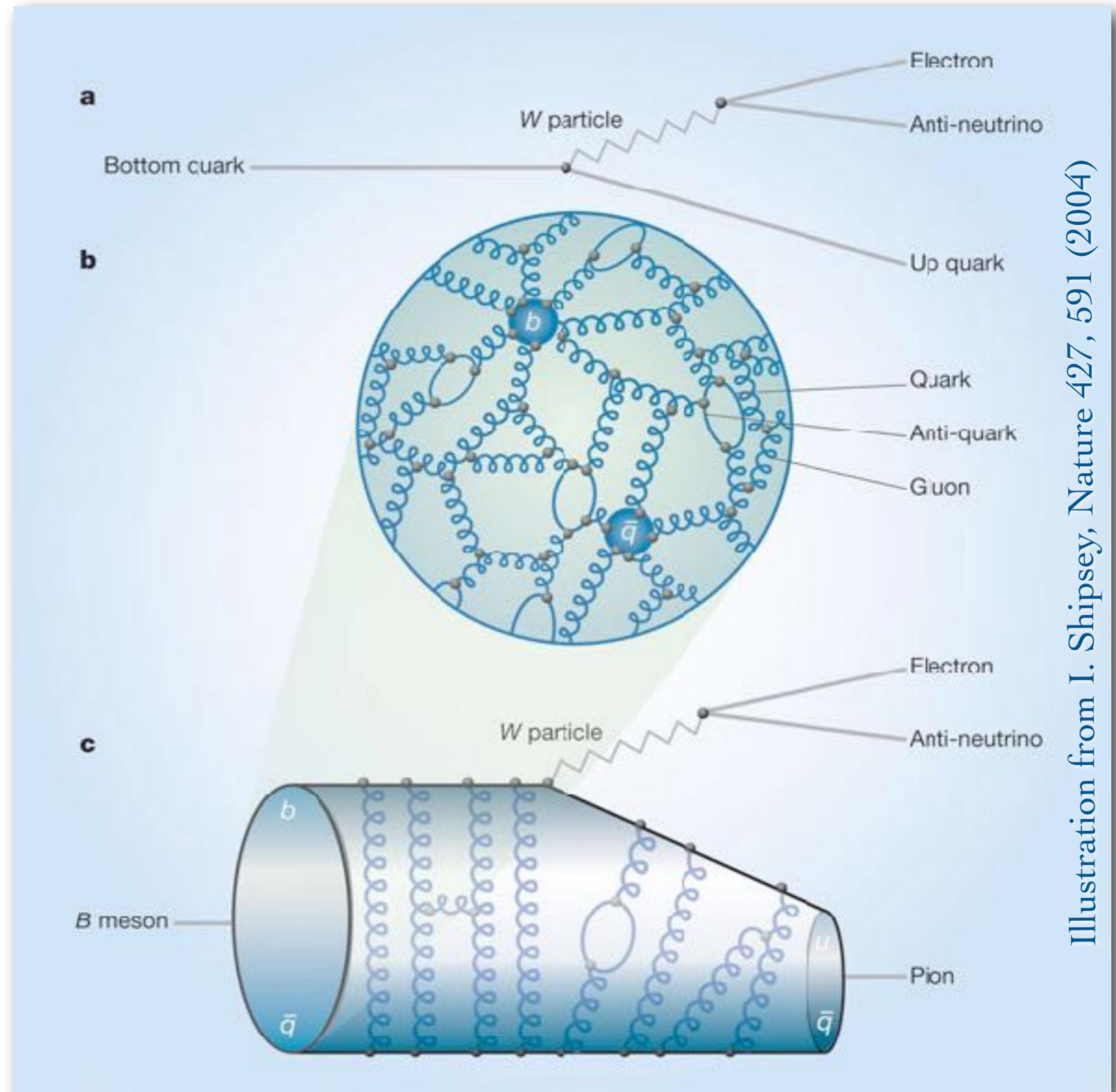


Illustration from I. Shipsey, Nature 427, 591 (2004)

Role of Lattice QCD

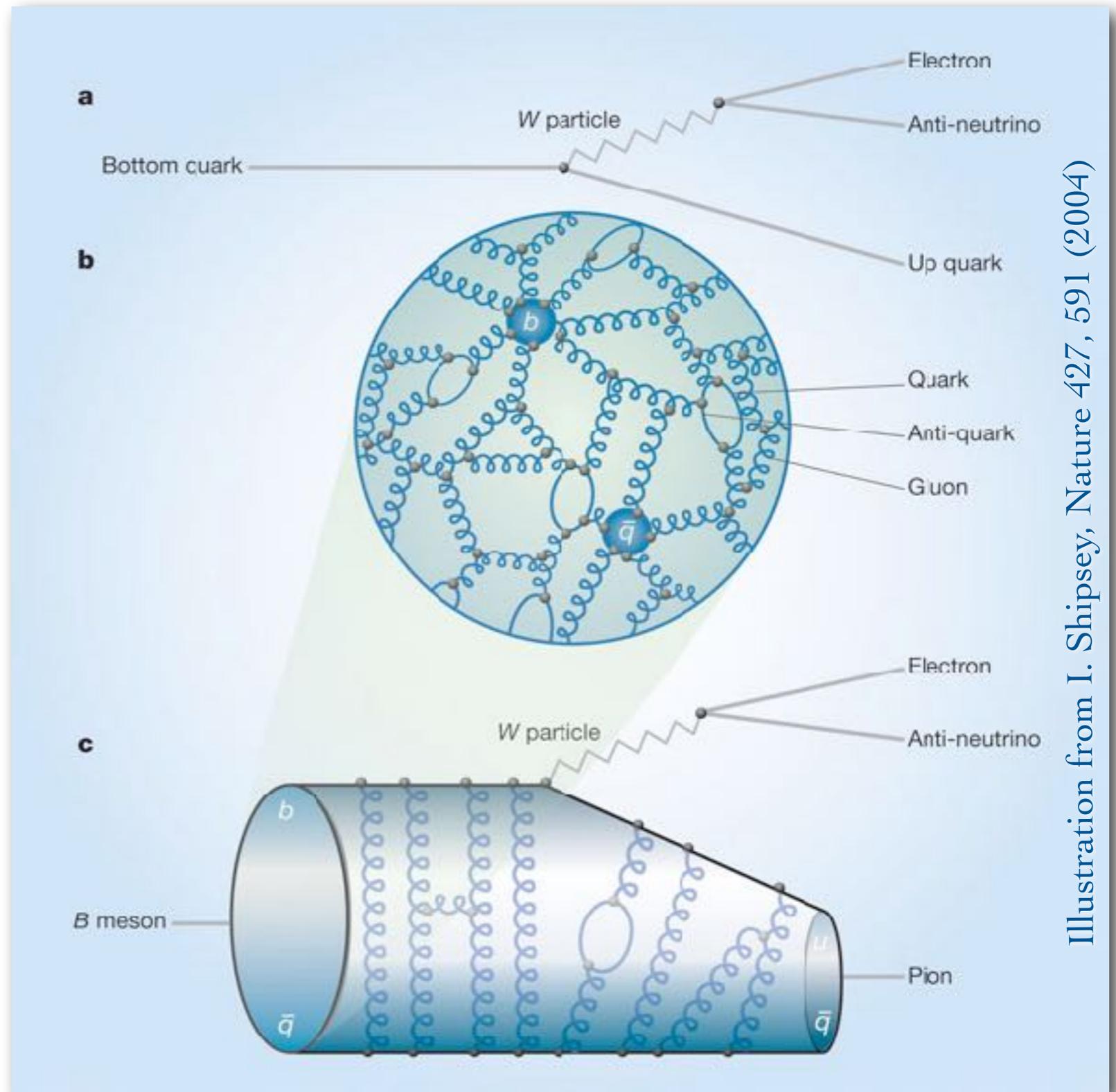
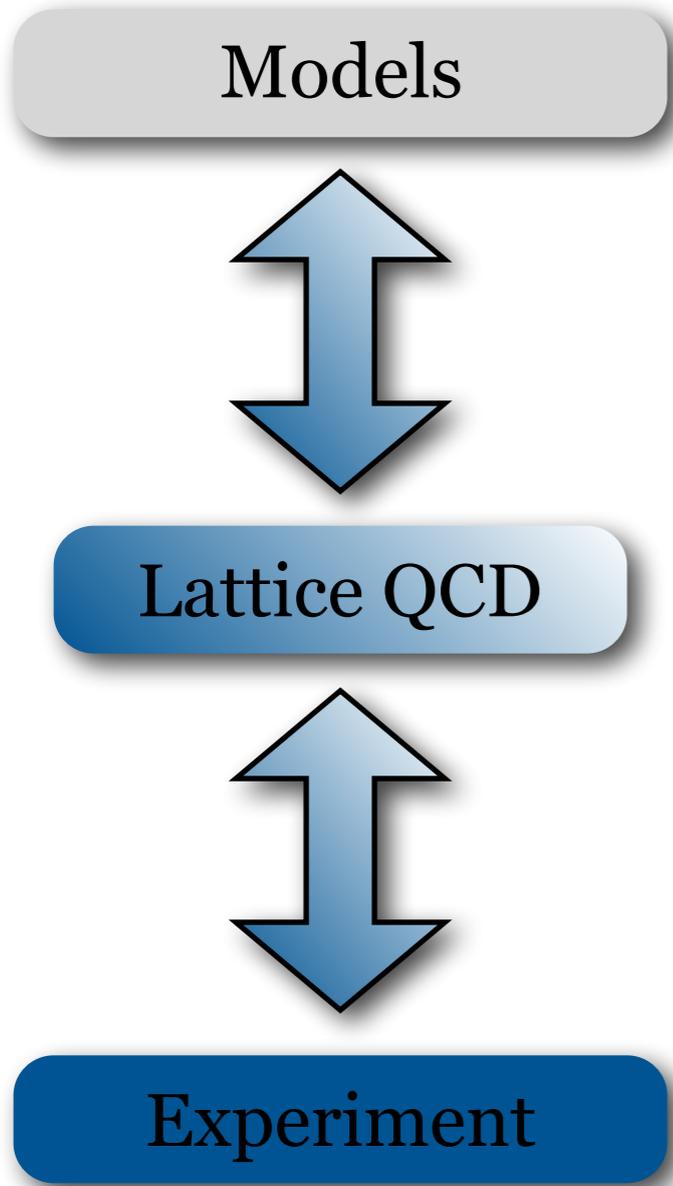


Illustration from I. Shipsey, Nature 427, 591 (2004)

Lattice QCD

- Use methods of effective field theory and renormalization to turn a **quantum physics** problem into a **statistical physics** problem
- **Quarks** propagating through strongly interacting QCD **glue** + **sea** of quark-antiquark bubbles
- Numerically evaluate path integrals using Monte Carlo methods: **importance sampling** & **correlation functions**
- Numerical challenge: solving $Mx = b$ where M is big and has a **diverging condition number** as $am_q \rightarrow 0$ (vanishing lattice spacing \times light quark mass)

Lattice QCD in a nutshell

QFT : Imaginary-time path integral

$$\langle J(z')J(z) \rangle = \frac{1}{Z} \int [d\psi][d\bar{\psi}][dU] J(z')J(z) e^{-S_E}$$

SFT : Sum over all microstates

$$\langle J(z')J(z) \rangle = \frac{1}{Z} \text{Tr} [J(z')J(z) e^{-\beta H}]$$

Use the same numerical methods!

Monte Carlo Calculation : Find and use field
“configurations” which dominate the integral/sum

Lattice QCD in a nutshell

Gluonic expectation values

$$\begin{aligned}\langle \Theta \rangle &= \frac{1}{Z} \int [d\psi][d\bar{\psi}][dU] \Theta[U] e^{-S_g[U] - \bar{\psi}Q[U]\psi} \\ &= \frac{1}{Z} \int [dU] \Theta[U] \det Q[U] e^{-S_g[U]}\end{aligned}$$

Fermionic expectation values

$$\langle \bar{\psi}\Gamma\psi \rangle = \int [dU] \frac{\delta}{\delta \bar{\zeta}} \Gamma \frac{\delta}{\delta \zeta} e^{-\bar{\zeta}Q^{-1}[U]\zeta} \det Q[U] e^{-S_g[U]} \Big|_{\zeta, \bar{\zeta} \rightarrow 0}$$

Lattice QCD in a nutshell

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Probability weight

Lattice QCD in a nutshell

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Probability weight

Determinant in probability weight difficult

- 1) Requires nonlocal updating;
- 2) Matrix becomes singular

Lattice QCD in a nutshell

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Partial quenching =

different mass for valence Q^{-1} than for sea $\det Q$

Lattice QCD

$$\langle \Phi_\pi(z) V_\mu(y) \Phi_B(x) \rangle = \frac{1}{Z} \int [d\psi][d\bar{\psi}][dU] \Phi_\pi(z) V_\mu(y) \Phi_B(x) e^{-S[\psi, \bar{\psi}, U]}$$

- Imaginary time formulation: path integrands real, non-negative
- Discrete lattice points: regulates field theory
- Sharply peaked path integrand: permits importance sampling

Systematic error

Controllable limit

Theory

Lattice volume

$$L \gg 1/m_\pi$$

Chiral pert. th.
Brute force

Lattice spacing

$$a \ll 1/\Lambda_{\text{QCD}}$$

Symanzik EFT

Light quark mass

$$m_\pi \ll m_\rho, 4\pi f_\pi$$

Chiral pert. th.
Brute force

Heavy quark mass

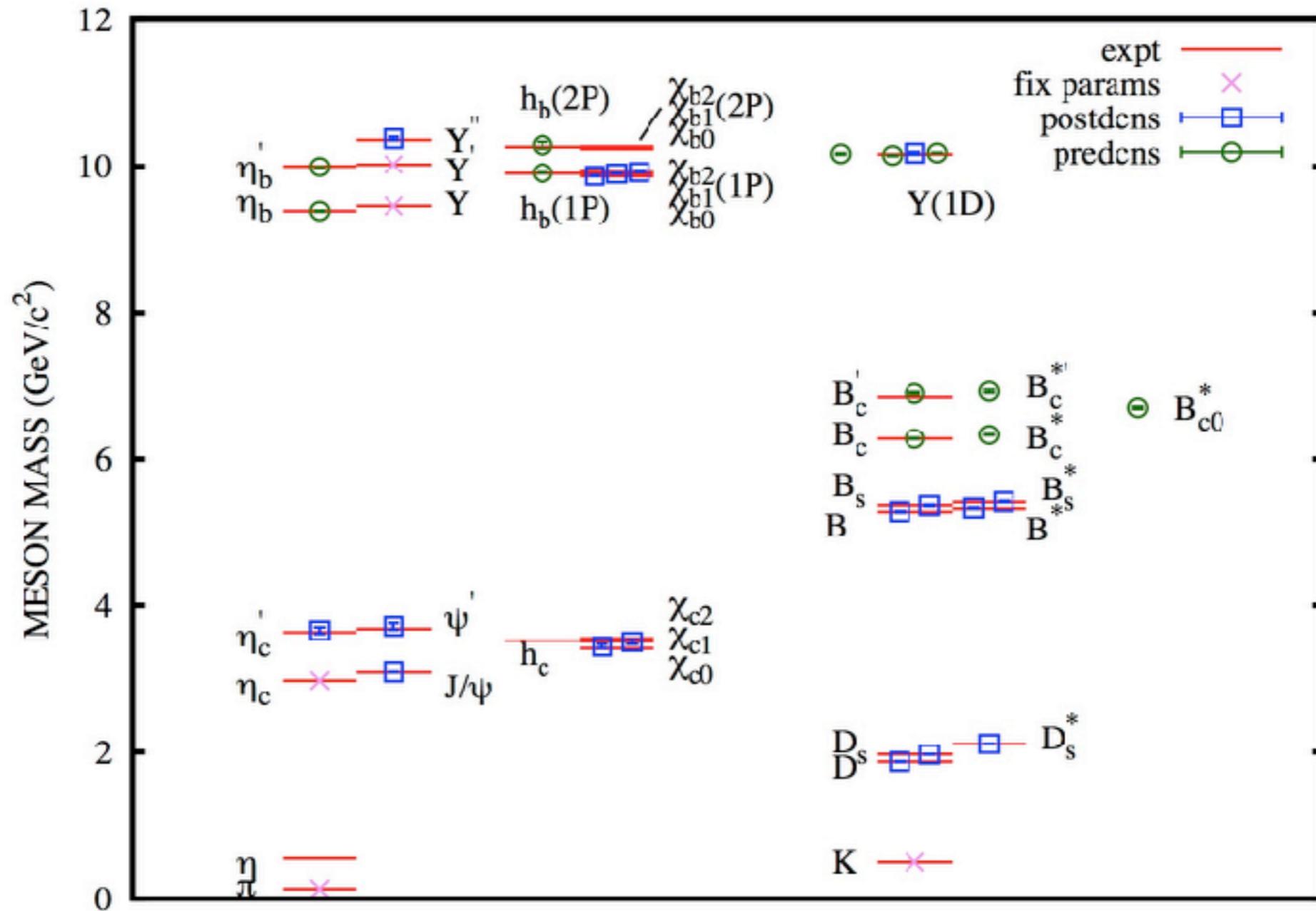
$$m_Q \gg 1/a$$

$$m_Q < 1/a$$

$$m_Q \approx 1/a$$

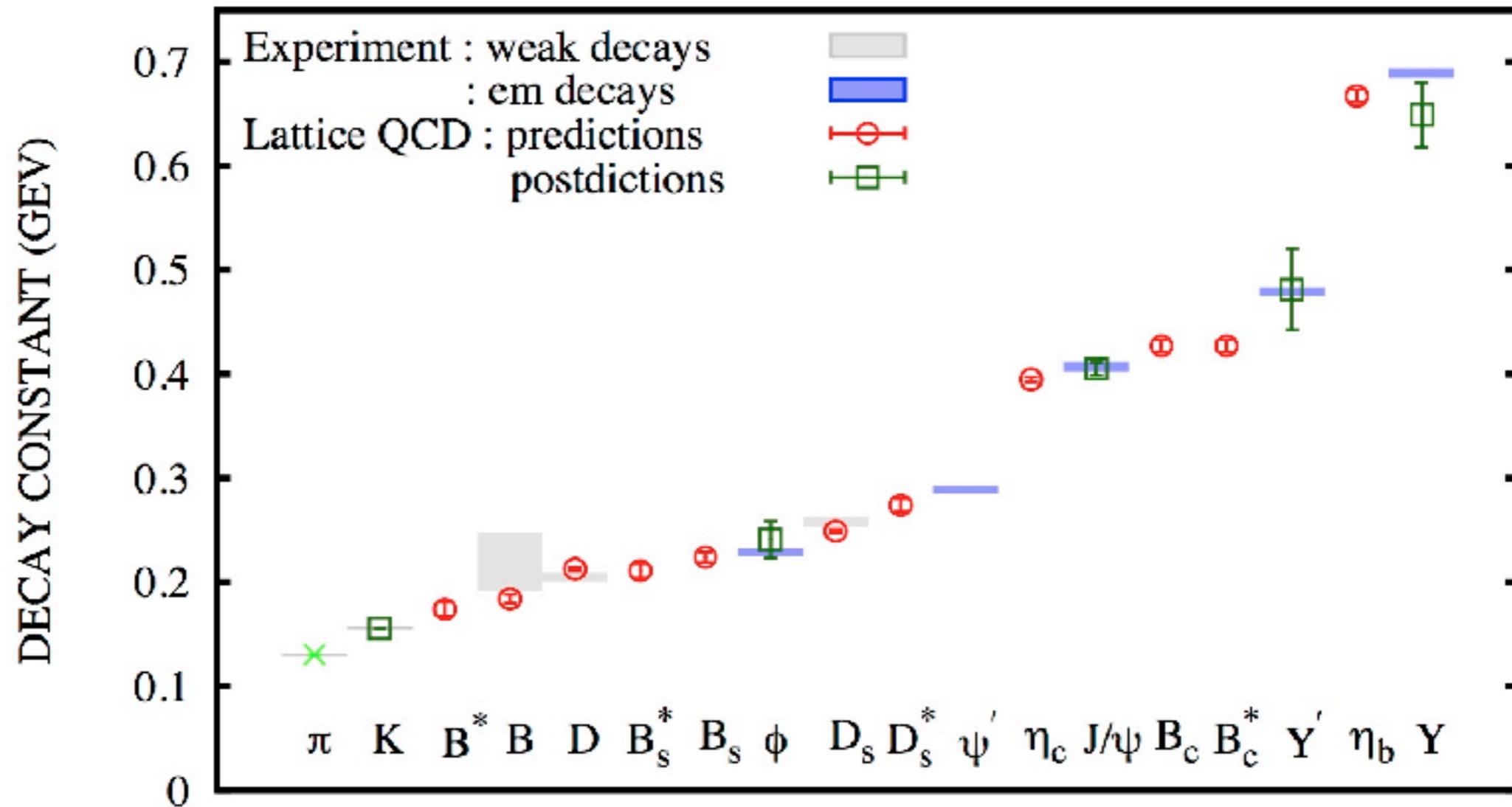
NRQCD, HQET
Extra-fine, extra-improvement
Fermilab

Meson mass splittings



CTH Davies, [HPQCD Collaboration [website](#)]

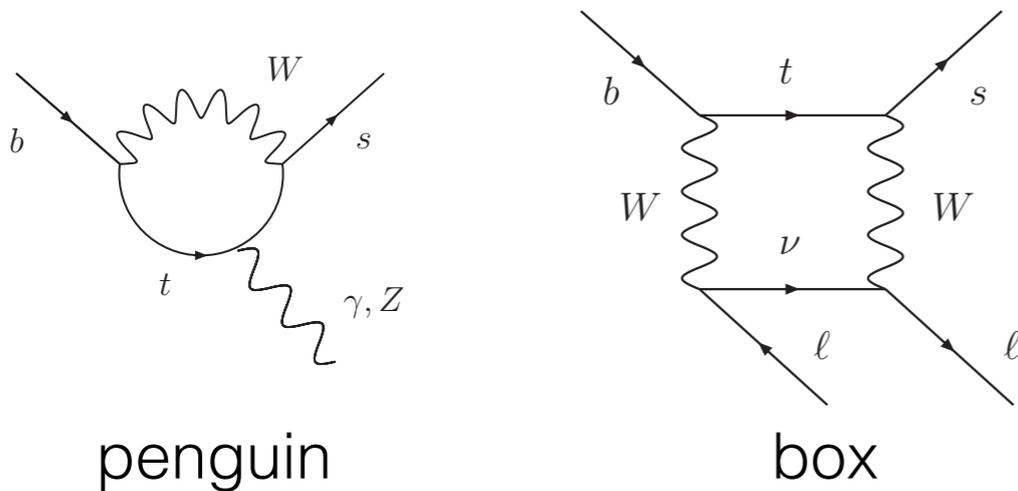
Decay constants



CTH Davies, [HPQCD Collaboration [website](#)]

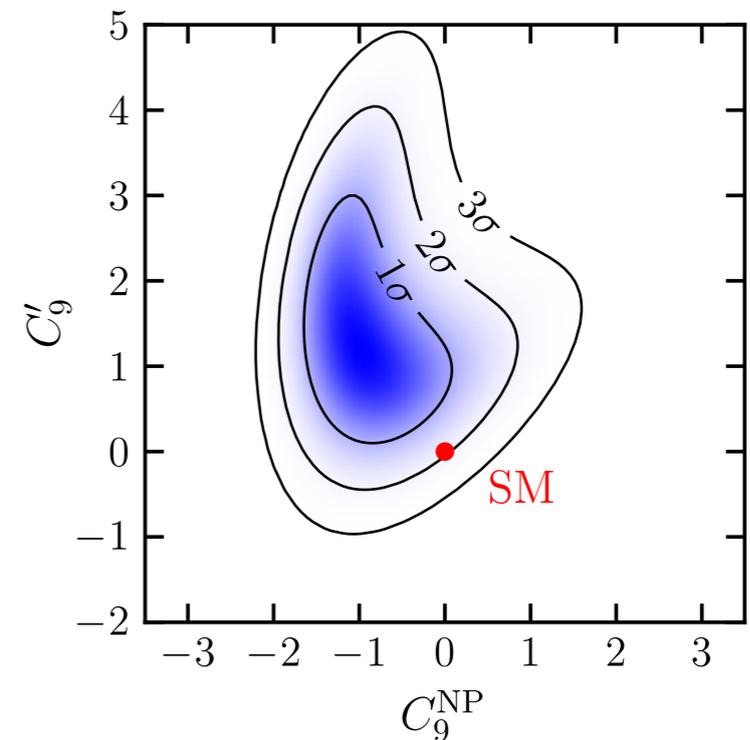
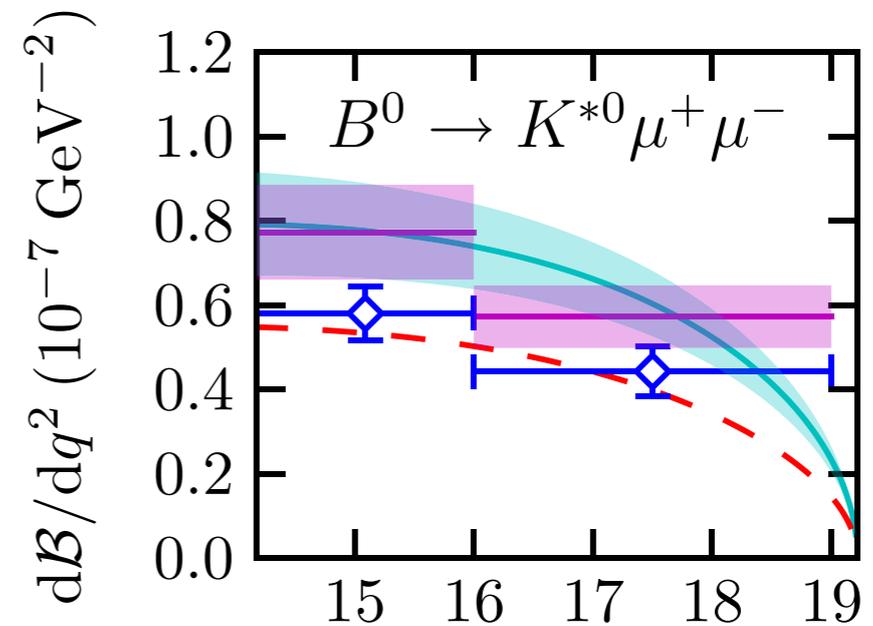
Rare b decays

Flavour changing neutral decays



$$B \rightarrow K^* l^+ l^- \quad B_s \rightarrow \phi l^+ l^-$$

Horgan et al., (HPQCD) [arXiv:1310.3722](https://arxiv.org/abs/1310.3722), [arXiv:1310.3887](https://arxiv.org/abs/1310.3887)



LQCD & DiRAC

UKQCD consortium

- 24 faculty at 8 UK institutions
- Membership/Leadership in several international collaborations (e.g. HPQCD, RBC-UKQCD, HadSpec, QCDSF, FastSum)
- Broad range of physics: quark flavour, hadron spectrum, hot/dense QCD; BSM theories of EWSB, dark matter
- Widespread impact: LHC, BES-III, Belle, JLab, J-PARC, FAIR, RHIC, NA62



Image credit: CIA World Factbook

DiRAC 2

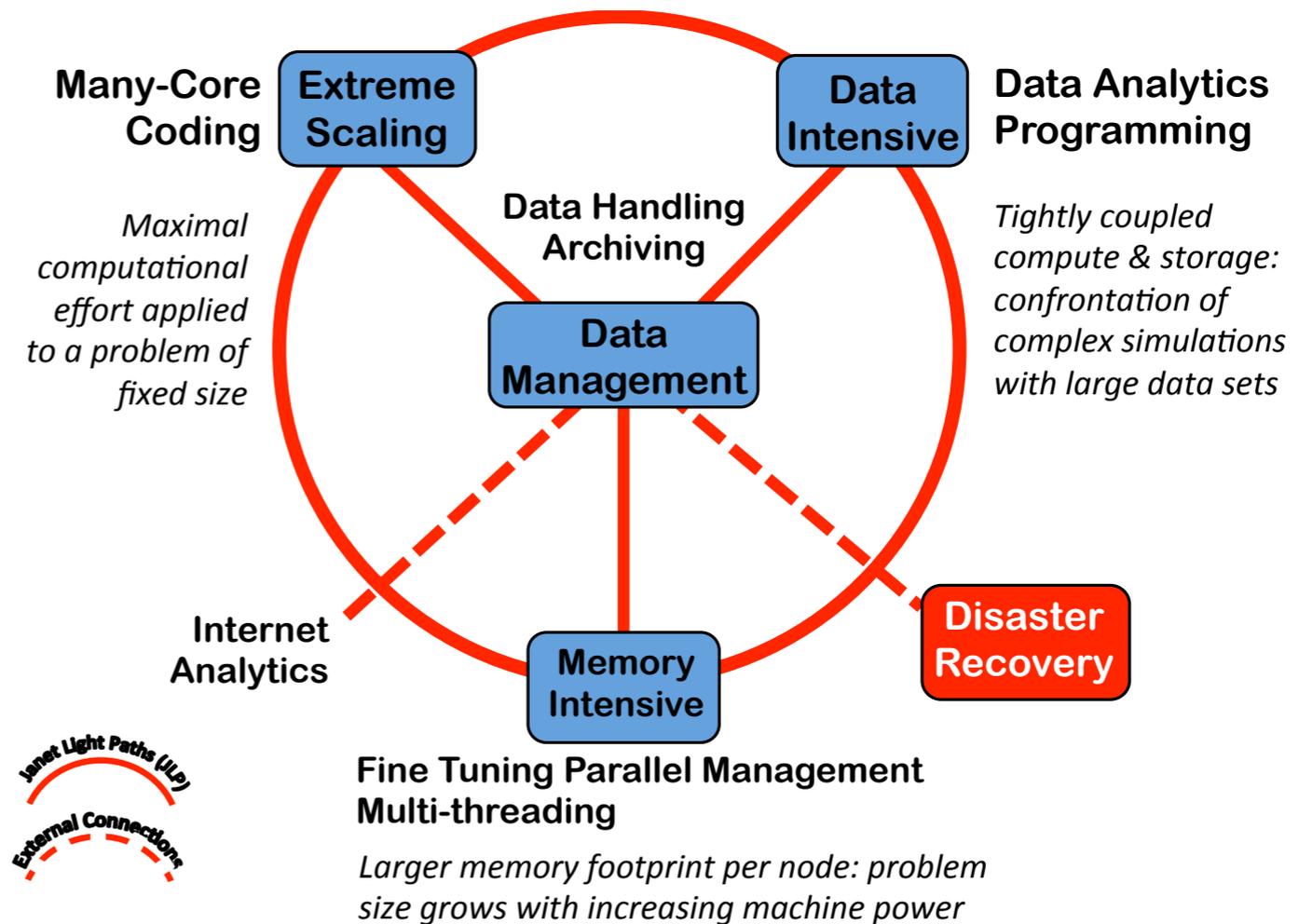
- 2011: £15M BIS investment in national distributed HPC facility for particle & nuclear physics, cosmology, & theoretical astrophysics. Recurrent costs funded by STFC
- 2012: 5 systems deployed:
 - **Extreme scaling:** 1.3 Pflop/s Blue Gene/Q (Edinburgh)
 - **Data Analytic/Data Centric/Complexity:** 3 tightly-coupled clusters with various levels of interconnectivity, memory, and fast I/O (Cambridge, Durham, Leicester)
 - **Shared Memory System (SMP)** (Cambridge)
- Service started 1 December 2012

DiRAC 2 outputs

- 106 lattice publications, with 1977 citations (as of 20/7/2017)
- 765 publications in a broad scientific range (PPAN) — 35,365 citations (as of 20/7/2017)
- Gravitational waves, cosmology, galaxy & planet formation, exoplanets, MHD, particle pheno, nuclear physics
- Valuable resource for PDRA's & PhD students
- Scientific results, training in high performance computing

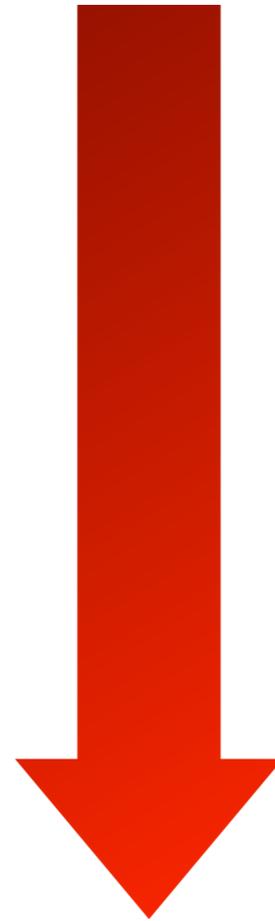
DiRAC 3

- Continued success requires continued investment
- Seek approx £25M capital investment to upgrade DiRAC-2 x10
- Running costs for staff and electricity
- Improve exploitation of research and HPC training impact with PDRA and PhD support (Big Data CDTs)
- Part of RCUK's e-Infrastructure roadmap



2011/12

DiRAC 2



Stop-gap funding:

2016/17 DiRAC 2.5

2017 DiRAC 2.5x

2018/19

DiRAC 3

DiRAC 2.5

After £1.67M capital injection

- **Extreme Scaling 2.5:** 1.3 Pflop/s Blue Gene/Q
- **Data Analytic 2.5:** Share of Peta5 system + continued access to Sandybridge system
 - Shared EPSRC/DiRAC/Cambridge: 25K Skylake cores + 1.0 Pflop/s GPU + 0.5 Pflop/s KNL service
- **Data Centric 2.5:** Over 14K cores, 128 GB RAM/node
- **Complexity 2.5:** 4.7K large-job cores + 3K small-job cores
- **SMP:** 14.8TB, 1.8K core shared memory service

DiRAC 2.5x

June 2017: £9M capital funding (BEIS), lifeline to DiRAC3:

- Planned investment
 - **Extreme scaling:** 1024-node, 2.5 Pflop/s system
 - **Memory intensive:** 144 nodes, 4.6K cores, 110 TB RAM
 - **Data analytic:** 128 nodes, 4K cores, 256GB/node; hierarchy of fat nodes (1-6 TB); NVMe storage for data intensive workflows
- Additional storage at all DiRAC sites
- Procurement procedure: November 2017
- Target for hardware availability: April 2018

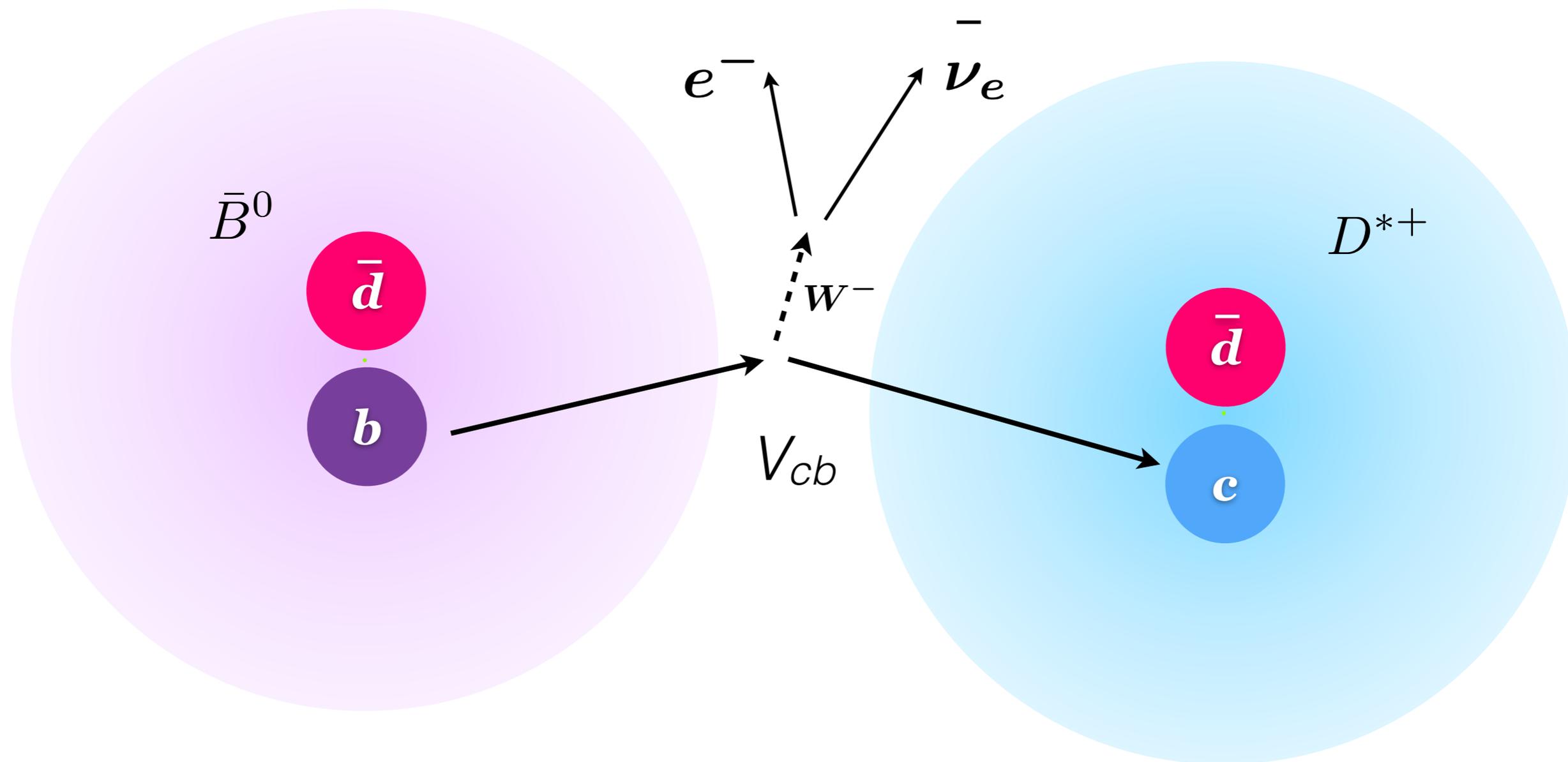
DiRAC & LQCD

- Capital expenditure has come directly from BIS/BEIS, running costs through STFC
- DiRAC has allowed the UK to be a major contributor to world-wide Lattice QCD (and BSM) efforts
- High precision theory needed to make the most of high precision experiment

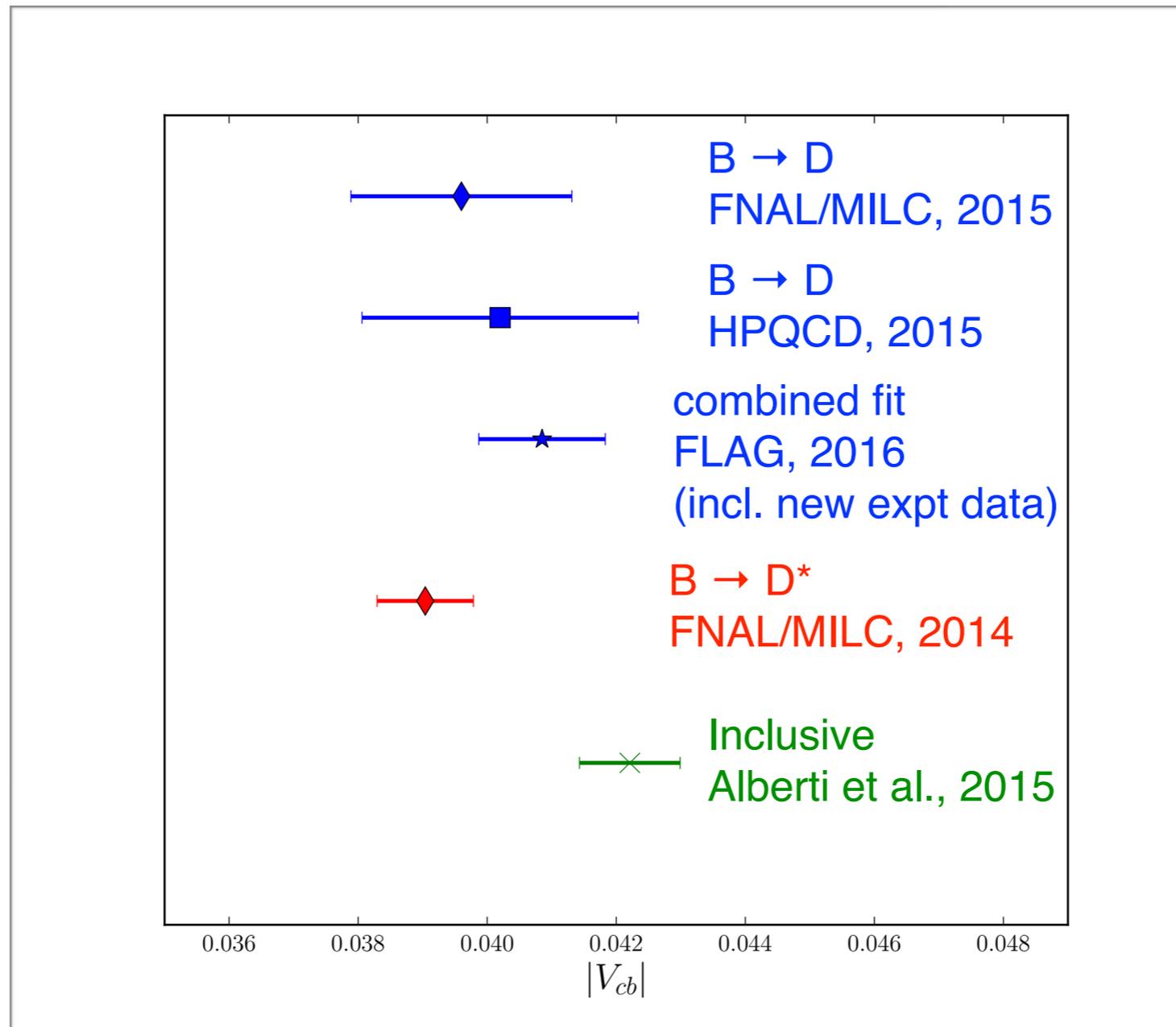
$B \rightarrow D^* | v$

and V_{cb}

$$B \rightarrow D^* | \nu$$

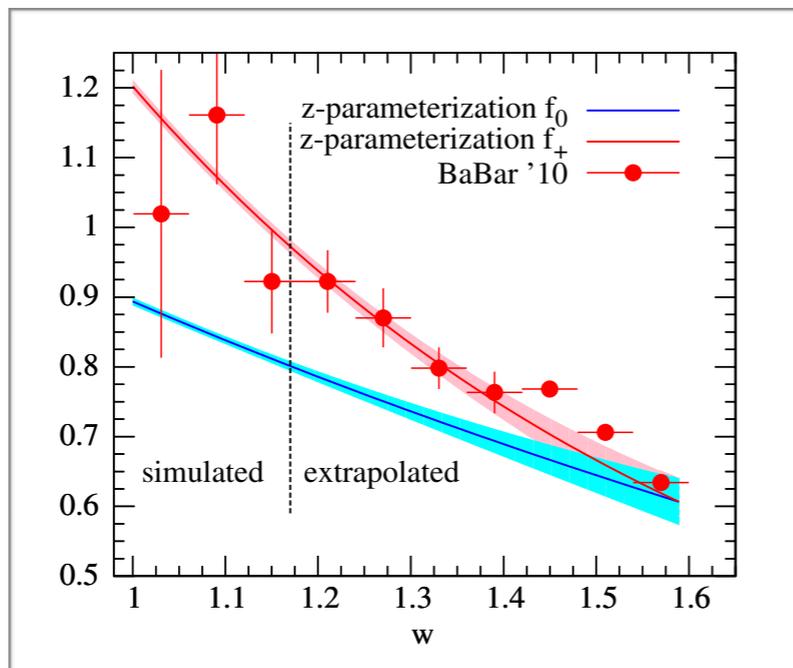


$|V_{cb}|$ (before 2/2017)



Published $B \rightarrow D$

Fermilab/
MILC



HPQCD

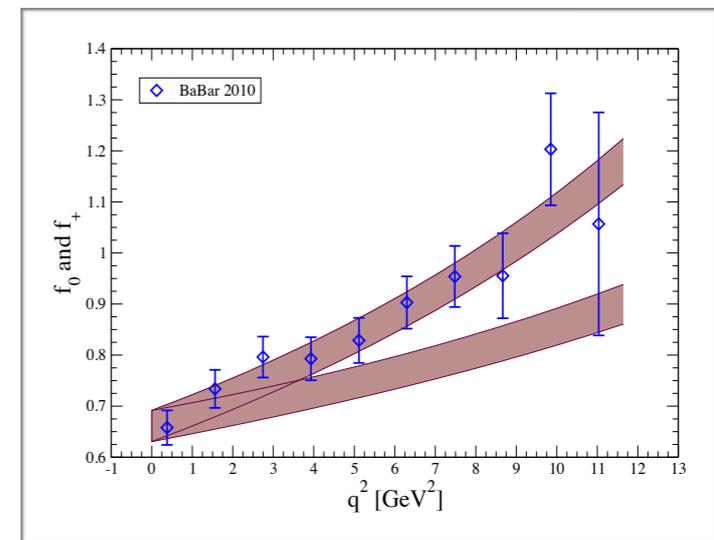


TABLE V. Error budget table for $|V_{cb}|$. The first three rows are from experiments, and the rest are from lattice simulations.

Source	$f_+(\%)$	$f_0(\%)$
Statistics+matching+ χ PT cont. extrapol.	1.2	1.1
(Statistics)	(0.7)	(0.7)
(Matching)	(0.7)	(0.7)
(χ PT/cont. extrapol.)	(0.6)	(0.5)
Heavy-quark discretization	0.4	0.4
Lattice scale r_1	0.2	0.2
Total error	1.2	1.1

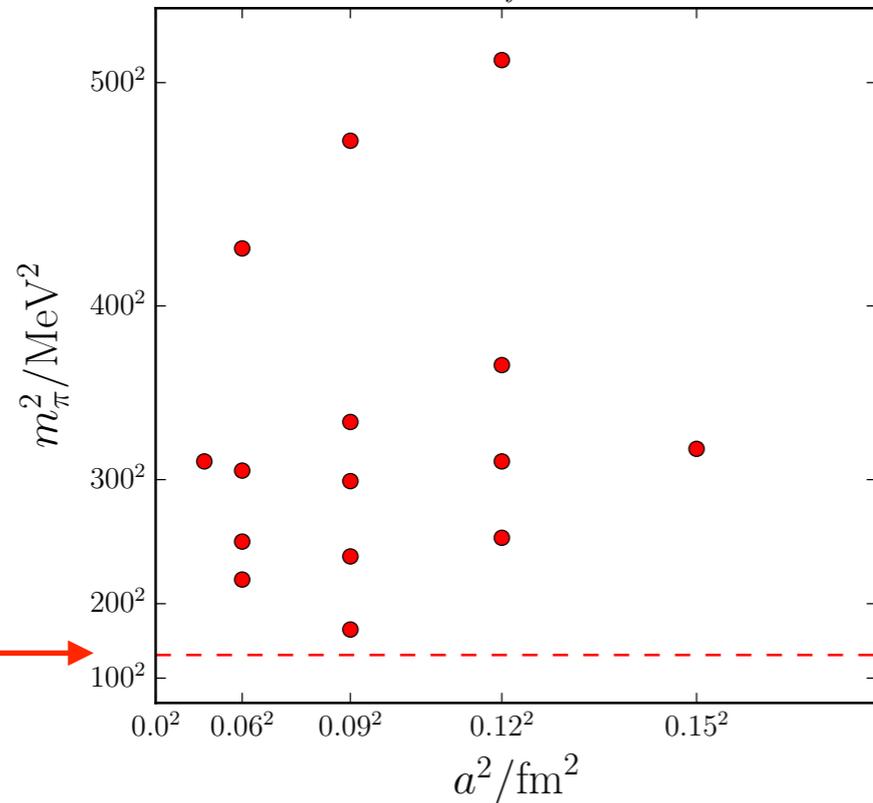
Type	Partial errors [%]
experimental statistics	1.55
experimental systematic	3.3
meson masses	0.01
lattice statistics	1.22
chiral extrapolation	1.14
discretization	2.59
kinematic	0.96
matching	2.11
electro-weak	0.48
finite size effect	0.1
total	5.34

Bailey et al. (FNAL/MILC), arXiv:1503.07237

Na et al. (HPQCD), arXiv:1505.03925

Published $B \rightarrow D^*$

MILC $n_f = 2 + 1$

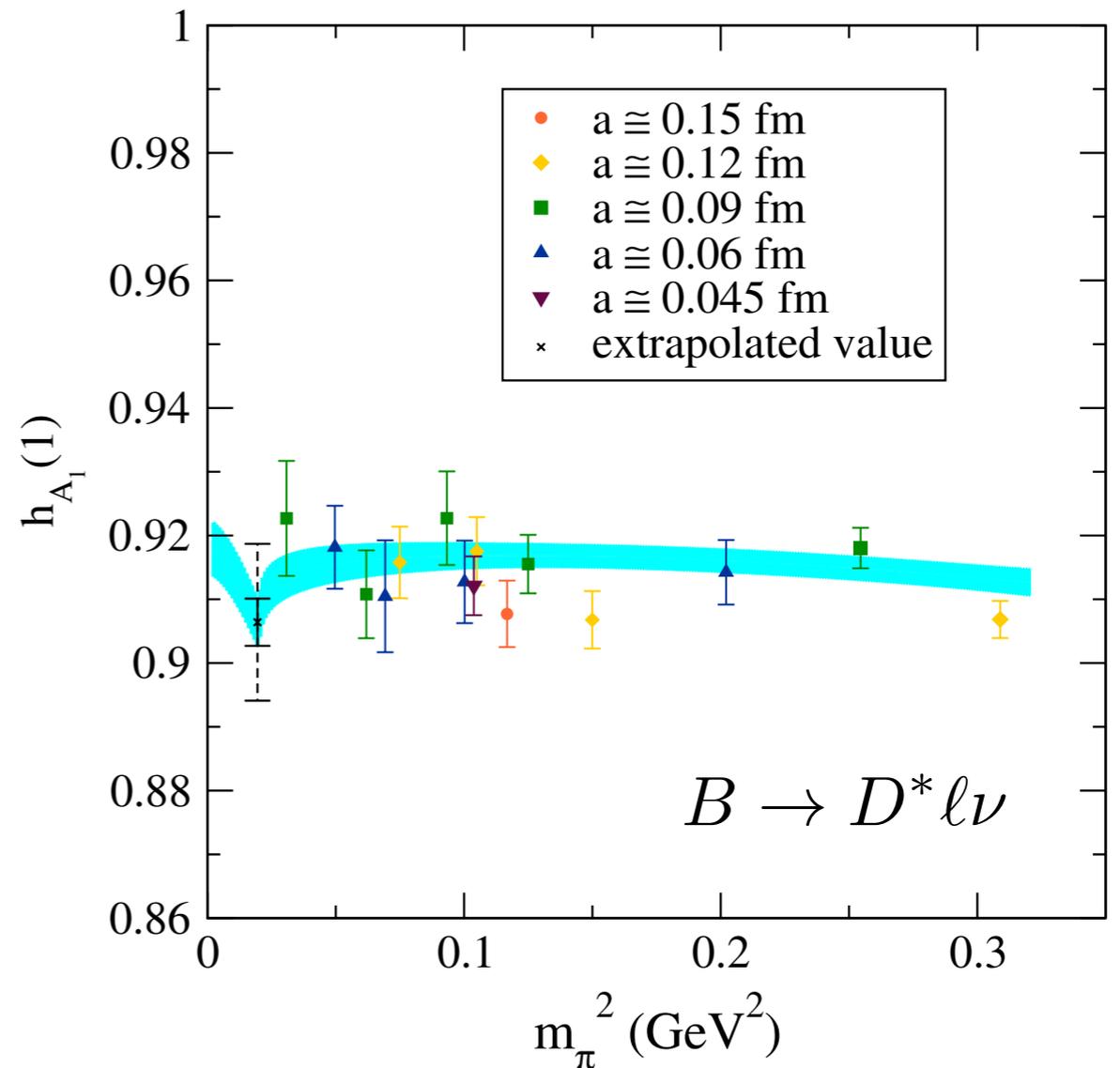


nature

TABLE X. Final error budget for $h_{A_1}(1)$ where each error is discussed in the text. Systematic errors are added in quadrature and combined in quadrature with the statistical error to obtain the total error.

Uncertainty	$h_{A_1}(1)$
Statistics	0.4%
Scale (r_1) error	0.1%
χ PT fits	0.5%
$g_{D^* D \pi}$	0.3%
Discretization errors	1.0%
Perturbation theory	0.4%
Isospin	0.1%
Total	1.4%

FNAL/MILC

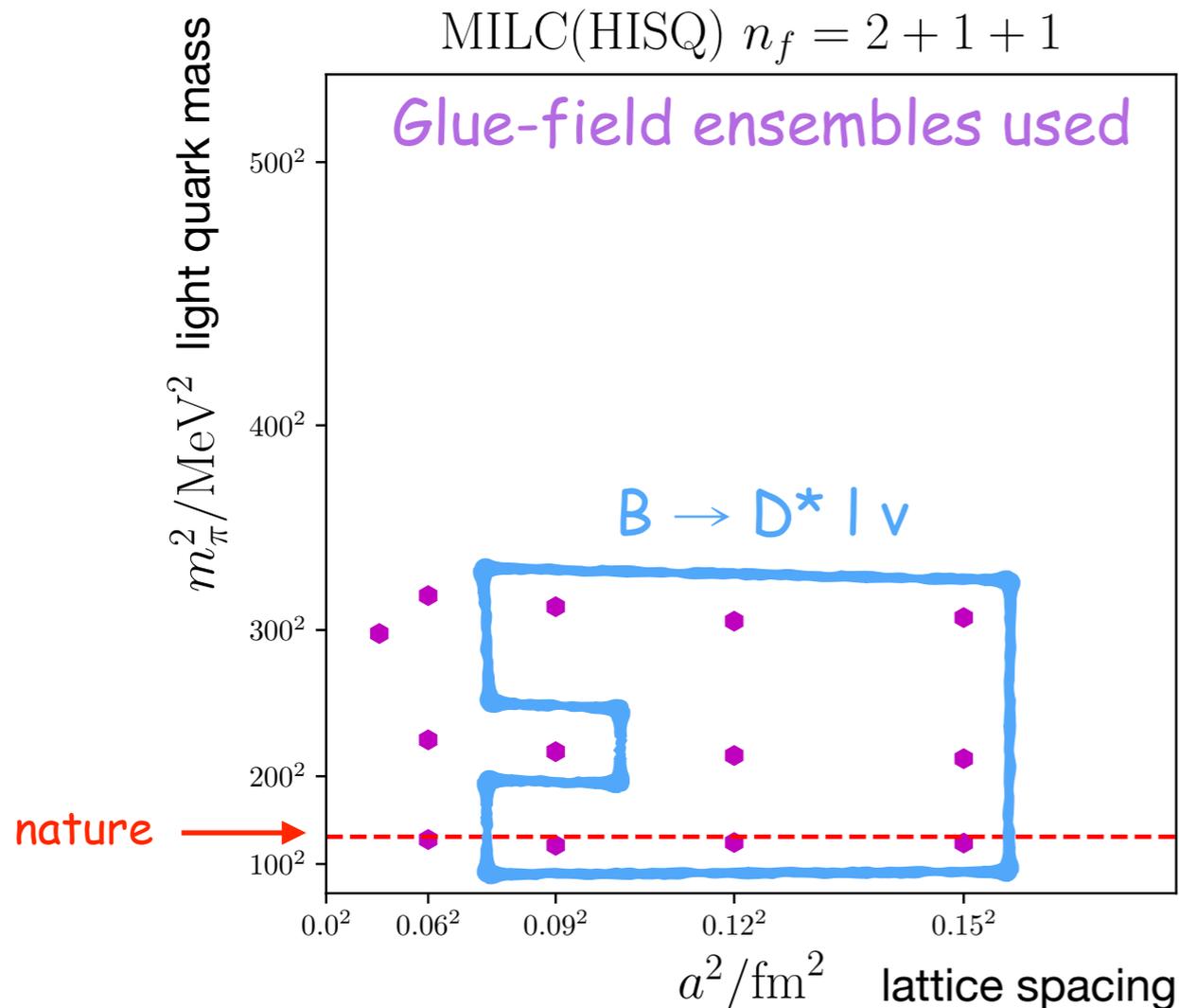


$B \rightarrow D^* \ell \nu$

Bailey et al. (FNAL/MILC), PRD89 (2014)

HPQCD calculation

Judd Harrison, Christine Davies, MBW (HPQCD), [arXiv:1711.11013](https://arxiv.org/abs/1711.11013)



- Statistically independent calculations from Fermilab/MILC
- HISQ vs. AsqTad light/strange
- HISQ vs. FNAL charm
- NRQCD vs. FNAL bottom

Zero recoil

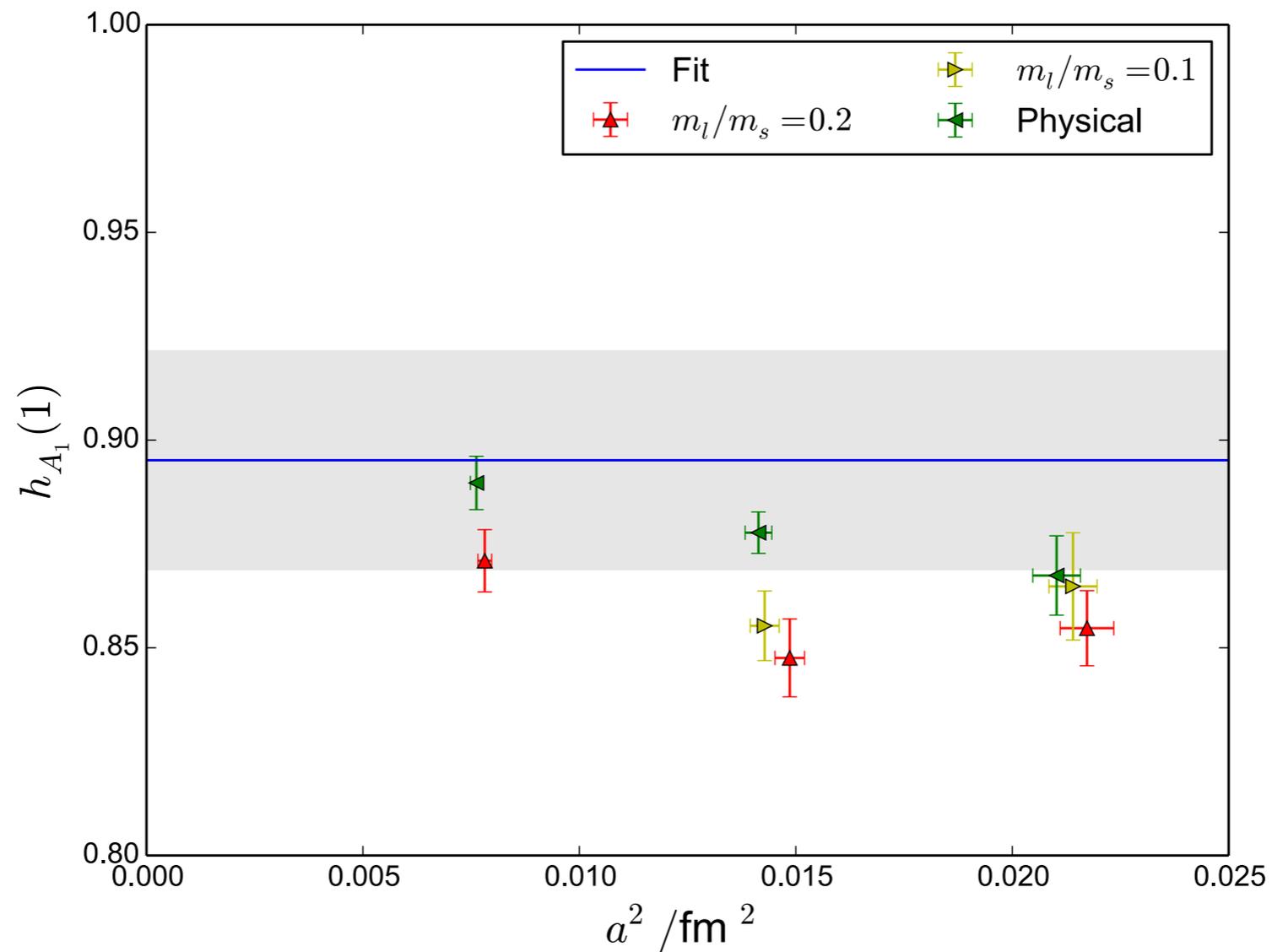
$$\frac{d\Gamma}{dw}(\bar{B}^0 \rightarrow D^{*+} l^- \bar{\nu}_l) = \frac{G_F^2 M_{D^*}^3 |\bar{\eta}_{EW} V_{cb}|^2}{4\pi^3} (M_B - M_{D^*})^2 \sqrt{w^2 - 1} \chi(w) |\mathcal{F}(w)|^2$$

$$\chi(1) = 1 \quad \mathcal{F}(1) = h_{A_1}(1) = \frac{M_B + M_{D^*}}{2\sqrt{M_B M_{D^*}}} A_1(q_{\max}^2)$$

$$\langle D^*(p', \epsilon) | \bar{q} \gamma^\mu \gamma^5 Q | B(p) \rangle = 2M_{D^*} A_0(q^2) \frac{\epsilon^* \cdot q}{q^2} q^\mu + (M_B + M_{D^*}) A_1(q^2) \left[\epsilon^{*\mu} - \frac{\epsilon^* \cdot q}{q^2} q^\mu \right] - A_2(q^2) \frac{\epsilon^* \cdot q}{M_B + M_{D^*}} \left[p^\mu + p'^\mu - \frac{M_B^2 - M_{D^*}^2}{q^2} q^\mu \right].$$

$$\langle D^*(p', \epsilon) | \bar{q} \gamma^\mu Q | B(p) \rangle = \frac{2iV(q^2)}{M_B + M_{D^*}} \epsilon^{\mu\nu\rho\sigma} \epsilon_\nu^* p'_\rho p_\sigma$$

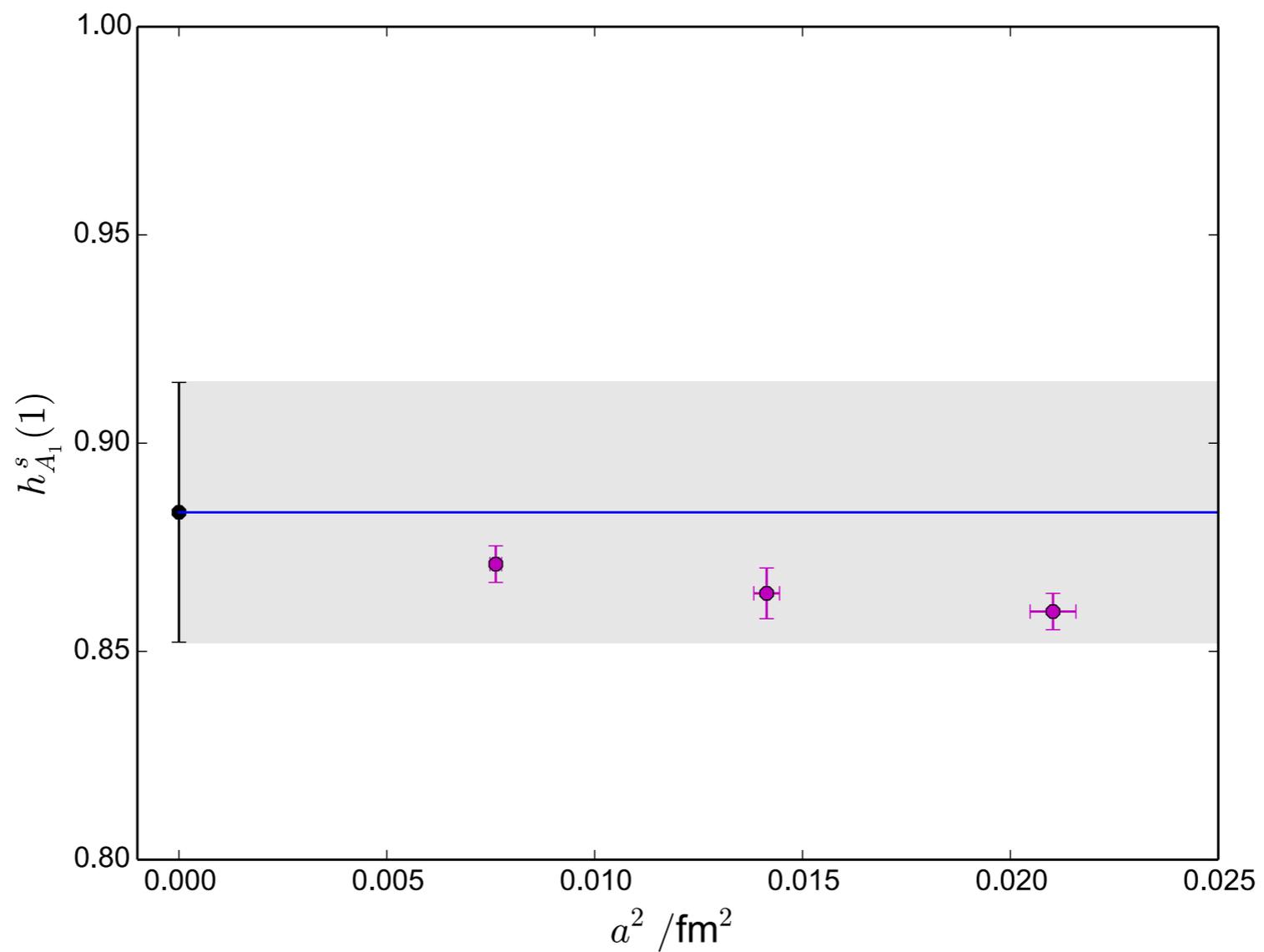
$B \rightarrow D^*$ — lattice spacing



Continuum-
physical mass
curve

$$\mathcal{F}^{B \rightarrow D^*}(1) = h_{A_1}(1) = 0.895(10)_{\text{stat}}(24)_{\text{sys}}$$

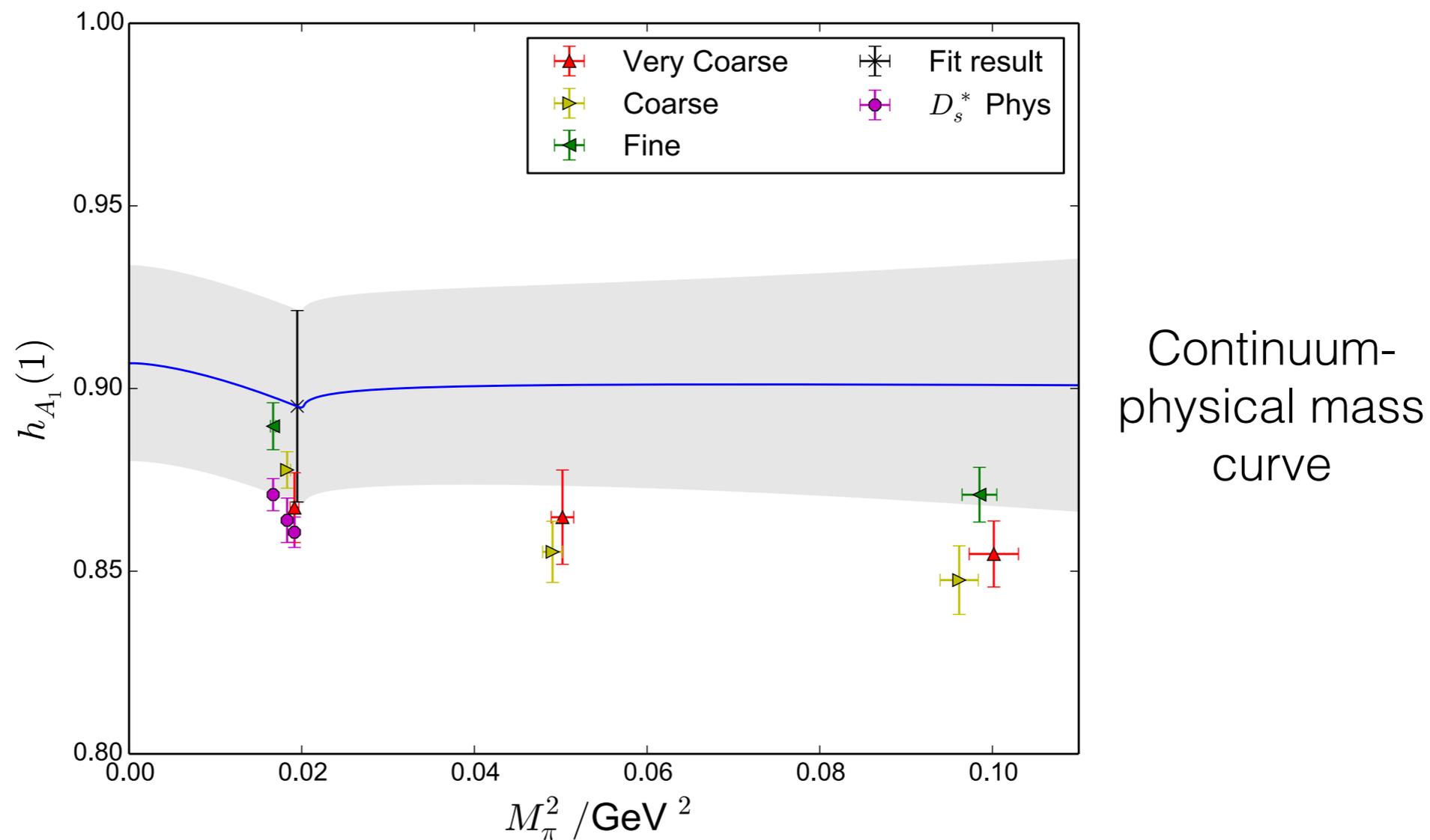
$B_s \rightarrow D_s^*$ — lattice spacing



Continuum-
physical mass
curve

$$\mathcal{F}^{B_s \rightarrow D_s^*}(1) = h_{A_1}^s(1) = 0.883(12)_{\text{stat}}(28)_{\text{sys}}$$

$B_{(s)} \rightarrow D_{(s)}^*$ — quark mass



Lattice results

$$\mathcal{F}^{B \rightarrow D^*}(1) = h_{A_1}(1) = 0.895(10)_{\text{stat}}(24)_{\text{sys}}$$

$$\mathcal{F}^{B_s \rightarrow D_s^*}(1) = h_{A_1}^s(1) = 0.883(12)_{\text{stat}}(28)_{\text{sys}}$$

$$\frac{\mathcal{F}^{B \rightarrow D^*}(1)}{\mathcal{F}^{B_s \rightarrow D_s^*}(1)} = \frac{h_{A_1}(1)}{h_{A_1}^s(1)} = 1.013(14)_{\text{stat}}(17)_{\text{sys}}$$

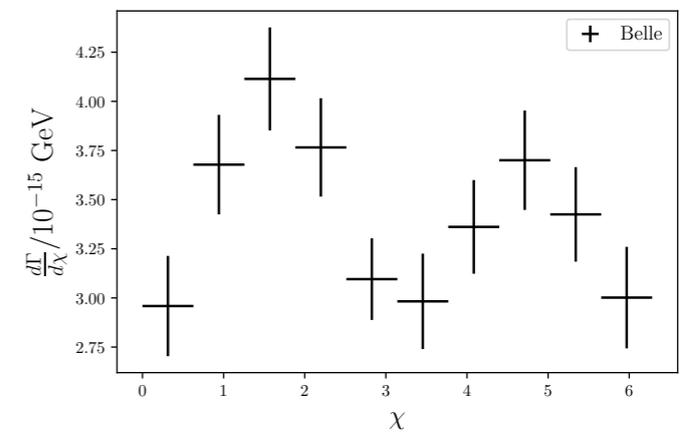
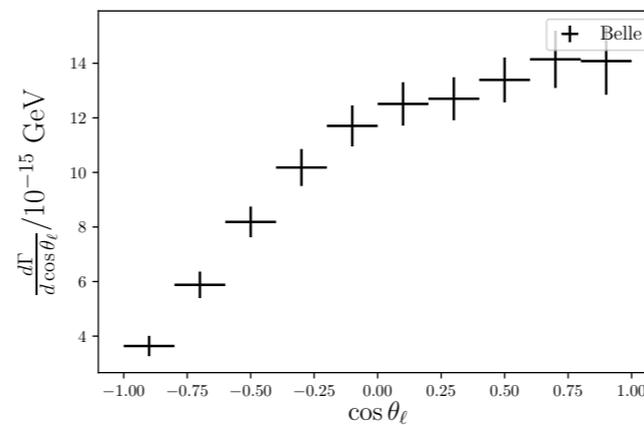
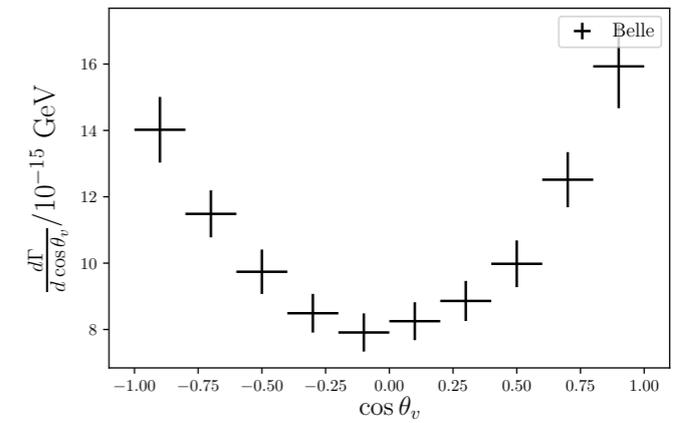
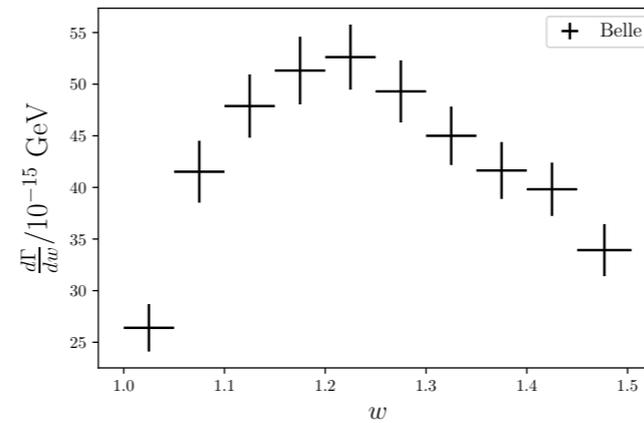
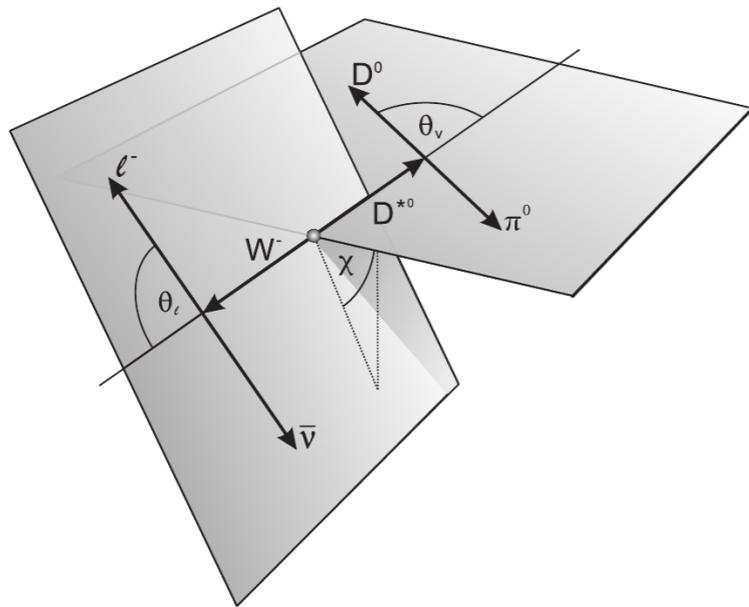
Uncertainty	$h_{A_1}(1)$	$h_{A_1}^s(1)$	$h_{A_1}(1)/h_{A_1}^s(1)$
α_s^2	2.1	2.5	0.4
$\alpha_s \Lambda_{\text{QCD}}/m_b$	0.9	0.9	0.0
$(\Lambda_{\text{QCD}}/m_b)^2$	0.8	0.8	0.0
a^2	0.7	1.4	1.4
$g_{D^* D \pi}$	0.2	0.03	0.2
Total systematic	2.7	3.2	1.7
Data	1.1	1.4	1.4
Total	2.9	3.5	2.2

- Good agreement with Fermilab/MILC result $h_{A_1}(1) = 0.906(4)(12)$
- Independent lattices
- Different heavy quark formulations

Implications for V_{cb}

unfolded Belle data

$$B^- \rightarrow D^*(\rightarrow D\pi)\ell^-\bar{\nu}$$



Abdesselam et al., arXiv:1702.01521

CLN parametrization

Form factors entering helicity amplitudes (massless leptons)

$$h_{A_1}(w) = h_{A_1}(1)[1 - 8\rho^2 z + (r_{h2r}\rho^2 + r_{h2})z^2 + (r_{h3r}\rho^2 + r_{h3})z^3]$$

$$R_1(w) = R_1(1) + r_{11}(w - 1) + r_{12}(w - 1)^2$$

$$R_2(w) = R_2(1) + r_{21}(w - 1) + r_{22}(w - 1)^2 \quad w = v \cdot v'$$

Fixed:

$$r_{h2r} = 53, r_{h2} = -15, r_{h3r} = -231, r_{h3} = 91$$

$$r_{11} = -0.12, r_{12} = 0.05, r_{21} = 0.11, r_{22} = -0.06$$

Using this “tight” CLN parametrization

$$I = |\bar{\eta}_{EW} V_{cb}| h_{A_1}(1)$$

$$I_{\text{Belle}} = 0.0348(12) \quad (\text{unfolded})$$

$$I_{\text{HFLAV}} = 0.03561(11)(44)$$

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CLN uncertainties

$$h_{A_1}(w) = h_{A_1}(1)[1 - 8\rho^2 z + (r_{h2r}\rho^2 + r_{h2})z^2 + (r_{h3r}\rho^2 + r_{h3})z^3]$$

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$$R_2(w) = R_2(1) + r_{21}(w - 1) + r_{22}(w - 1)^2$$

Coefficients calculated through Λ/m using HQET & sum rules

$$r_{h2r} = 53, r_{h2} = -15, r_{h3r} = -231, r_{h3} = 91$$

BIG!

$$r_{11} = -0.12, r_{12} = 0.05, r_{21} = 0.11, r_{22} = -0.06$$

small!

Ratios

$$V(q^2) = \frac{R_1(w)}{r'} h_{A_1}(w) \quad A_2(q^2) = \frac{R_2(w)}{r'} h_{A_1}(w)$$

What are the uncertainties for the r 's? 20%? 100%?

See papers by Bigi, Gambino, Schacht; Grinstein & Kobach; Bernlochner et al.; Jaiswal, et al.

z-expansion

Series (z) expansion

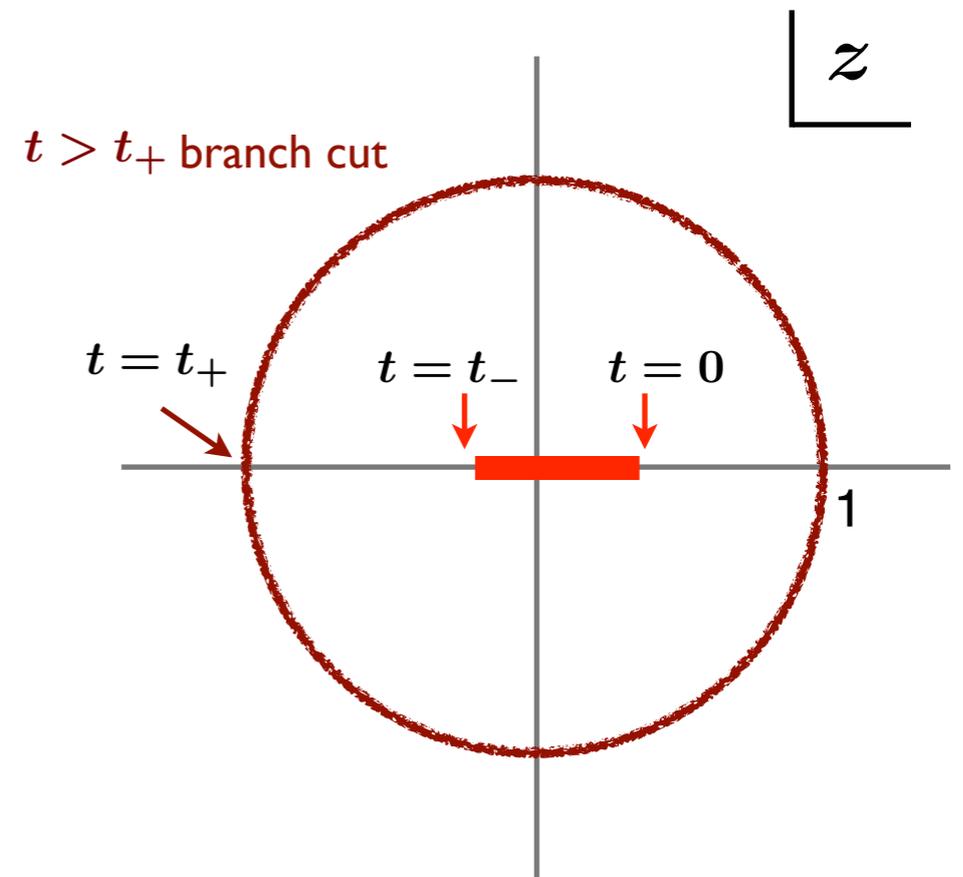
$$t = q^2 \quad t_{\pm} = (m_B \pm m_F)^2$$

Choose, e.g. $t_0 = t_-$

$$z = \frac{\sqrt{t_+ - t} - \sqrt{t_+ - t_0}}{\sqrt{t_+ - t} + \sqrt{t_+ - t_0}}$$

Simplified series expansion

$$F(t) = \frac{1}{1 - t/m_{\text{res}}^2} \sum_n a_n z^n$$



BGL parametrization

$$F(t) = Q_F(t) \sum_{k=0}^{K_F-1} a_k^{(F)} z^k(t, t_0) \quad Q_F(t) = \frac{1}{B_n(z)\phi_F(z)}$$

Blaschke factor

$$B_n(z) = \prod_{i=1}^n \frac{z - z_{P_i}}{1 - z z_{P_i}} \quad z_{P_i} = z(M_{P_i}^2, t_-)$$

Unitarity bounds

$$S_{fF} = \sum_{k=0}^{K_f-1} [(a_k^{(f)})^2 + (a_k^{(F_1)})^2] \leq 1 \quad S_g = \sum_{k=0}^{K_g-1} (a_k^{(g)})^2 \leq 1$$

Predictions for B_c vector & axial vector resonances

$$M_B + M_{D^*} = 7.290 \text{ GeV}$$

M_{1-}/GeV	method	Ref.	M_{1+}/GeV	method	Ref.
6.335(6)	lattice	[77]	6.745(14)	lattice	[77]
6.926(19)	lattice	[77]	6.75	model	[79, 80]
7.02	model	[79]	7.15	model	[79, 80]
7.28	model	[81]	7.15	model	[79, 80]

BCL parametrization

Simple form which uses less theoretical information.

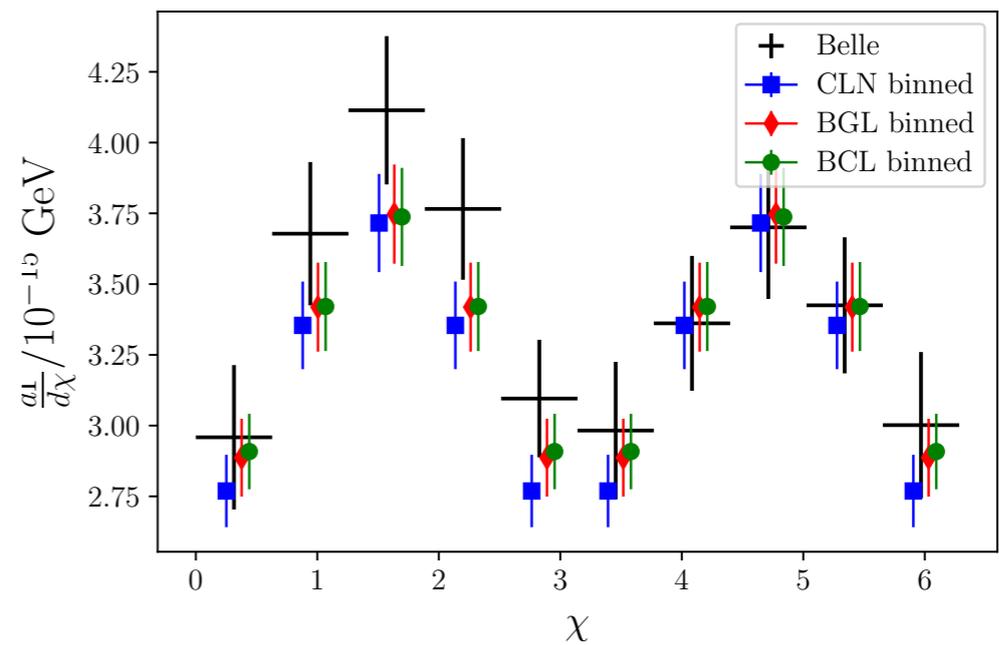
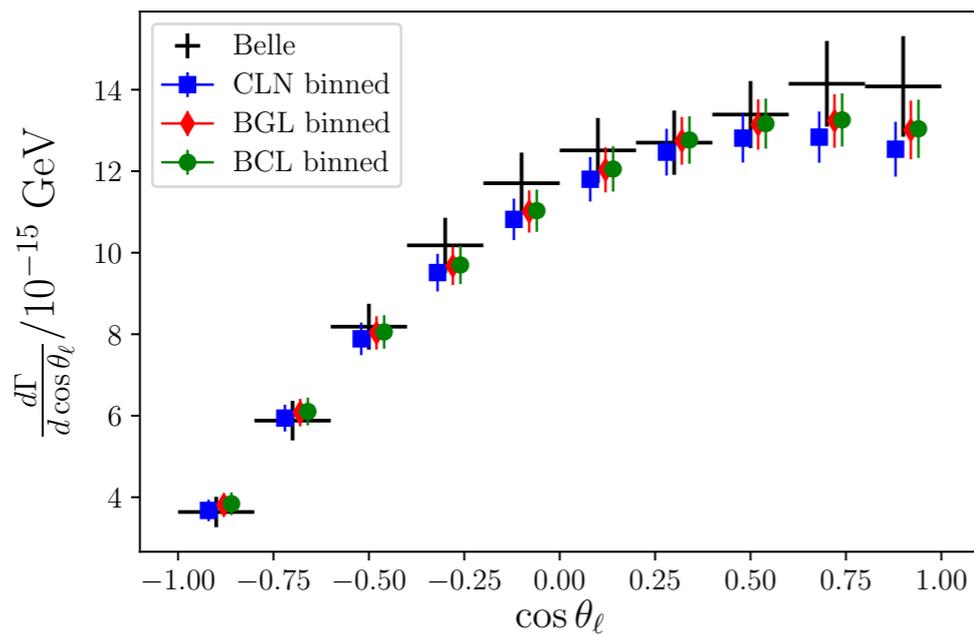
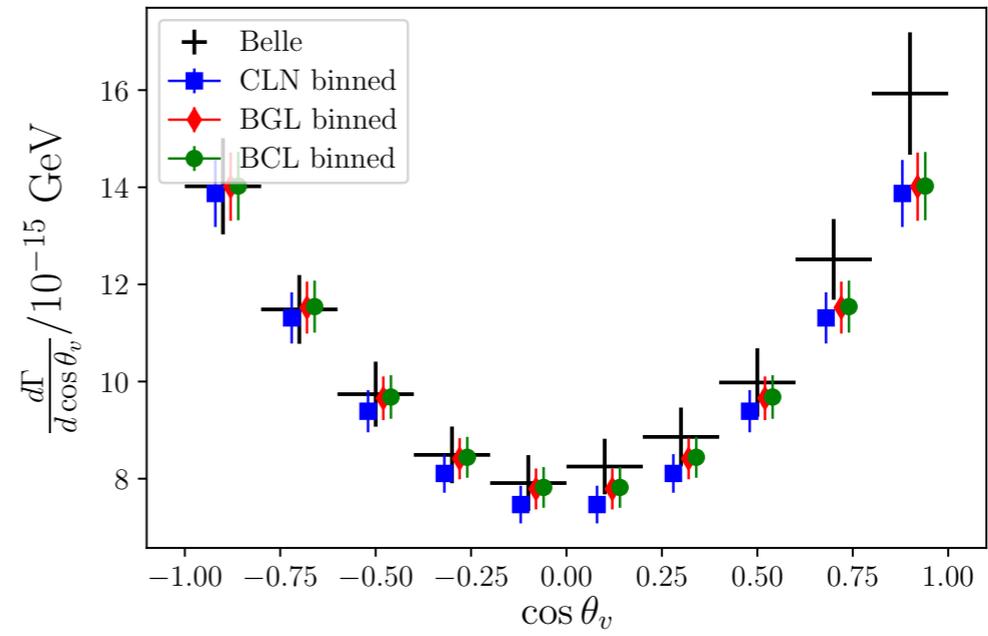
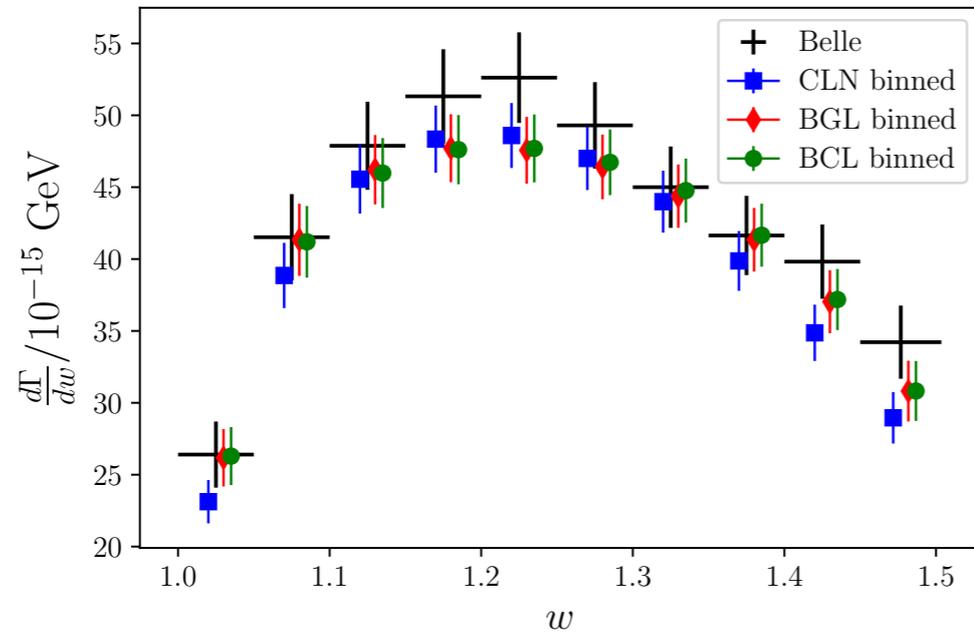
$$F(t) = Q_F(t) \sum_{k=0}^{K_F-1} a_k^{(F)} z^k(t, t_0) \quad Q_F(t) = \frac{N_F}{1 - \frac{t}{M_P^2}}$$

Using BGL as a guide, choose $N_f = 300$, $N_{F1} = 7000$, $N_g = 5$

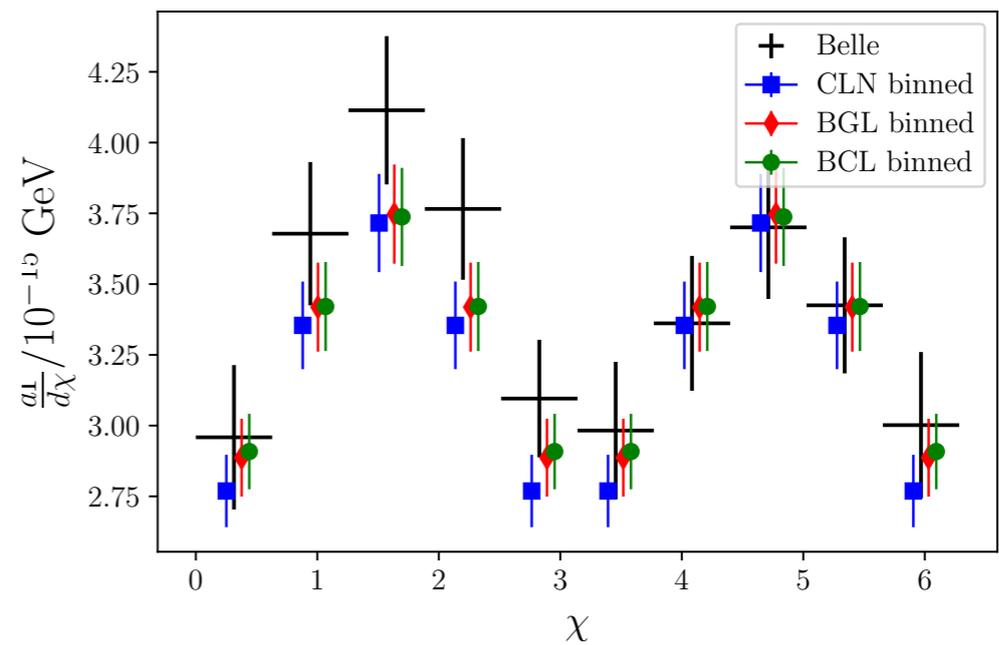
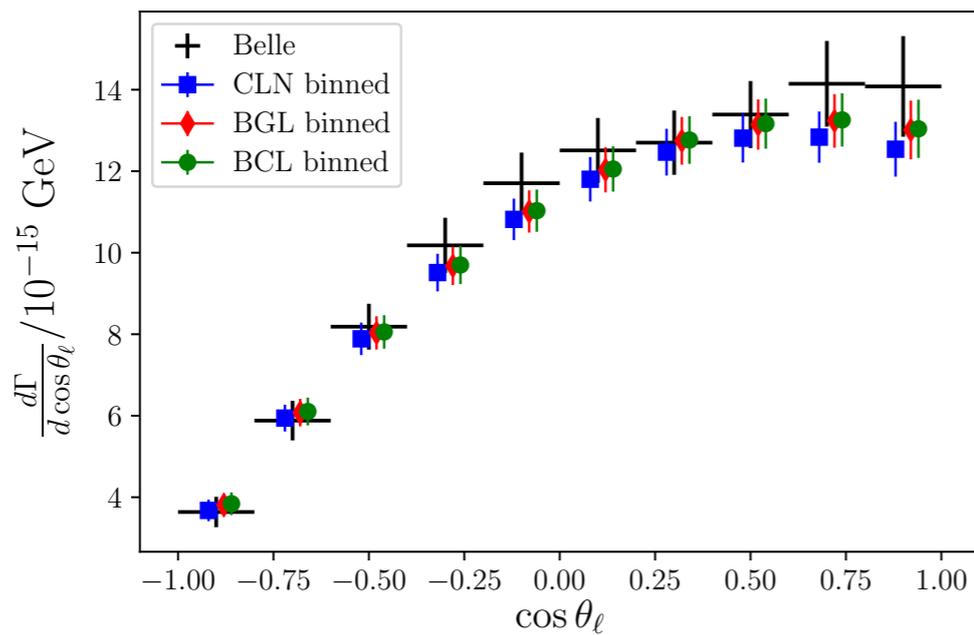
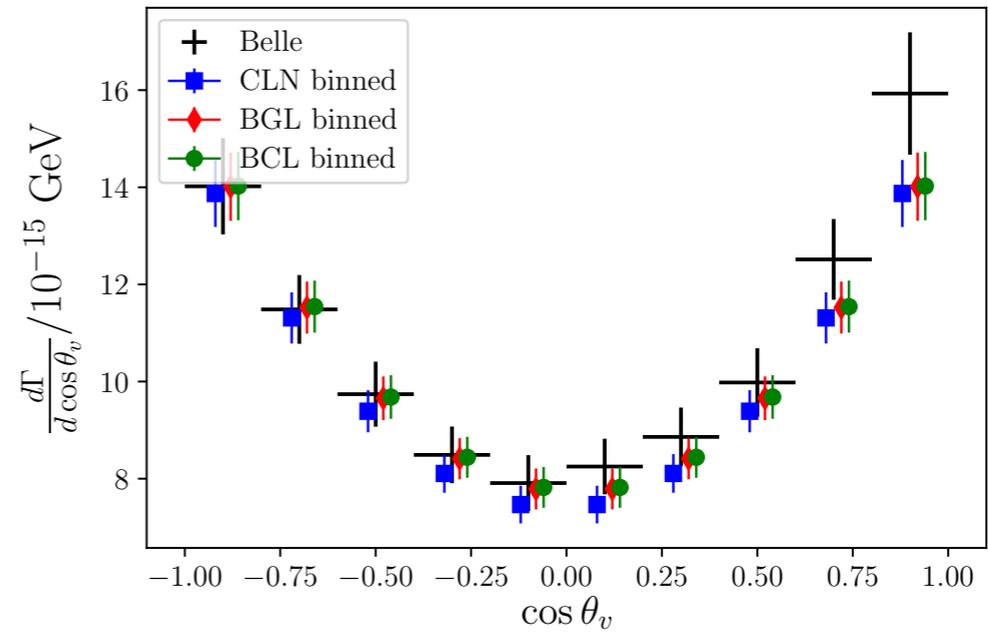
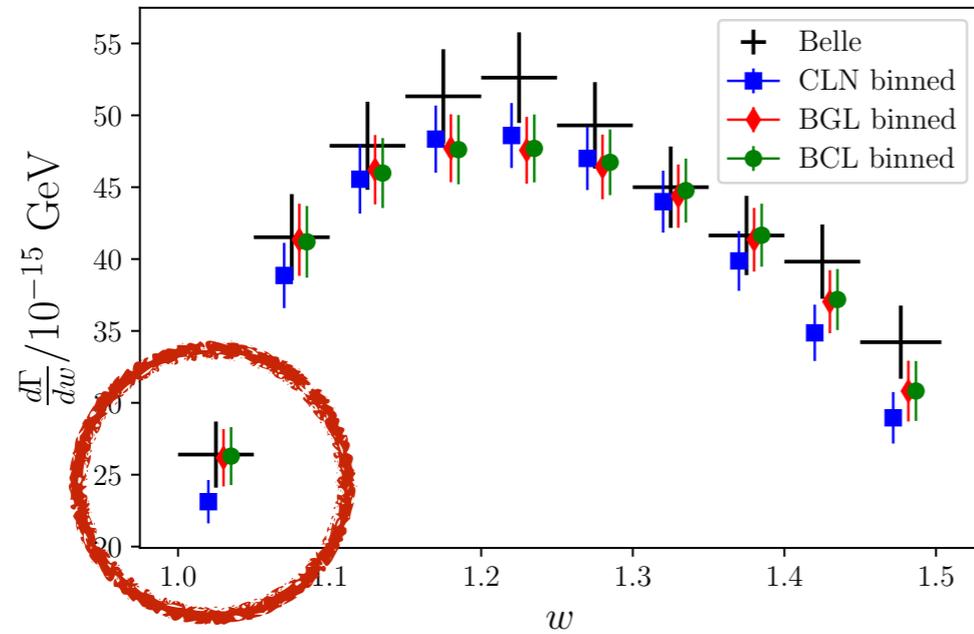
Clean baseline, against which affects of theoretical input (HQET, unitarity bounds) can be measured

fit	n_B^+	n_B^-	K	I	$a_0^{(f)}$	$a_1^{(f)}$	$a_0^{(F1)}$	$a_1^{(F1)}$	$a_0^{(g)}$	$a_1^{(g)}$	S_{fF}	S_g
BCL	-	-	2	0.0367(15)	0.01496(19)	-0.047(27)	0.002935(37)	-0.0029(27)	0.027(13)	0.77(44)	0.0025(26)	0.60(69)
BCL	-	-	3	0.0378(17)	0.01496(19)	-0.065(40)	0.002935(37)	-0.0135(82)	0.026(13)	0.82(46)	0.08(38)	0.67(75)
BCL	-	-	4	0.0382(18)	0.01497(19)	-0.310(42)	0.002936(37)	-0.0151(83)	0.109(16)	-0.29(37)	0.143(67)	0.10(22)
BCL	-	-	5	0.0382(18)	0.01497(19)	-0.310(42)	0.002936(37)	-0.0151(83)	0.109(16)	-0.29(37)	0.143(67)	0.10(22)

Fits to Belle data

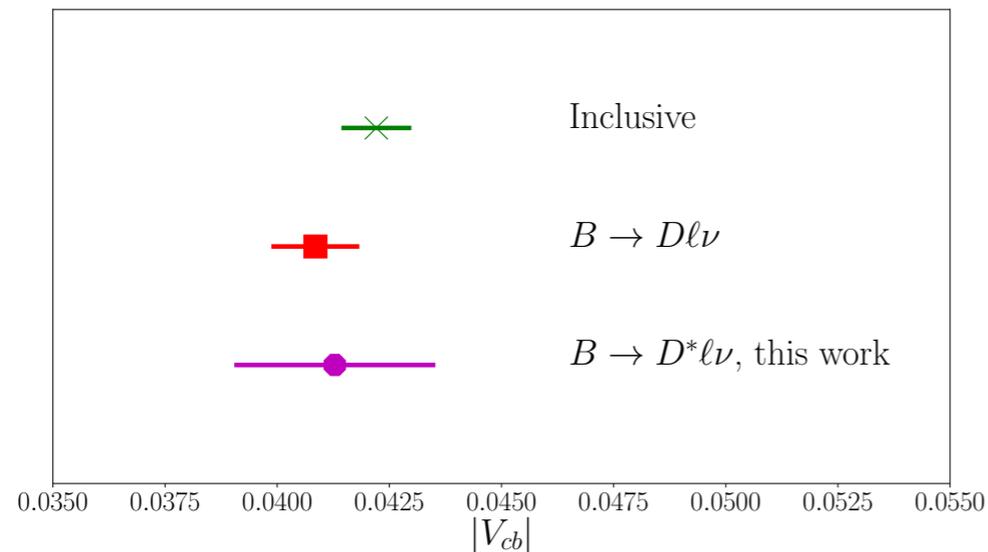
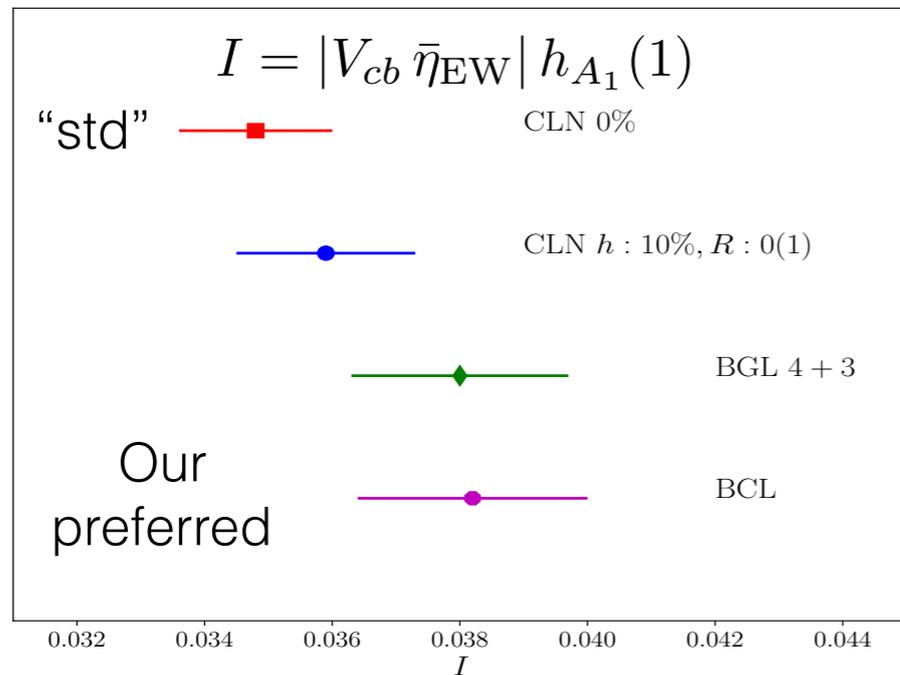


Fits to Belle data



Implications for V_{cb}

Different fit Ansätze



- Removal of theory assumptions resolves inclusive/exclusive tension, at least in Belle data
- Look forward to BaBar analysis
- Look forward to LQCD results at non-zero recoil

Conclusions

- Lattice field theory: nonperturbative, numerical approach connecting hadronic observables and fundamental quark interactions
- Lattice QCD plays an important role in studies of quark flavour
- Case study: $B \rightarrow D^* l \nu$
- Projects underway: more B semileptonic decay form factors, B mixing matrix elements, ...

back-up

NRQCD matching

$$\langle \mathcal{J}^i \rangle = (1 + \alpha_s(\eta - \tau)) \langle J_{\text{latt}}^{(0)i} \rangle + \langle J_{\text{latt}}^{(1)i} \rangle + e_4 \frac{\Lambda_{\text{QCD}}^2}{m_b^2}$$

$$J_{\text{latt}}^{(0)i}(x) = \bar{c} \gamma^i \gamma^5 Q$$

$$J_{\text{latt}}^{(1)i}(x) = -\frac{1}{2am_b} \bar{c} \gamma^i \gamma^5 \gamma \cdot \Delta Q$$

1-loop coefficients η & τ from Monahan, Shigemitsu, Horgan, PRD87 (2013)

Truncation errors enter at order: $\frac{\Lambda_{\text{QCD}}^2}{m_b^2}$ included as Gaussian noise

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$\alpha_s \tau \langle J_{\text{latt}}^{(0)} \rangle$	$\langle J_{\text{latt}}^{(1)} \rangle$		
0.00559(8)	0.0078(66)	very coarse	& physical sea quark masses
0.0064(1)	0.0055(48)	coarse	
<u>0.0080(9)</u>	<u>0.0048(6)</u>	fine	

Cancellation expected from Luke's theorem

Chiral-continuum fit

Fit function:

$$\begin{aligned}
 h_{A_1}(1) = & \overset{\text{Static}}{\underset{\text{limit}}{(1 + \delta_a^B)}} B + \overset{\delta_a: \text{disc}}{\underset{\text{errors}}{\delta_a^g}} \frac{g^2}{48\pi^2 f^2} \times \text{chiral logs} + \overset{\text{light quark}}{\underset{\text{mass}}{C}} \frac{M_\pi^2}{\Lambda_\chi^2} \\
 & + e_1 \alpha_s^2 \left[1 + e_5 (am_b - 2)/2 + e_6 ((am_b - 2)/2)^2 \right] J_{\text{latt}}^{(0)} \\
 & \text{2-loop matching error}
 \end{aligned}$$

with $g^2 = 0.53(8)$

The α_s^2 uncertainty is the largest, by a factor of 2, compared to others