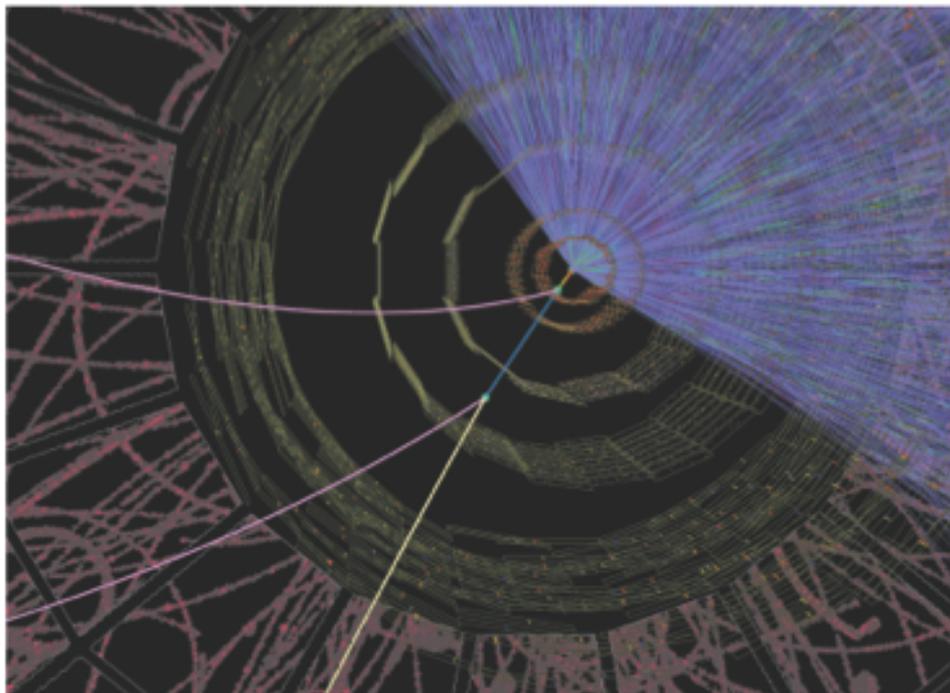




# $K_S^0$ & $\Lambda$

# Production in ALICE



Luke Hanratty

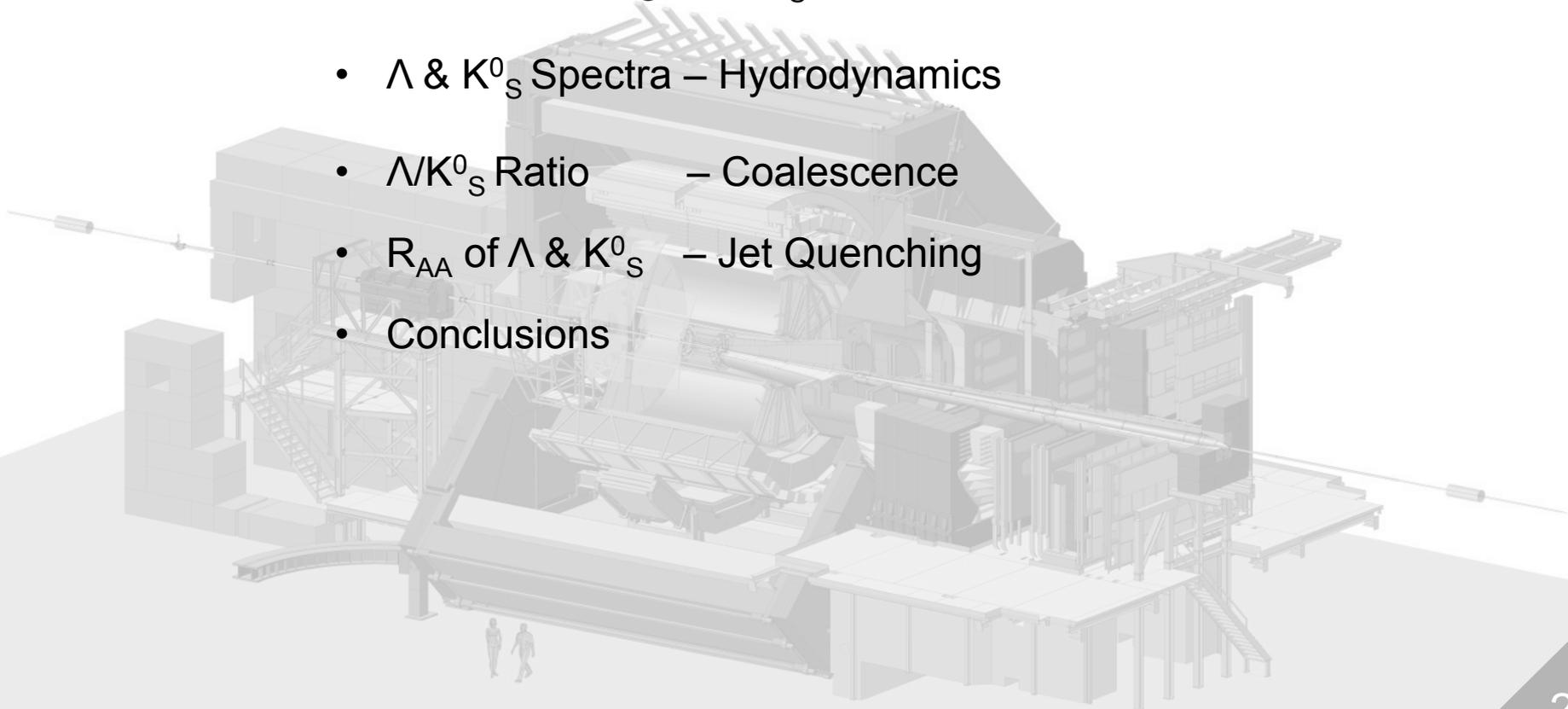
University of Birmingham

for the ALICE Collaboration



# Outline

- The ALICE Experiment
- Quark-Gluon Plasma
- Reconstructing  $\Lambda$  &  $K_S^0$
- $\Lambda$  &  $K_S^0$  Spectra – Hydrodynamics
- $\Lambda/K_S^0$  Ratio – Coalescence
- $R_{AA}$  of  $\Lambda$  &  $K_S^0$  – Jet Quenching
- Conclusions



# ALICE

TPC:

(Time Projection Chamber)  
tracking & PID

$|\eta| < 0.9, 85 < r < 247 \text{ cm}$

$\sigma_p/p \sim 0.7\% (1 \text{ GeV}/c)$

$\sigma_p/p \sim 6.5\% (10 \text{ GeV}/c)$

$\sigma_{dE/dx}/dE/dx \sim 6.5\%$

ITS:

(Silicon strips / pixels)  
tracking & triggering

$|\eta| < 0.9, 4 < r < 44 \text{ cm}$

$\sigma_{vertex} \sim 100 \mu\text{m}$

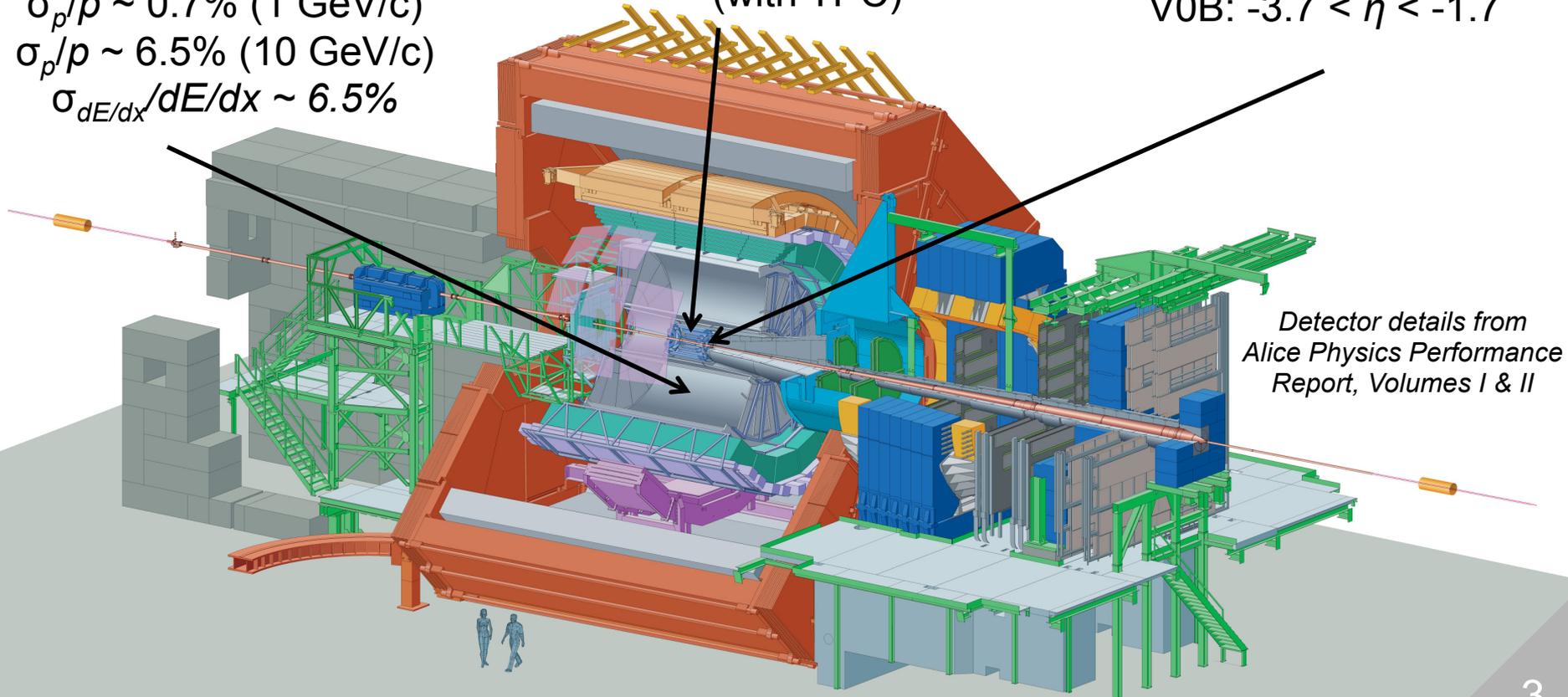
$\sigma_p/p \sim 3.5\% (100 \text{ GeV}/c)$   
(with TPC)

VZERO:

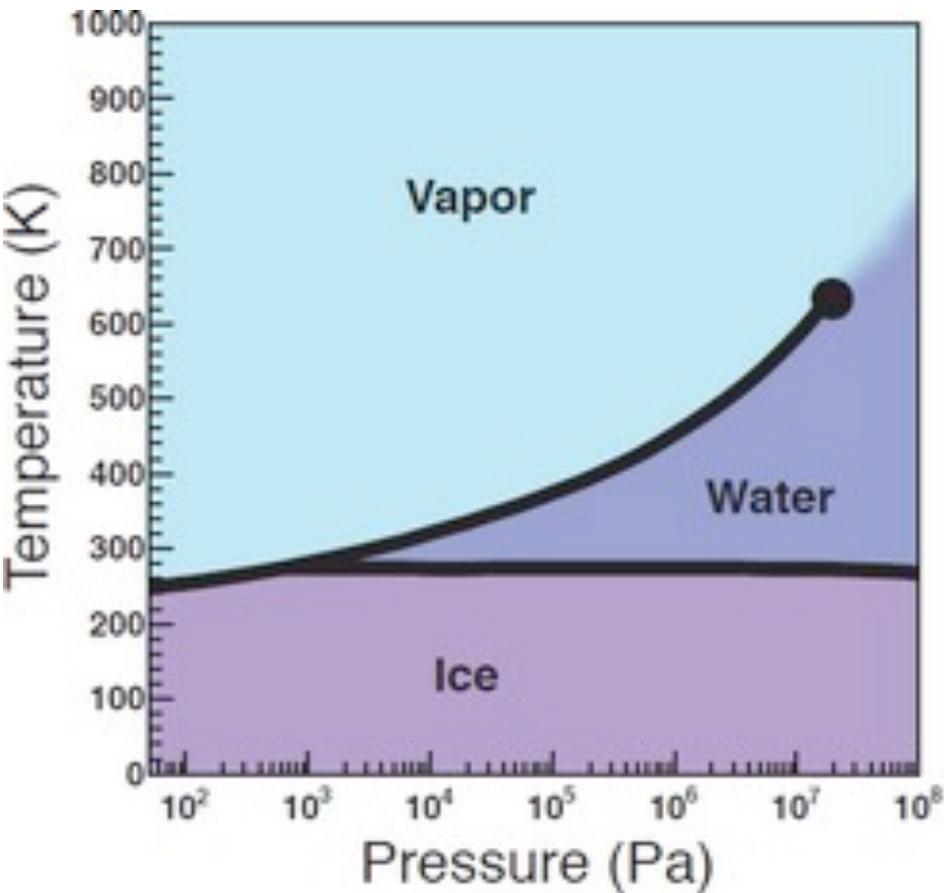
(Scintillator Arrays)  
centrality selection &  
triggering

V0A:  $2.8 < \eta < 5.1$

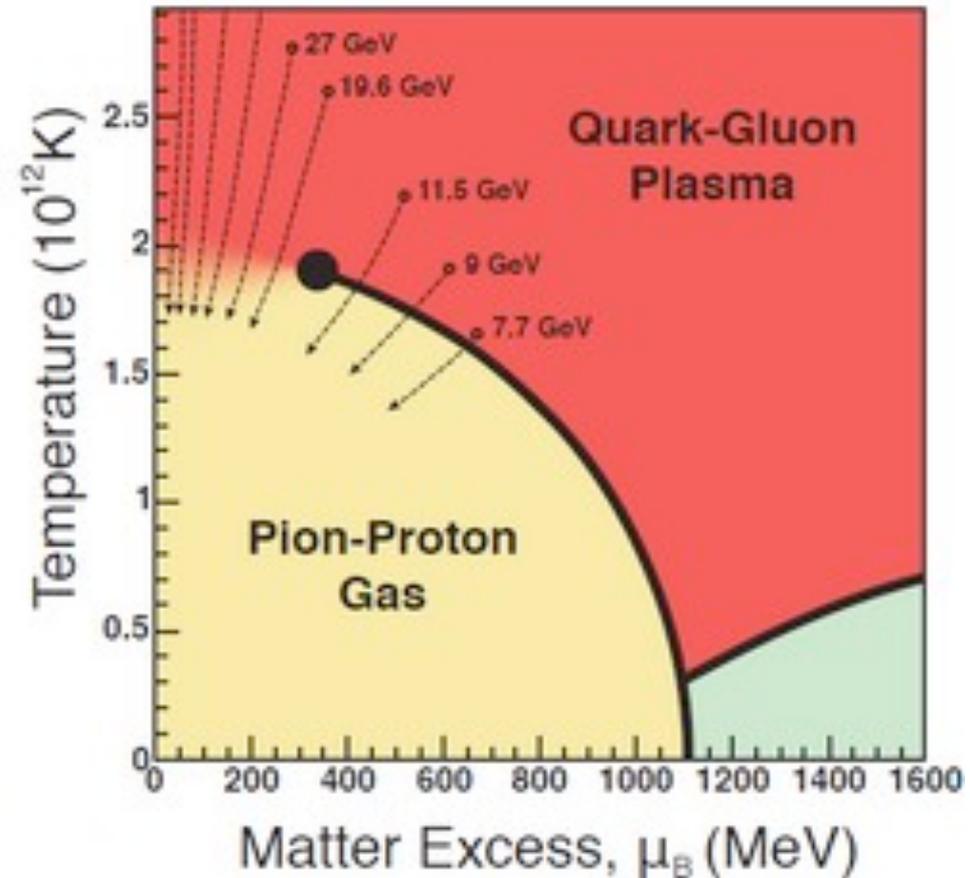
V0B:  $-3.7 < \eta < -1.7$



# QGP

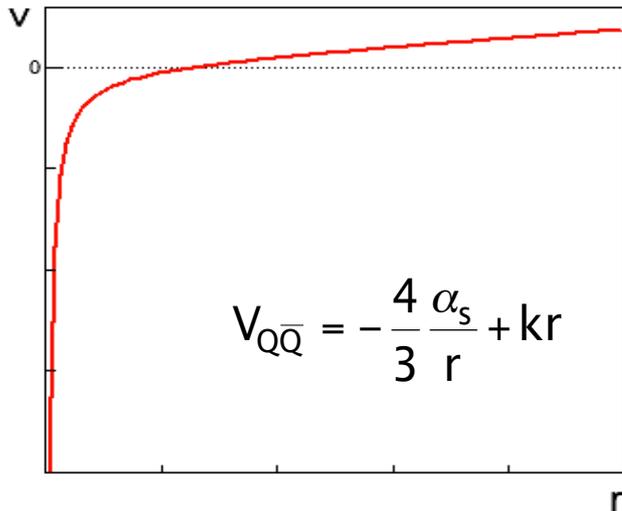


<http://alicematters.web.cern.ch/?q=content/node/515>

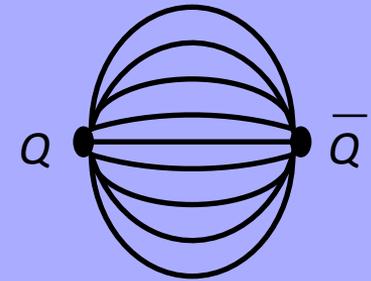


# QGP

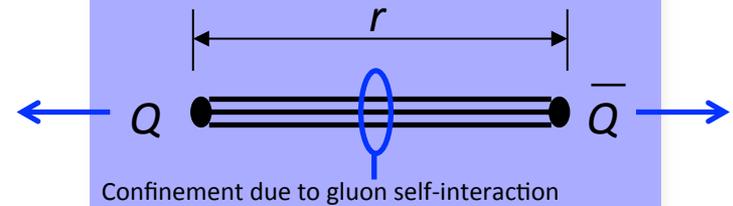
- Color charges are usually considered to be **confined**



a) QED or QCD ( $r < 1$  fm)



b) QCD ( $r > 1$  fm)

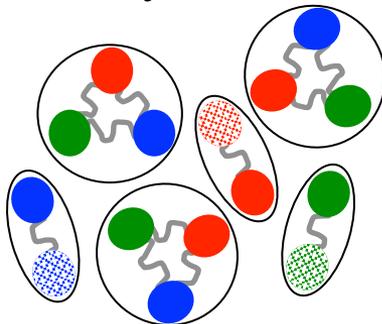


- But in a dense medium, the charge is screened (**Debye Screening**)

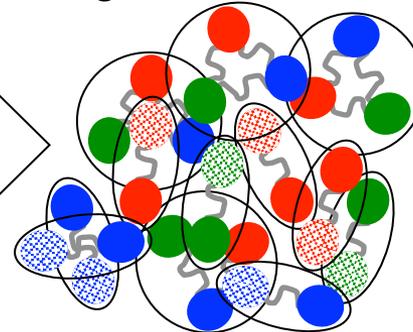
$$V_{Q\bar{Q}} \approx -\frac{4}{3} \frac{\alpha_s}{r} \exp\left(-\frac{r}{r_D}\right)$$

- So in a sufficiently hot or dense medium, the charges can be effectively free

**Hadron  
Gas**



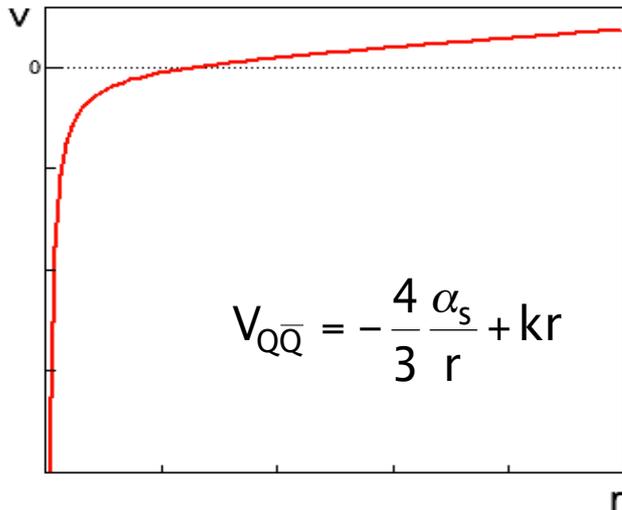
Heat,  
Compression



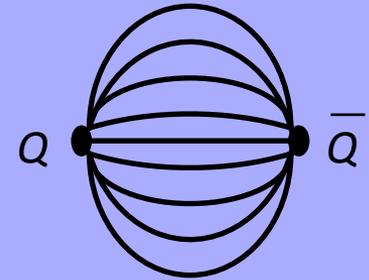
?

# QGP

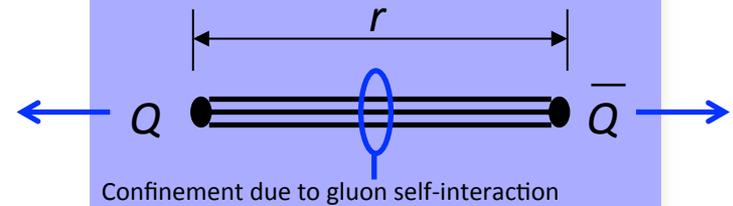
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a) QED or QCD ( $r < 1$  fm)



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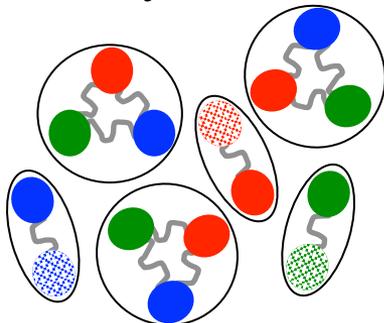


- But in a dense medium, the charge is screened (**Debye Screening**)

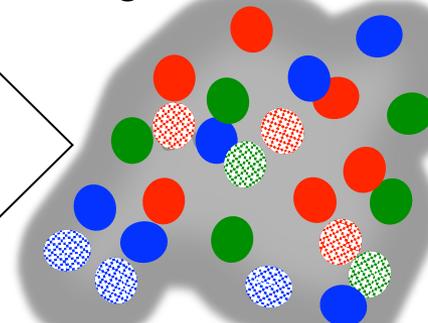
$$V_{Q\bar{Q}} \approx -\frac{4}{3} \frac{\alpha_s}{r} \exp\left(-\frac{r}{r_D}\right)$$

- So in a sufficiently hot or dense medium, the charges can be effectively free

**Hadron Gas**



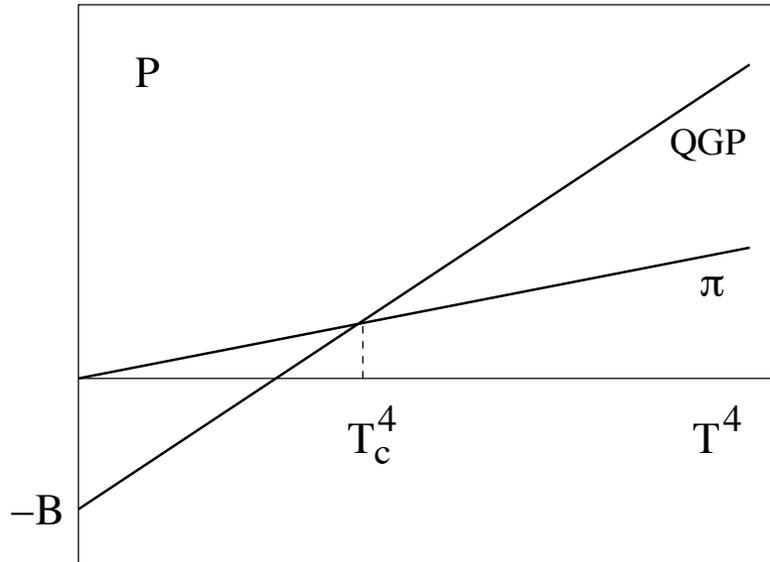
Heat,  
Compression



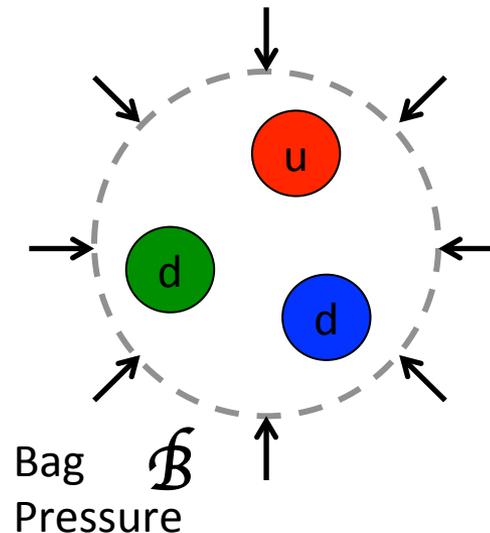
**Quark-gluon Plasma**

# QGP

- Quark-gluon plasma is a **thermalised** system of **deconfined** quarks and gluons
- As a first approximation, we can treat Hadron Gas / QGP as an ideal gas:



Satz, Helmut arXiv:0903.2778



$$P = g \frac{\pi^2}{90} T^4$$

$$g_{\text{HadronGas}} \sim 3$$

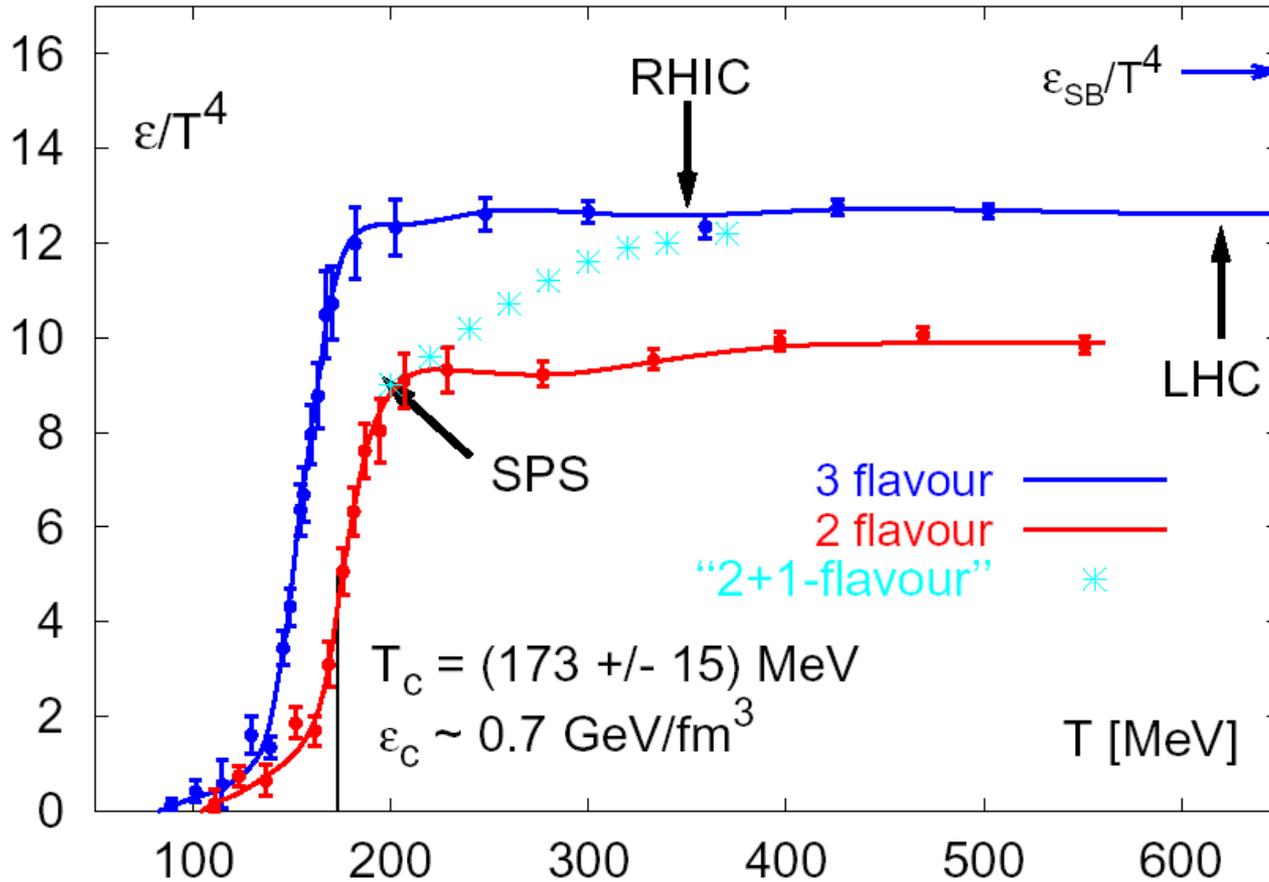
$$g_{\text{QGP}} \sim 48$$

A deconfined state can be attained if the QGP pressure is greater than the Bag pressure and the Hadron Gas pressure

→ Transition temperature  $\sim 140\text{MeV}$

# QGP

- We can test this prediction using Lattice QCD



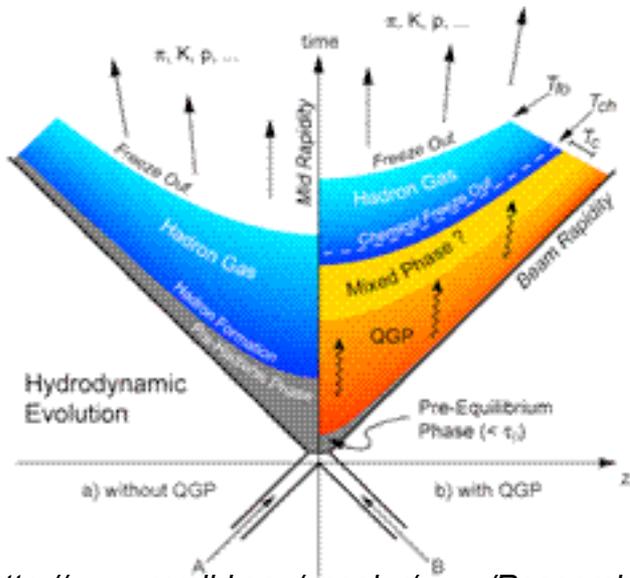
F Karsch: Quark Gluon Plasma 3 (World Scientific)

$$\frac{\epsilon}{T^4} = g \frac{\pi^2}{30}$$

$$g_{\text{HadronGas}} \sim 3$$

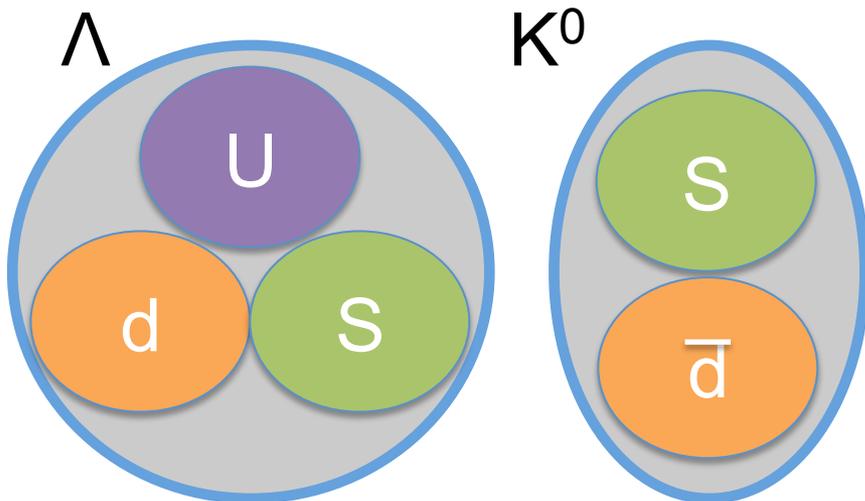
$$g_{\text{QGP}} \sim 48$$

# Motivation

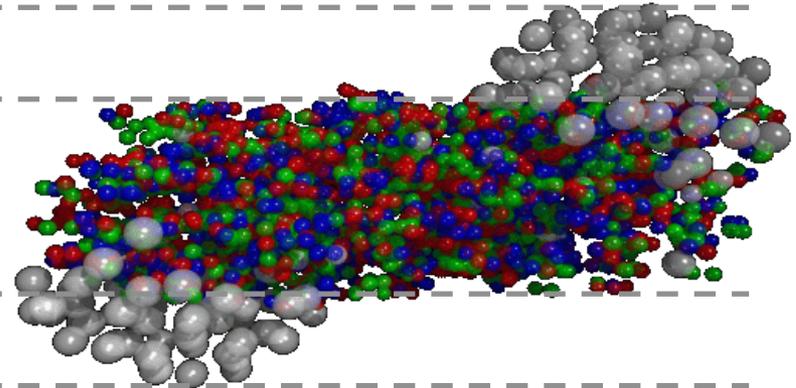
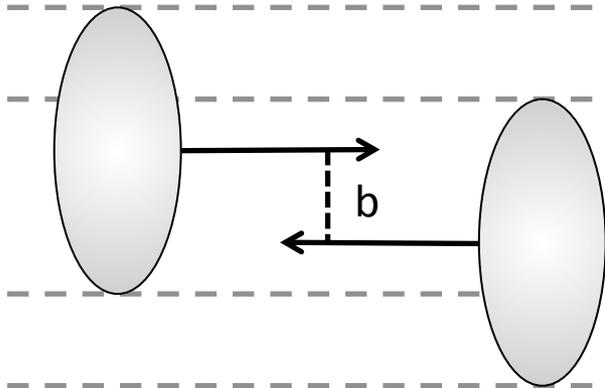


- The QGP cools very rapidly in the lab  
 $\rightarrow \tau \sim 11 \text{ fm}/c \sim 10^{-23} \text{ s}$
- Most information in Hadron Spectra at freezeout
- $\Lambda$  &  $K_S^0$  are the lightest strange baryon & meson
- Studying different hadrons allows understanding of mass & flavour dependence of Jet Quenching, Flow...

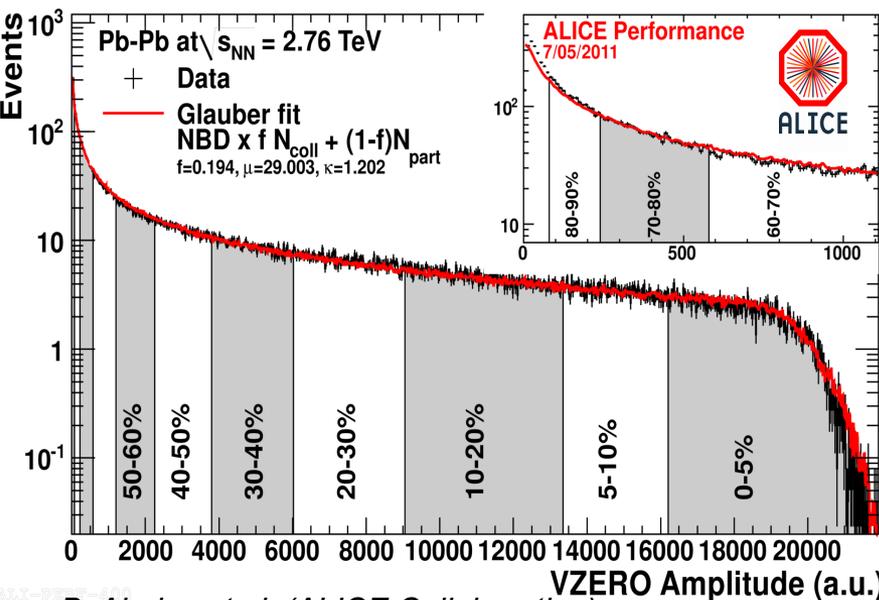
<http://www-rnc.lbl.gov/~ssalur/www/Research3.html>



# Centrality

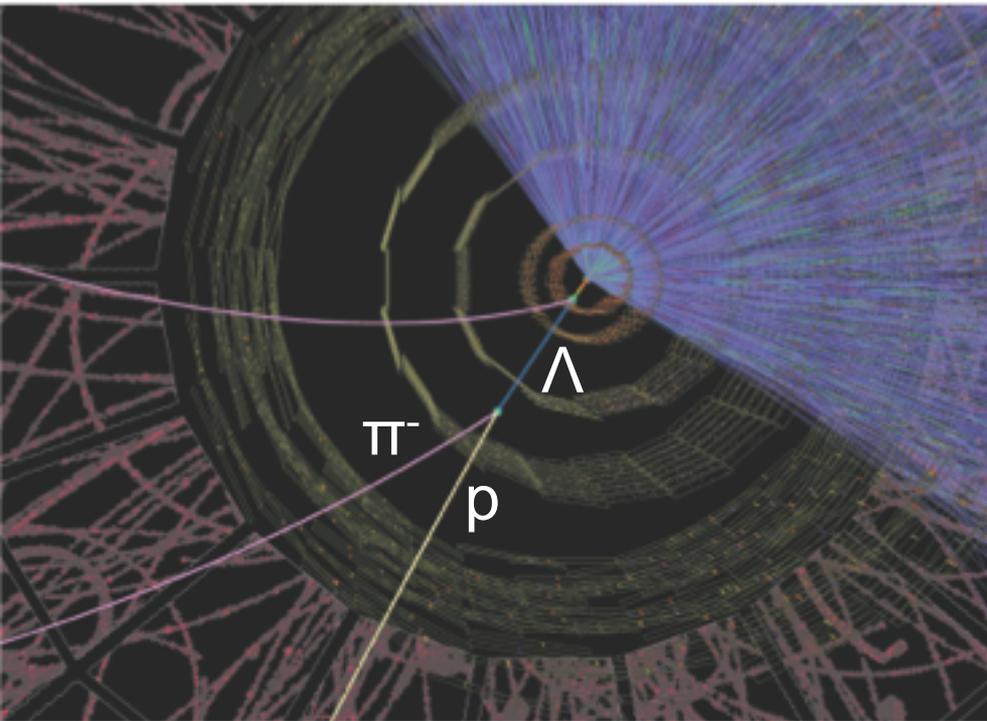


- Experimentally, we need to distinguish between head-on collisions, and glancing collisions.



- Events are classified by their **centrality** – a % of the total nuclear interaction cross section
- The number of charged particles produced in a collision is dependent on both  $N_{participants}$  and  $N_{collisions}$
- The number of charged particles seen at the VZERO is fitted with a prediction from MC, and split into centrality regions

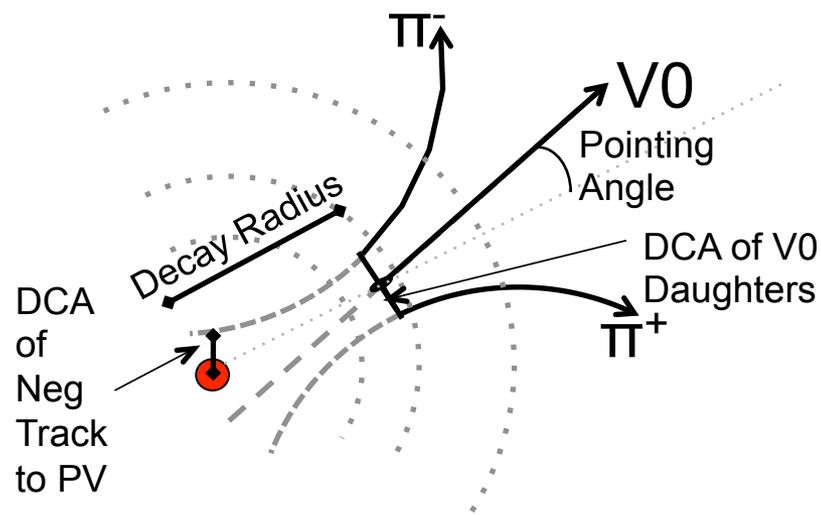
# Reconstruction – V0s



Pb-Pb 5.5TeV Hijing MC Event, not all tracks shown  
 Alice Physics Performance Report, Volume II (Figure IV)

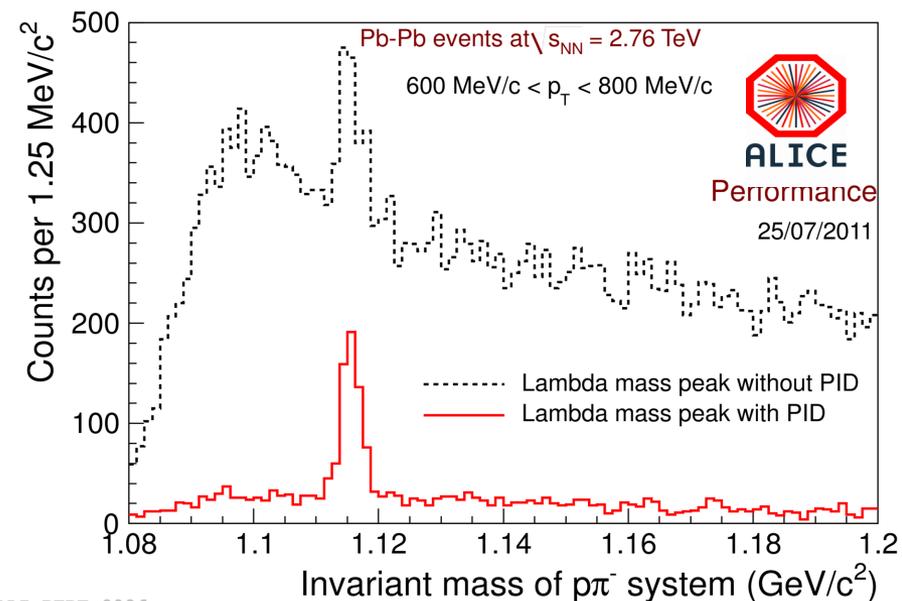
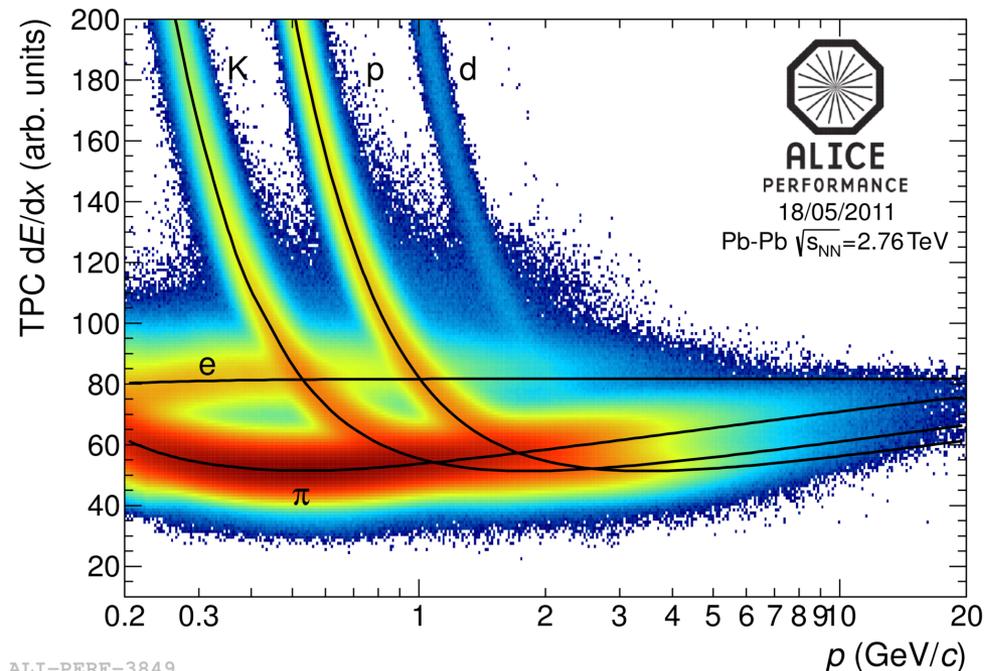
True 'V0's can be distinguished from combinatorial background by geometrical cuts on:

- DCA of daughters to each other
- DCA of daughters to primary vertex
- Cosine of pointing angle
- Decay radius



Decay	Branching Ratio
$K_S^0 \Rightarrow \pi^+\pi^-$	69.2%
$\Lambda \Rightarrow p\pi^-$	63.9%

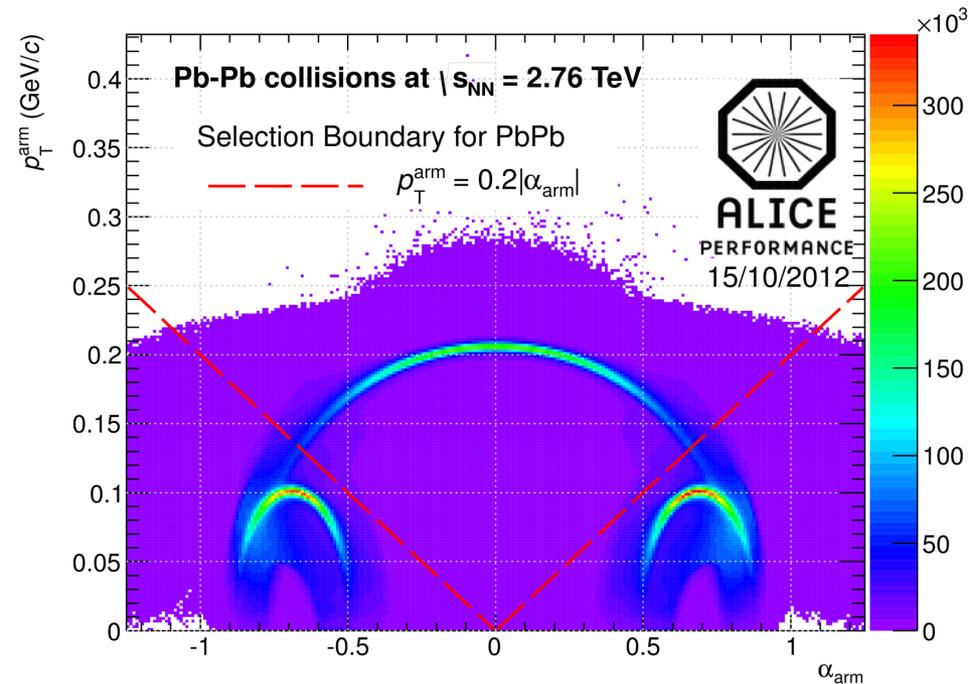
# Reconstruction - PID



At low  $p_T$ ,  $dE/dx$  information from the TPC allows separation of the signal from the pion dominated combinatorial background.

- $3\sigma$  cut on TPC  $dE/dx$  for proton daughters.

# Reconstruction – Armenteros Plot



The Armenteros-Podolanski diagram is essentially another way of visualizing the invariant mass distribution of V0s

Momentum of the daughters relative to the mother are plotted, where:

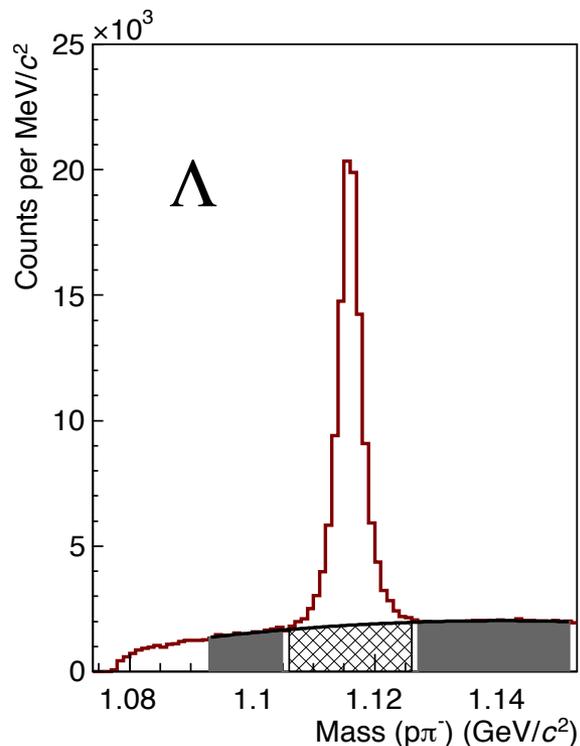
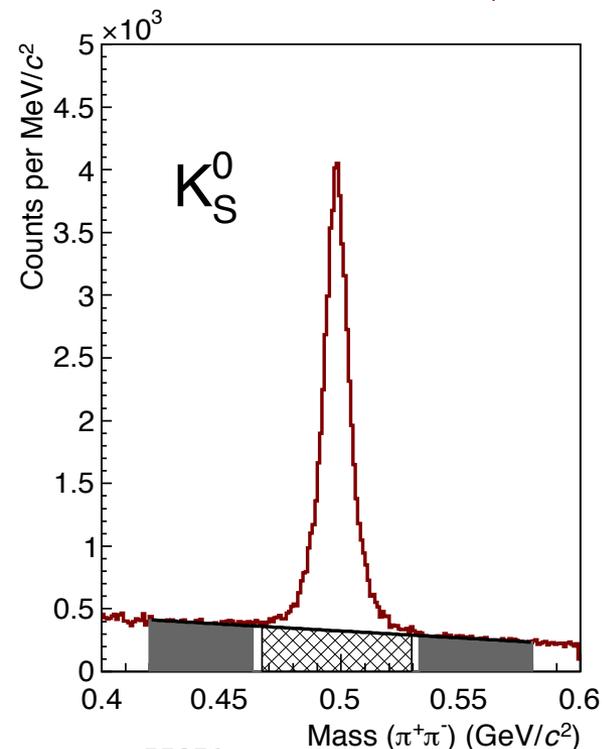
$$\alpha = \frac{p_{\parallel}^{+} - p_{\parallel}^{-}}{p_{\parallel}^{+} + p_{\parallel}^{-}}$$

A cut on this diagram was used to remove  $\Lambda$  particles from the  $K_S^0$  background

# Reconstruction – Mass Peaks

Pb-Pb at  $\sqrt{s_{NN}} = 2.76$  TeV,  $|y| < 0.5$

$3.0 < p_T < 3.2$  GeV/c, 0-5% centrality



- Peak fitted with Gaussian + 2<sup>nd</sup> degree polynomial
- ‘Sideband regions’ fitted with 2<sup>nd</sup> degree polynomial
- Signal defined as the difference between the counts in the ‘peak region’ and the background fit

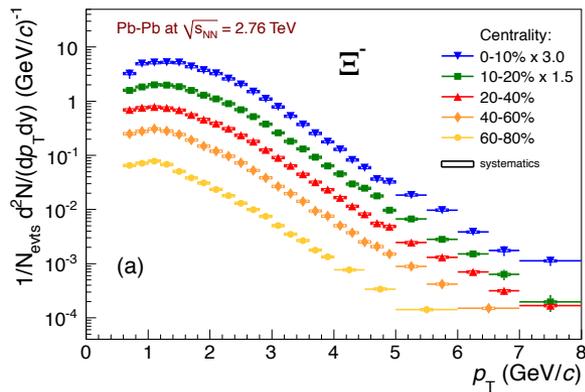
ALI-PUB-55071

ALICE Collaboration, (2013), *arXiv:nucl-ex/1307.5530*

# Reconstruction - Feeddown

Decay	Branching Ratio
$\Xi^- \Rightarrow \Lambda \pi^-$	$99.887 \pm 0.035\%$
$\Xi^0 \Rightarrow \Lambda \pi^0$	$99.525 \pm 0.012\%$

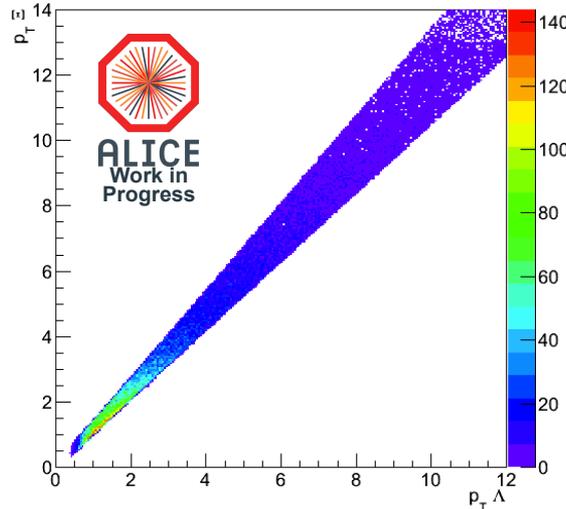
- $\Lambda$  coming from weak decays of  $\Xi$  are removed



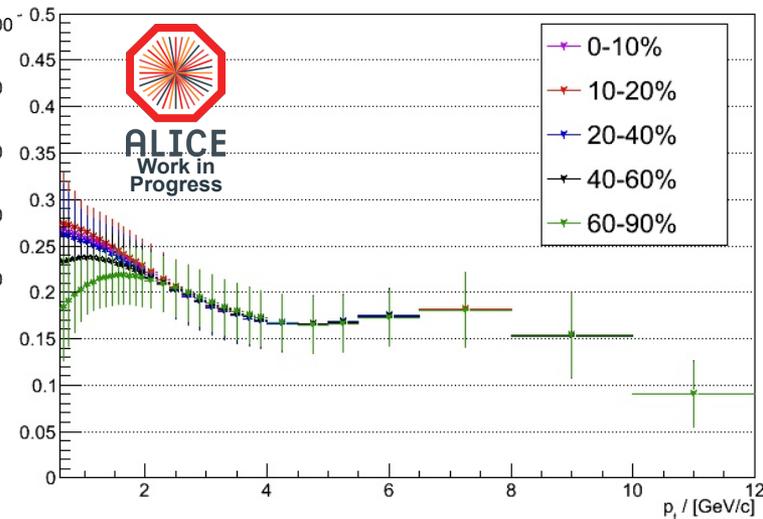
Measured  $\Xi$  Spectra

B. Abelev et al. (ALICE Collaboration), *Physics Letters B* 728 (2014) 216–227

MC Feeddown Matrix

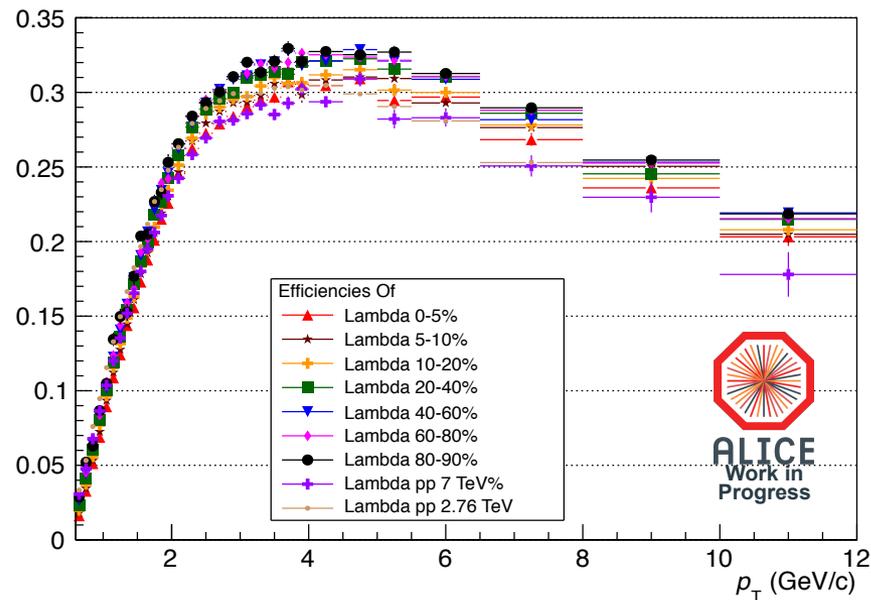
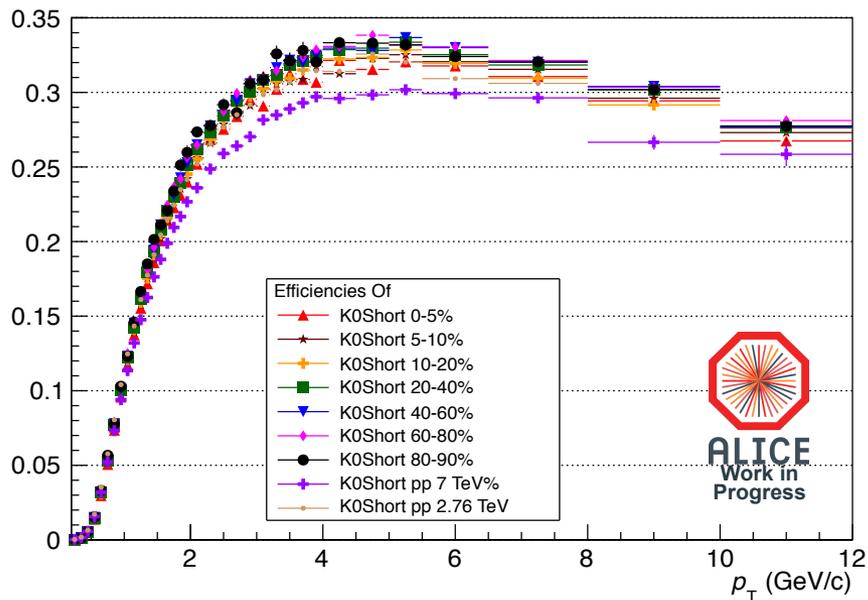


Feeddown Fraction



- $\Omega$  decay also considered, but found to be negligible

# Efficiencies

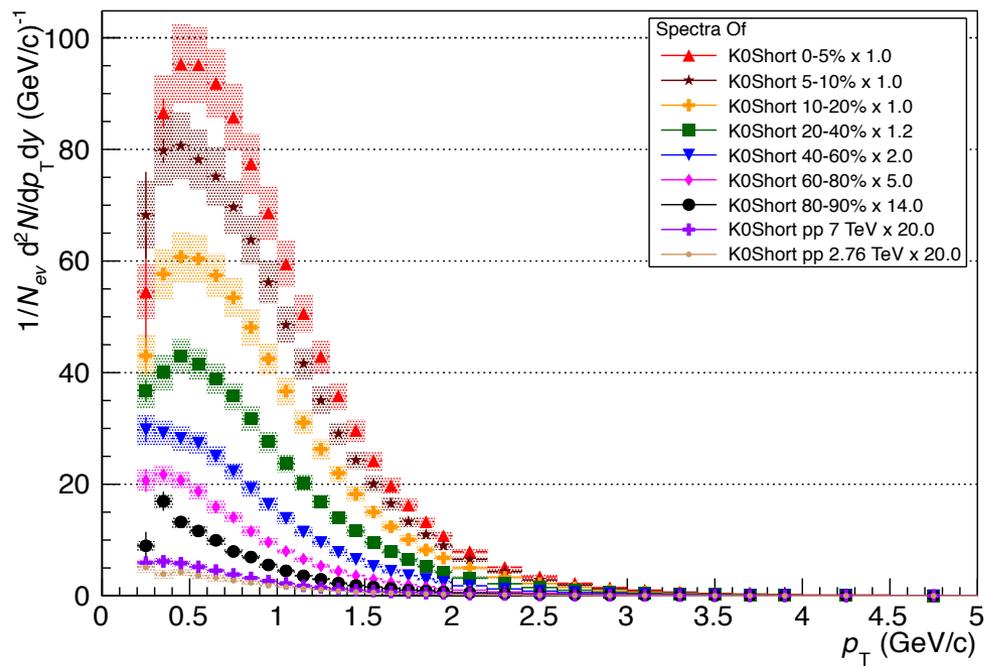
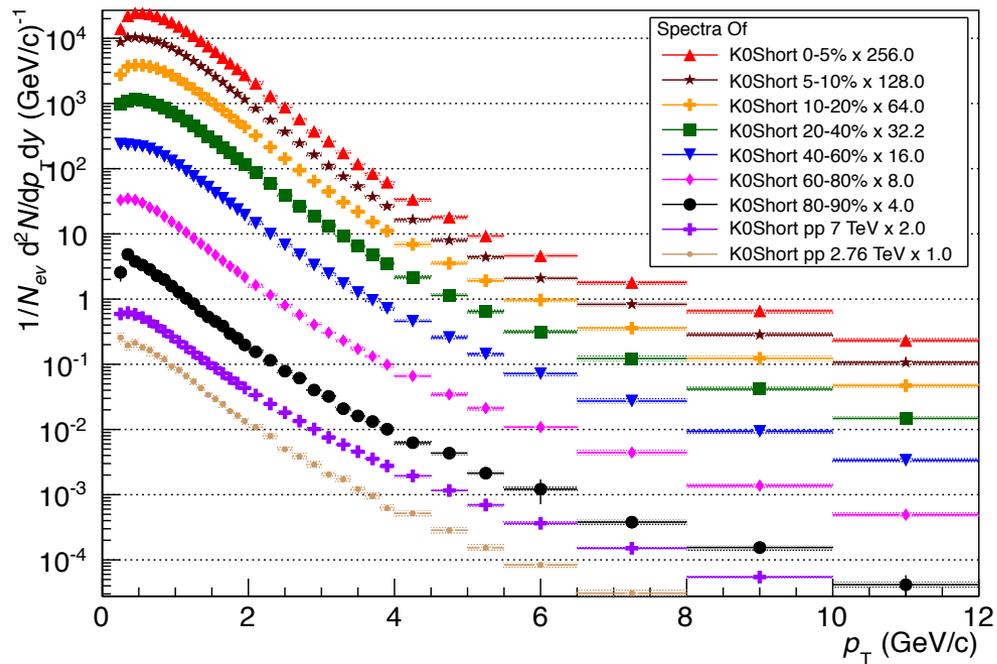


- Reconstruction efficiency  $\sim 20\text{-}30\%$
- Low  $p_T$  is more difficult  $\rightarrow$  low  $p_T$  cutoff;  $\Lambda > 0.6$  GeV/c,  $K_S^0 > 0.4$  GeV/c
- Slight variations with centrality,  $\sqrt{s}$ , but largely consistent

# Spectra

 $K^0_S$ 

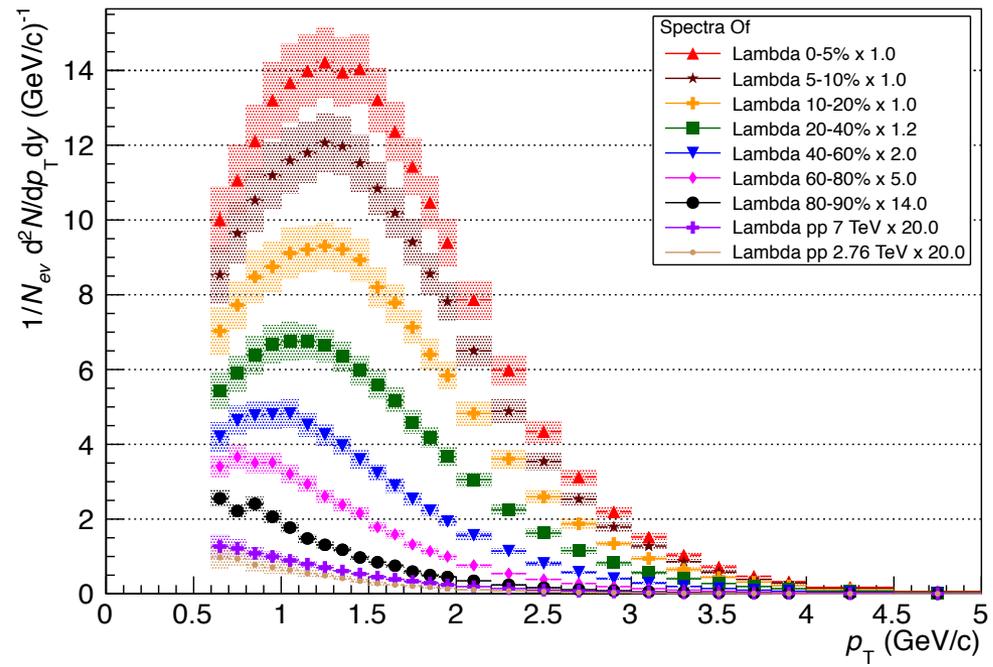
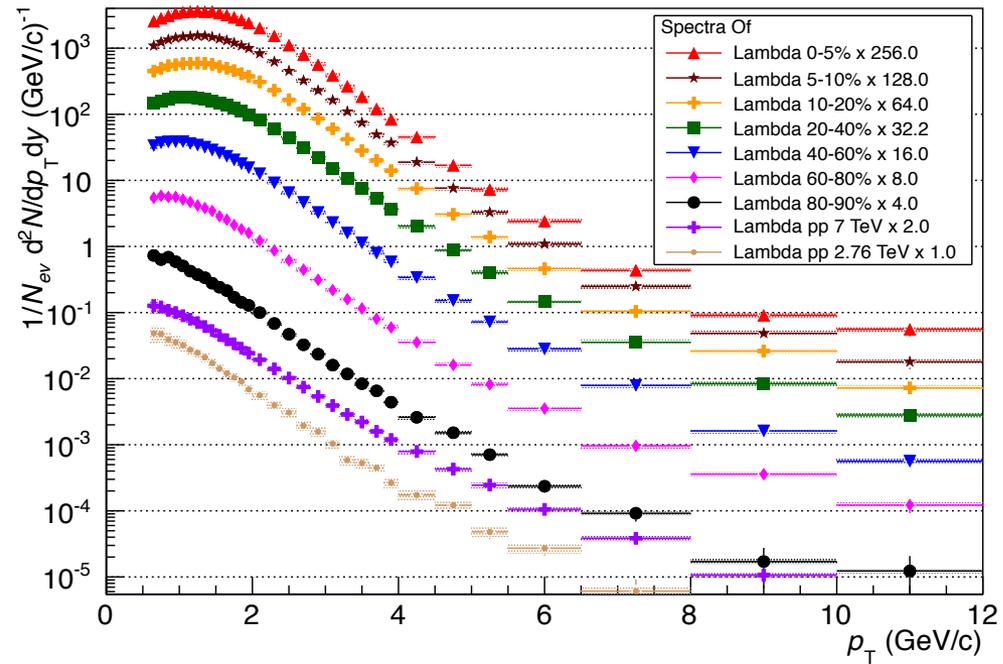
- Plots scaled to separate centralities; ordering unchanged



# Spectra

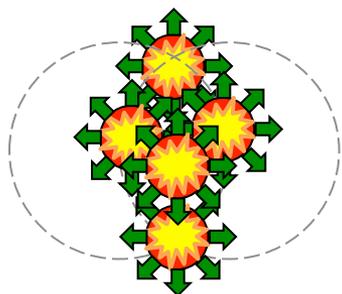


- Plots scaled to separate centralities; ordering unchanged

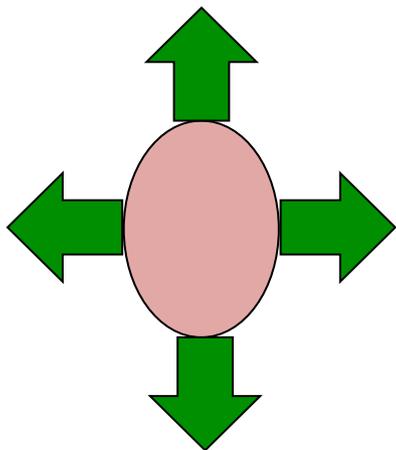


# Hydrodynamics

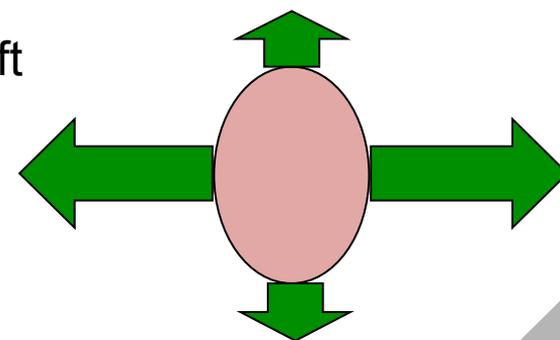
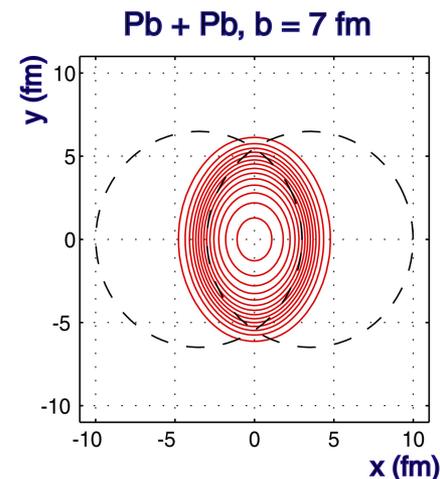
## Superposition of independent p-p



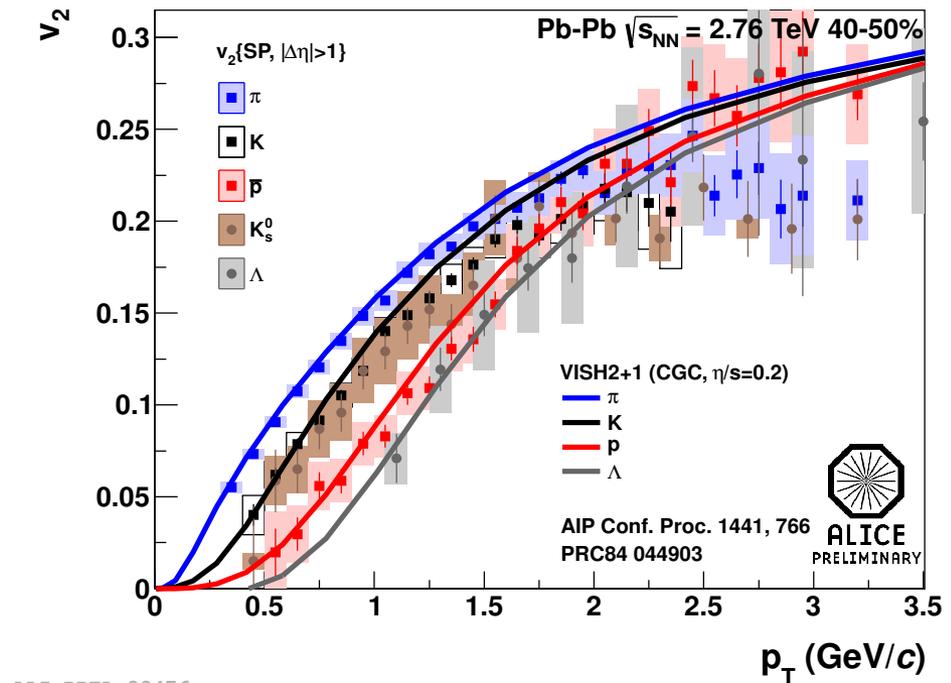
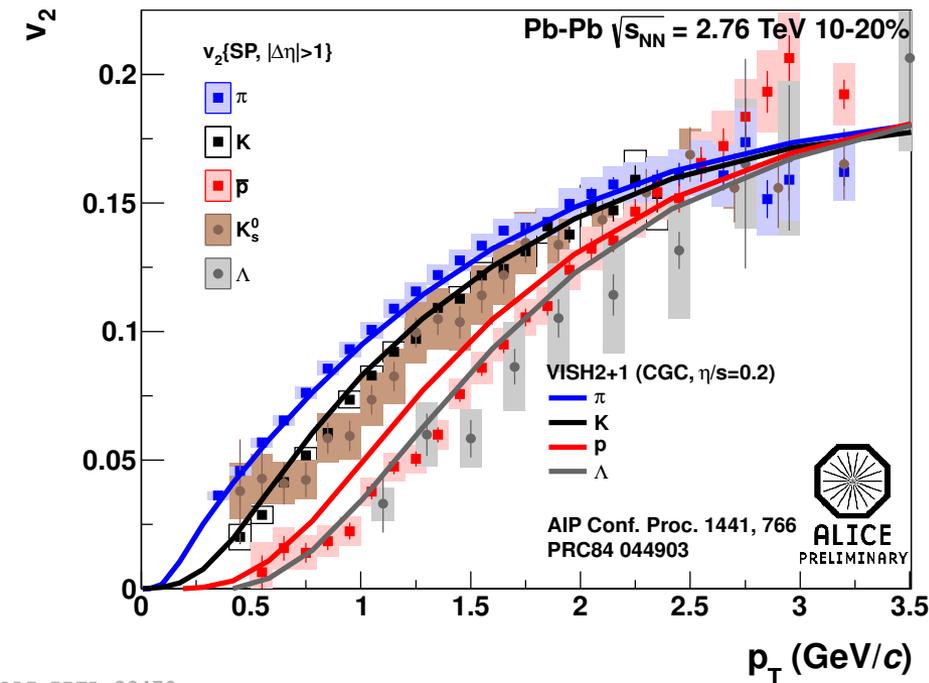
- If a thermalised medium is formed, one can have collective flow
- Initial spatial anisotropy leads final momentum anisotropy
- Typically studied by breaking down into Fourier components  
→ radial, elliptical...
- Radial flow causes a Doppler-shift of the momentum distributions of particles – larger for heavier particles



## Collective Behaviour - flow



# Hydrodynamics



ALI-PREL-28470

ALI-PREL-28476

- $v_2$  is a measure of the elliptical flow of the system
- Viscous hydrodynamics predicts its behaviour well

$$E \frac{d^3N}{d^3p} = \frac{1}{2\pi} \frac{d^2N}{p_T dp_T dy} \left( 1 + \sum_{n=1}^{\infty} 2v_n \cos[n(\phi - \Psi_r)] \right) \quad v_2 = \langle \cos[2(\phi - \Psi_r)] \rangle$$

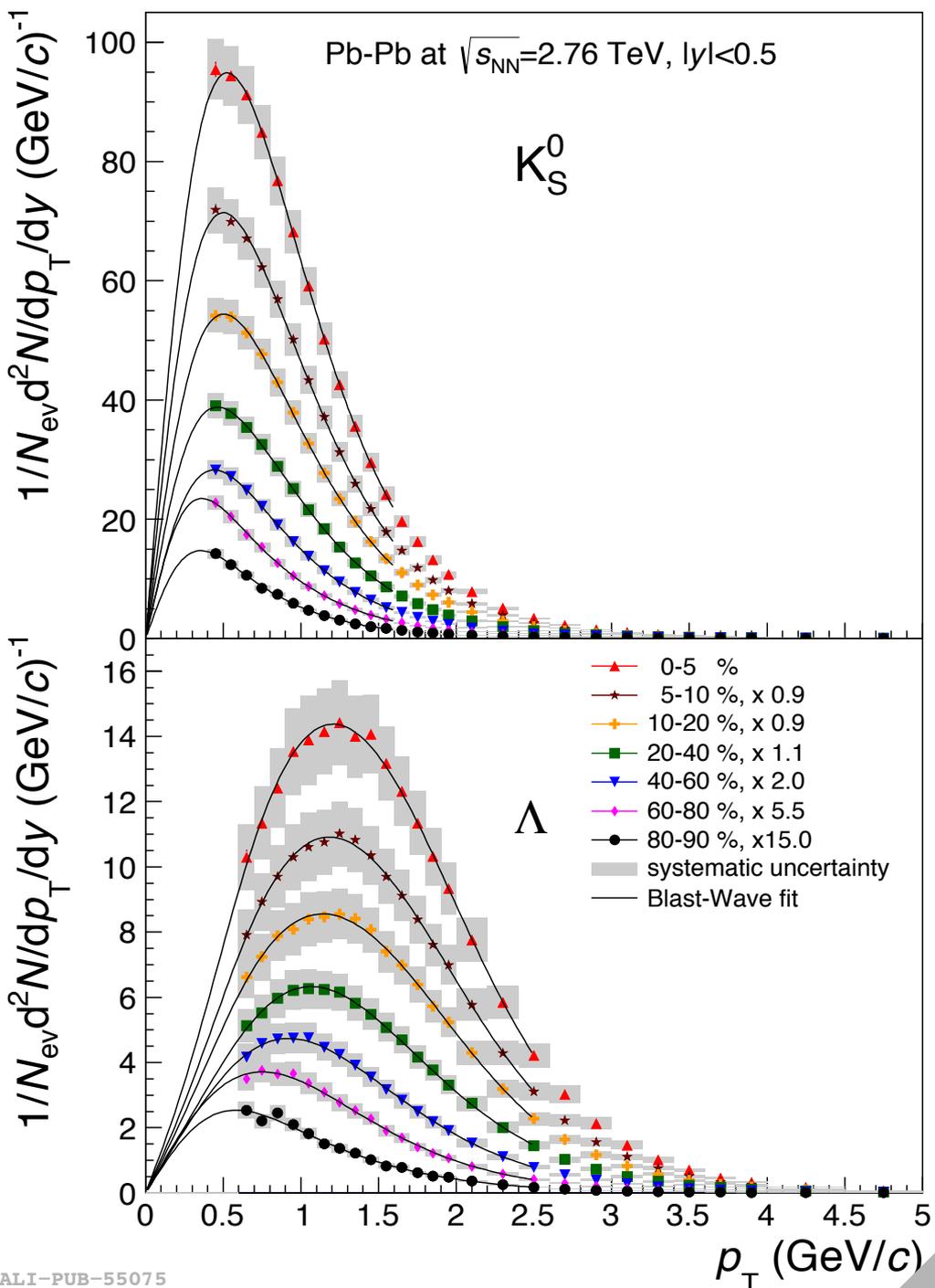
# Blastwave

- Doppler-shifted Boltzmann distribution, characterised by:
  - Collective Velocity  $\beta$
  - Temperature  $T$
- Allows extrapolation down to  $p_T=0$

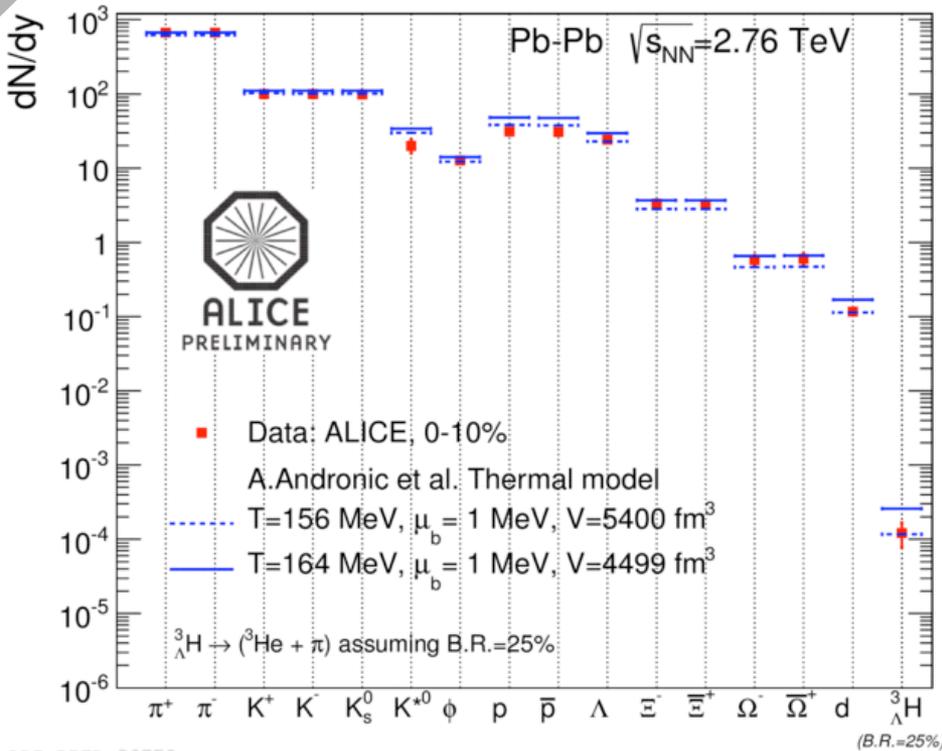
$$\frac{dN}{dp_T} \propto p_T \int_0^1 x m_T K_1\left(\frac{m_T \cosh(\eta)}{T}\right) I_0\left(\frac{p_T \sinh(\eta)}{T}\right) dx$$

$$\eta = \tanh^{-1}(\beta),$$

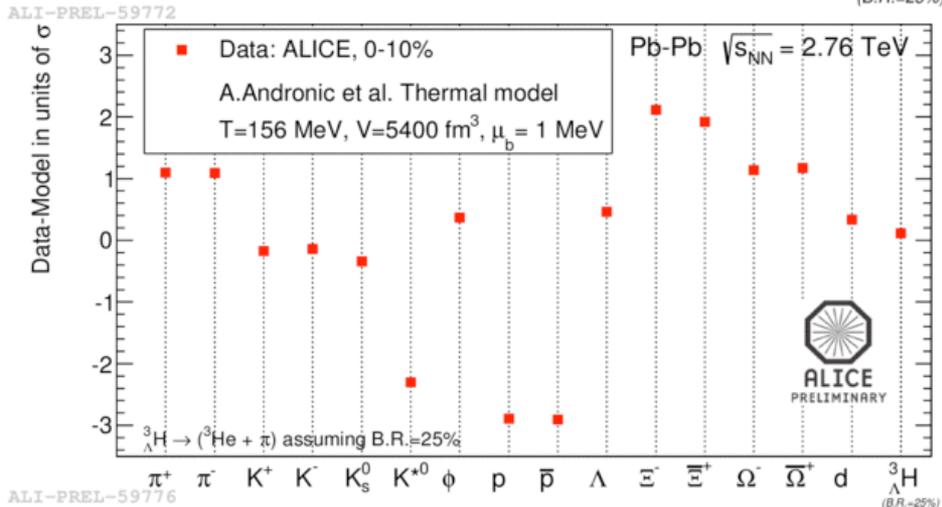
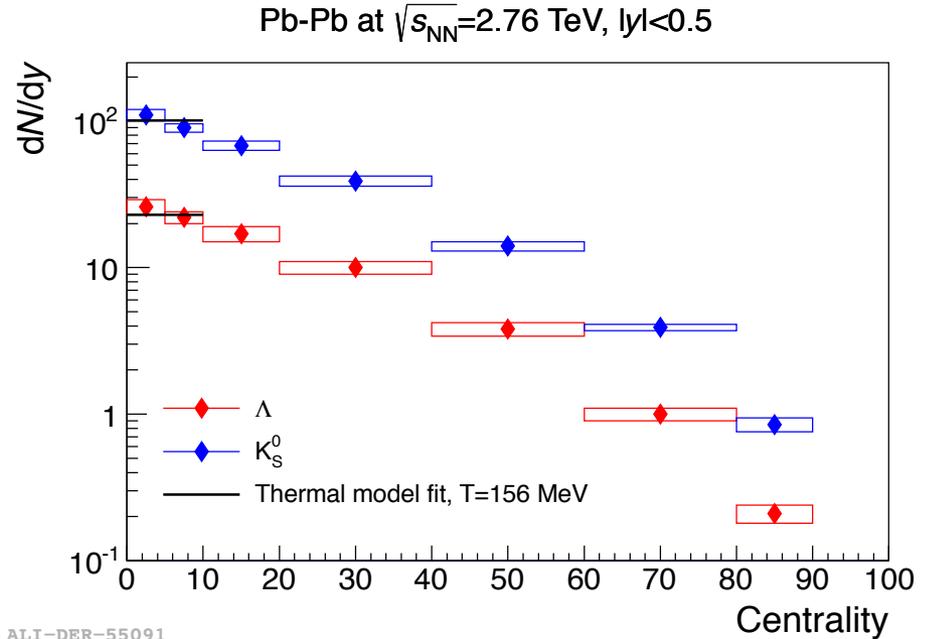
$$m_T = \sqrt{p_T^2 + m^2}$$



# Thermal Yields



Plot created from values in:  
ALICE Collaboration, (2013), arXiv:nucl-ex/1307.5530



ALI-DER-55091

Simultaneous thermal fit to all particles agrees well with  $\Lambda$  &  $K_S^0$  yields in central collisions

# Theory Models used

## **Hydro** – “Vish2+1”

- Viscous hydrodynamic model
- Sharp transition from fluid to hadrons at freeze-out temperature

*H. Song and U. W. Heinz, Phys. Lett. B658, 279 (2008), arXiv:nucl-th/0709.0742 .*  
*H. Song and U. W. Heinz, Phys. Rev. C77, 064901 (2008), arXiv:nucl-th/0712.3715 .*  
*H. Song and U. W. Heinz, Phys. Rev. C78, 024902 (2008), arXiv:nucl-th/0805.1756 .*

## **EPOS 2.17v3**

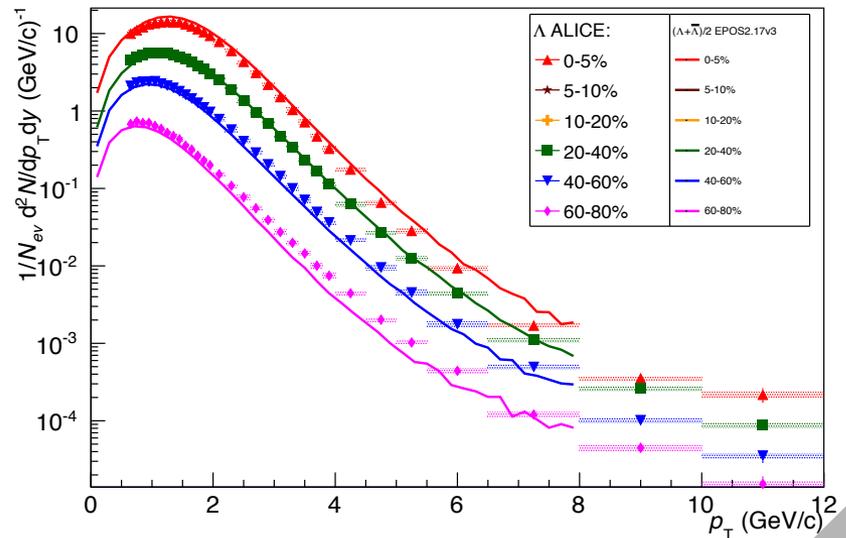
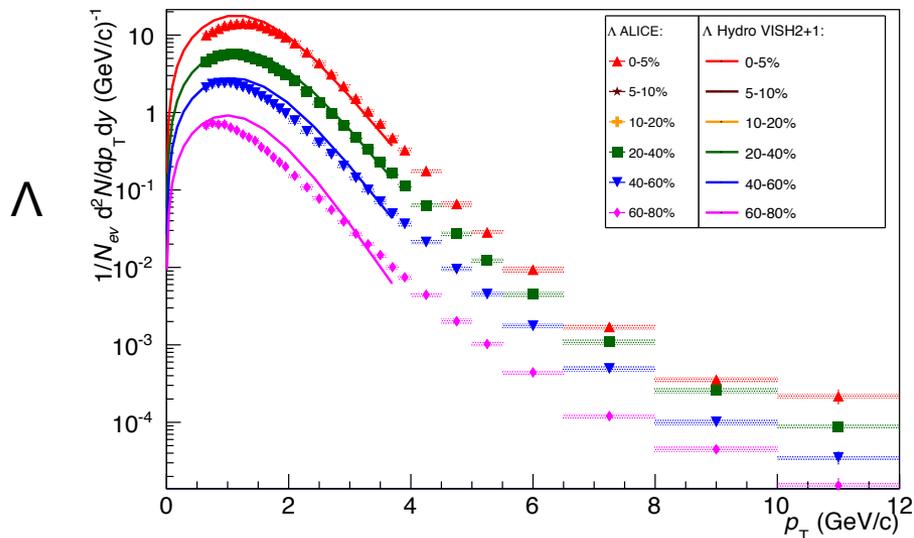
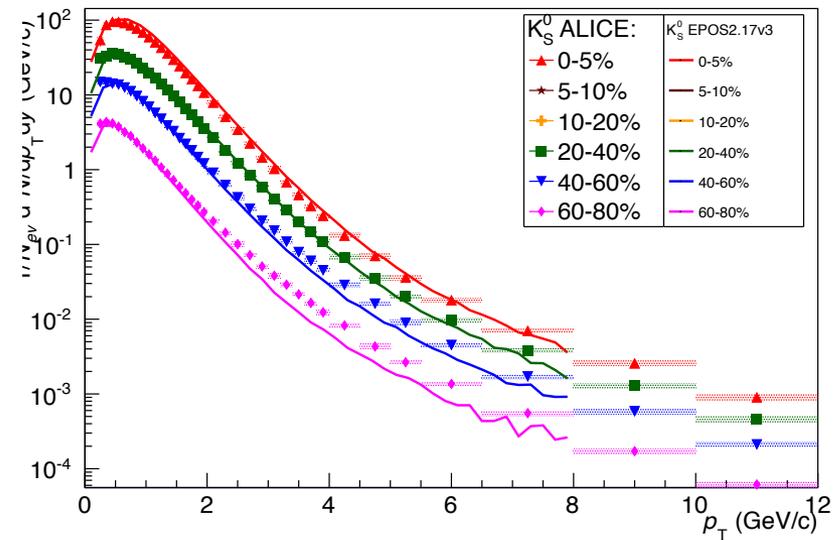
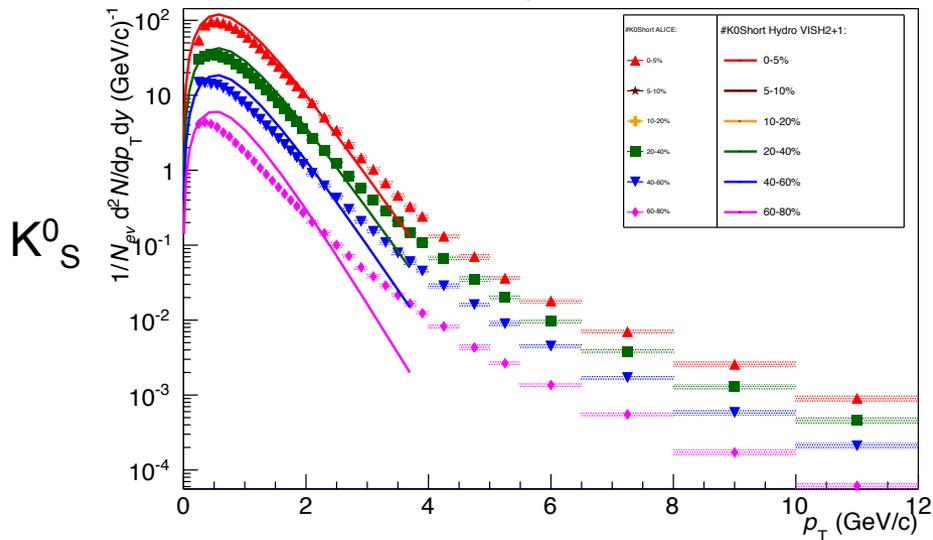
- String-breaking model
- Low energy strings form a ‘medium’
- Strings breaking in medium take quarks from medium
- Models ‘Core-Corona’ effect

*K. Werner, Phys. Rev. C 85, 064907 (2012) arXiv:1203.5704 .*  
*K. Werner, Phys. Rev. Lett. 109, 102301 (2012) arXiv:1204.1394*

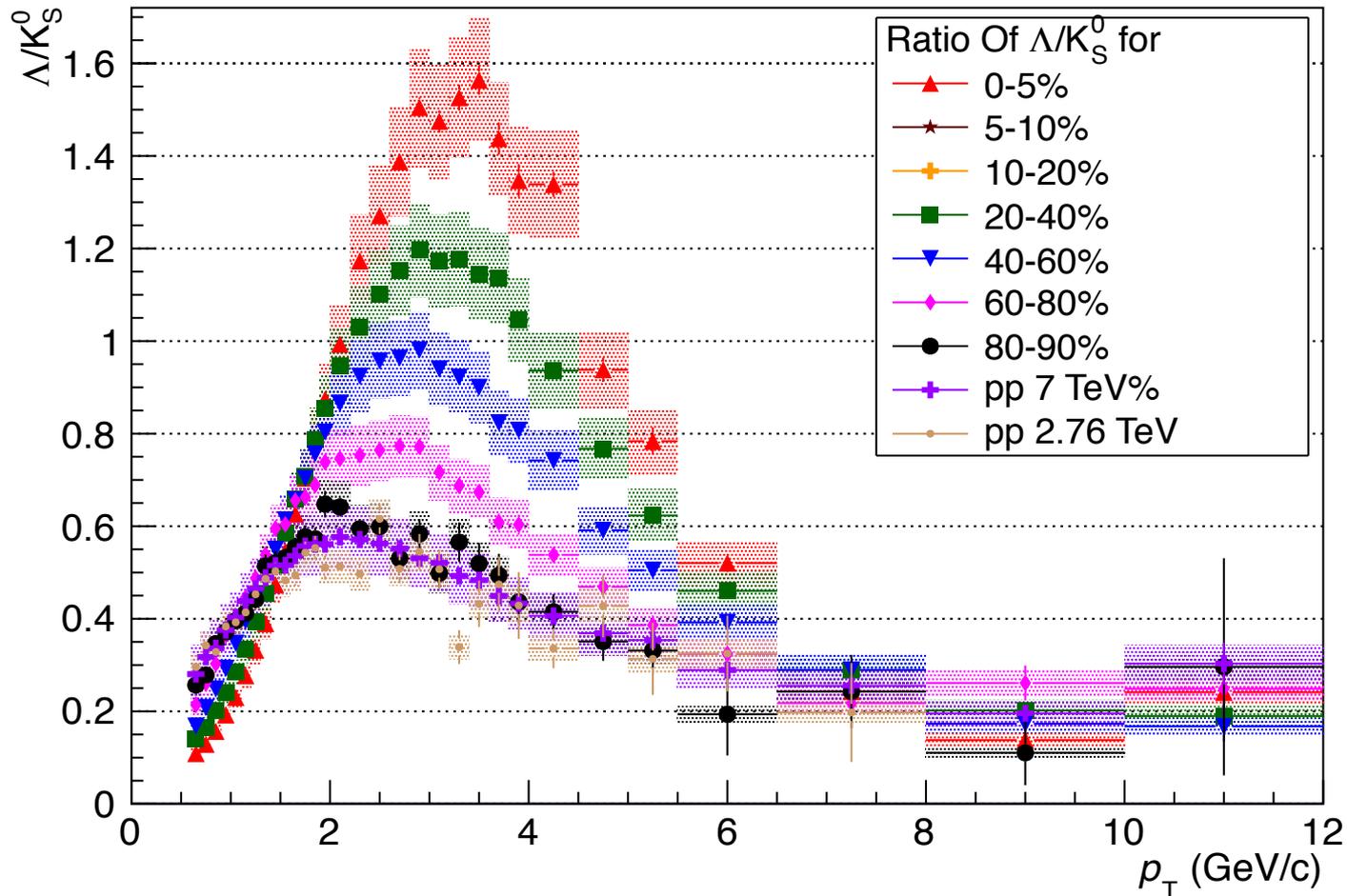
# Spectra compared to models

Hydro

EPOS

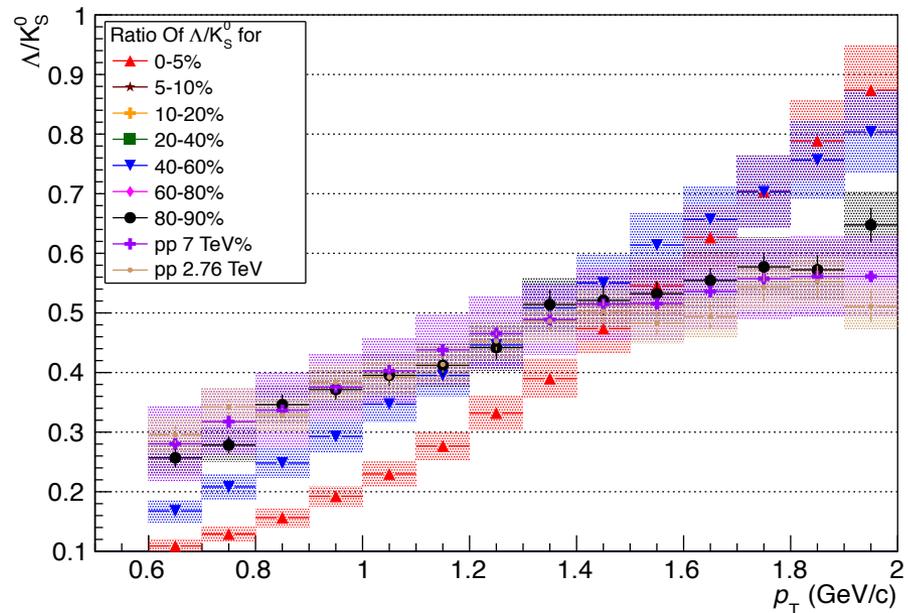
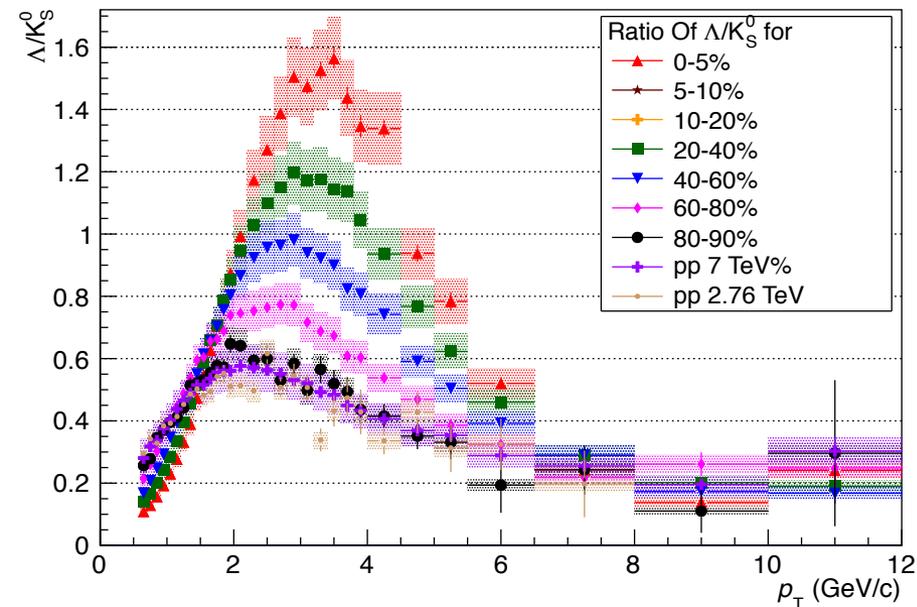


# $\Lambda/K_S^0$ Ratio

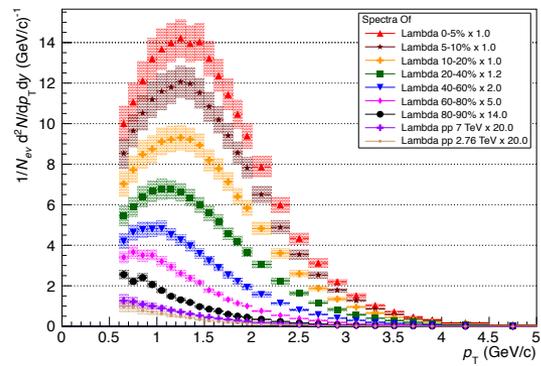


- Peripheral collisions, and p-p collisions at both energies agree well
- All centralities and systems consistent above  $p_T \sim 6.5$  GeV/c
- Magnitude of enhancement increases with centrality

# $\Lambda/K^0_S$ Ratio – low $p_T$

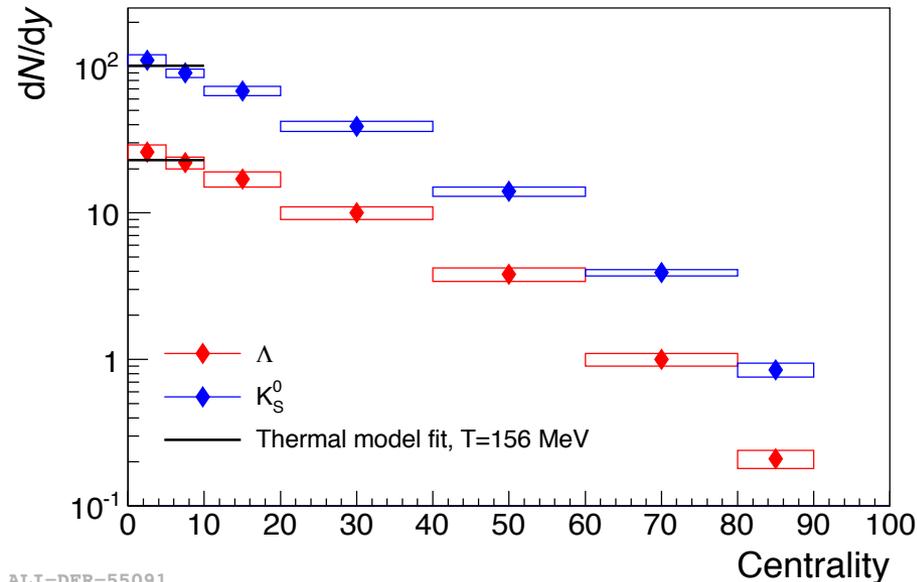


- At lowest momentum, more central collisions appear to show a suppression of the ratio
- This could hint towards a redistribution of particles from low to mid momentum



# $\Lambda/K_S^0$ ratio – integrated ratio

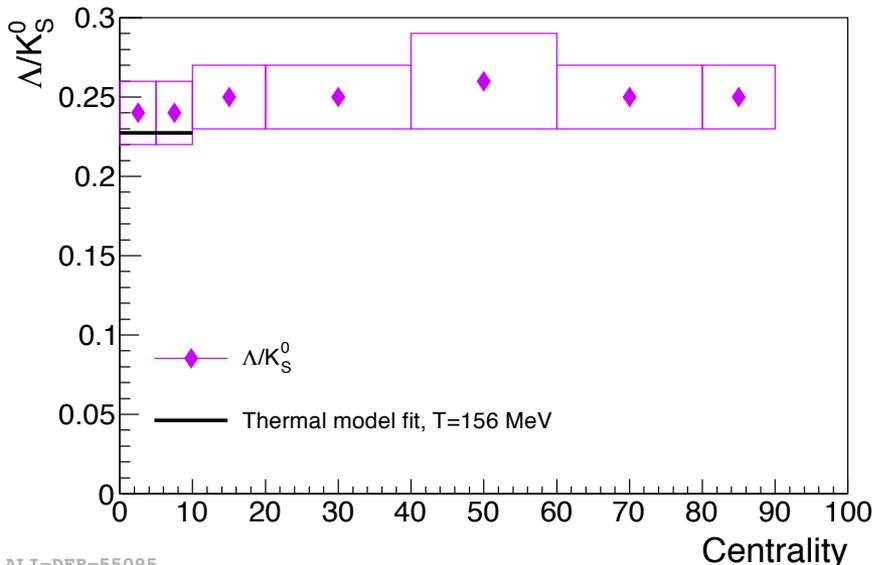
Pb-Pb at  $\sqrt{s_{NN}}=2.76$  TeV,  $|y|<0.5$



- Using a blastwave fit to extrapolate to low  $p_T$ , the yield of  $\Lambda$  &  $K_S^0$  can be integrated over  $p_T$
- The  $\Lambda/K_S^0$  ratio of these integrated yields appears constant with centrality
- Further supports that baryons / mesons are redistributed in  $p_T$  rather than enhanced / suppressed

ALI-DER-55091

Pb-Pb at  $\sqrt{s_{NN}}=2.76$  TeV,  $|y|<0.5$

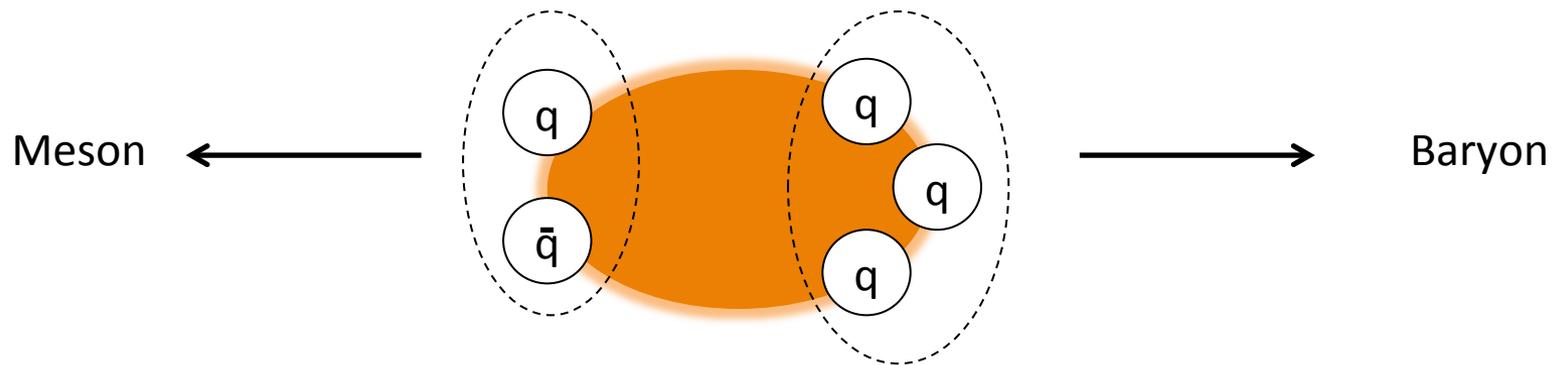


Plot created from values in:  
ALICE Collaboration,  
(2013), arXiv:nucl-ex/  
1307.5530

ALI-DER-55095

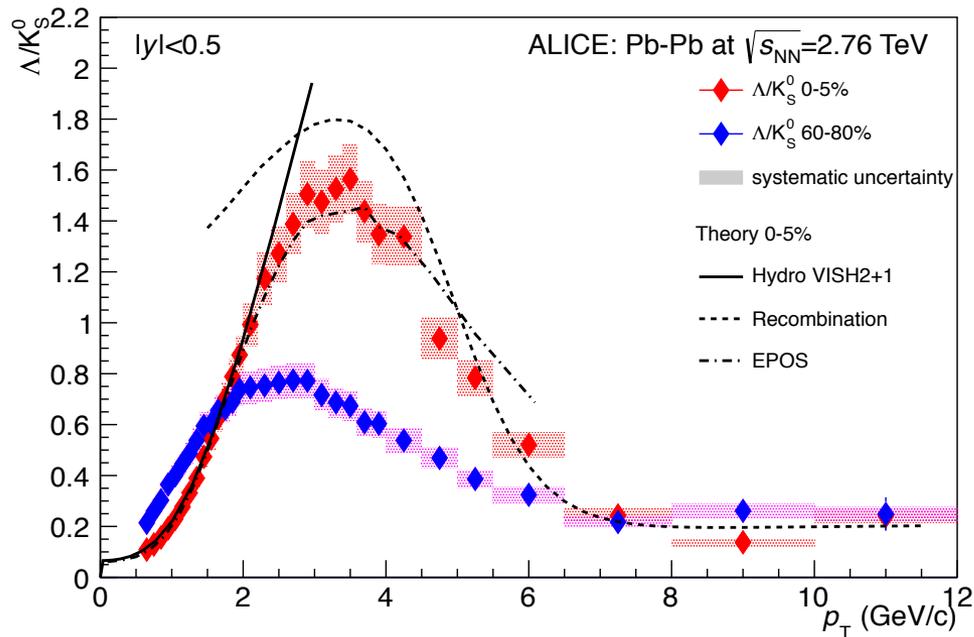
# Coalescence

- As a QGP cools, it must pass back over the phase boundary and undergo hadronization
- The distribution of hadrons formed in this way need not be the same as for vacuum fragmentation
- In particular, it may allow for the formation of baryons/mesons by grouping 3/2 quarks close in phase space

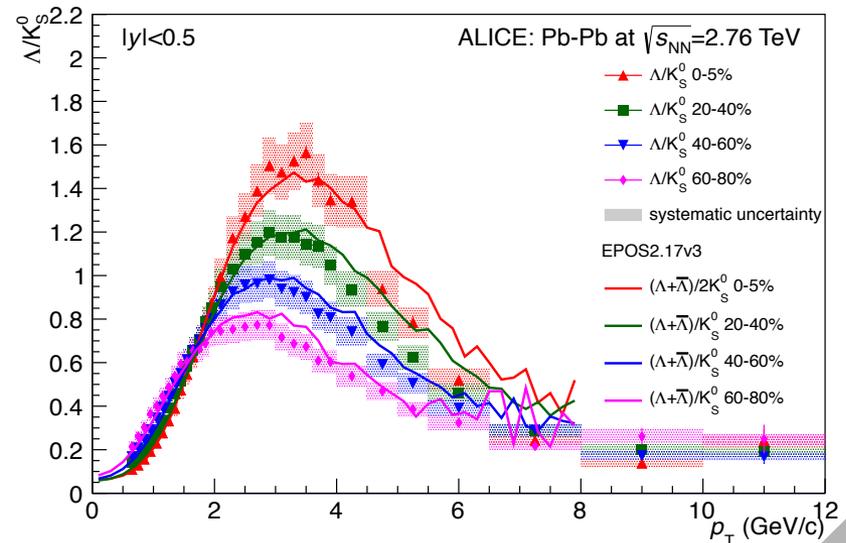
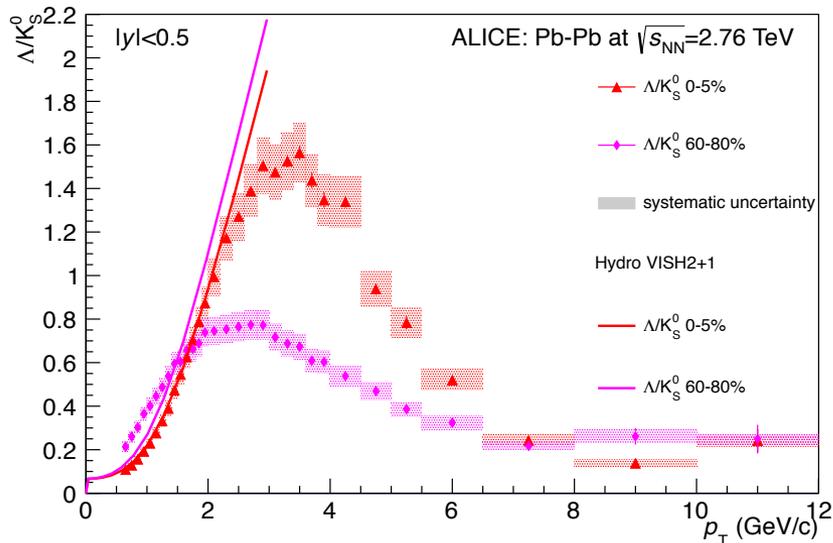


- From this process, one would expect the average momentum of the baryon to be  $\sim 3\langle p_q \rangle$ , while the meson would have  $\sim 2\langle p_q \rangle$
- So, for hadrons formed from the medium, one could expect the baryons to be formed at higher momenta than the mesons

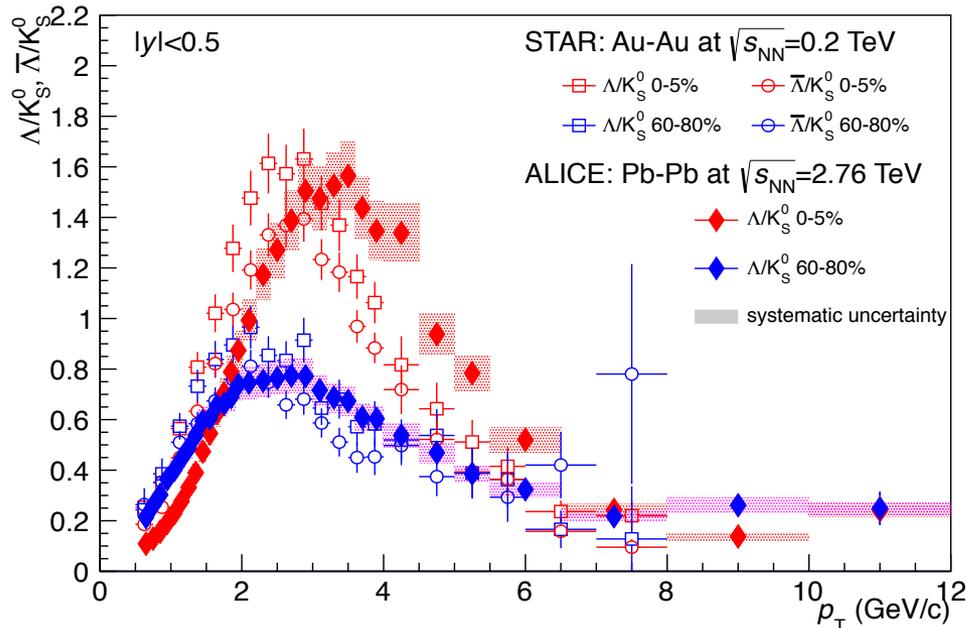
# $\Lambda/K_S^0$ Ratio - theory



- Hydrodynamics describes low- $p_T$ , central ratio very well
- Recombination describes general trend, but not normalisation
- Hydrodynamics less accurate for more peripheral collisions
- EPOS describes ratio well

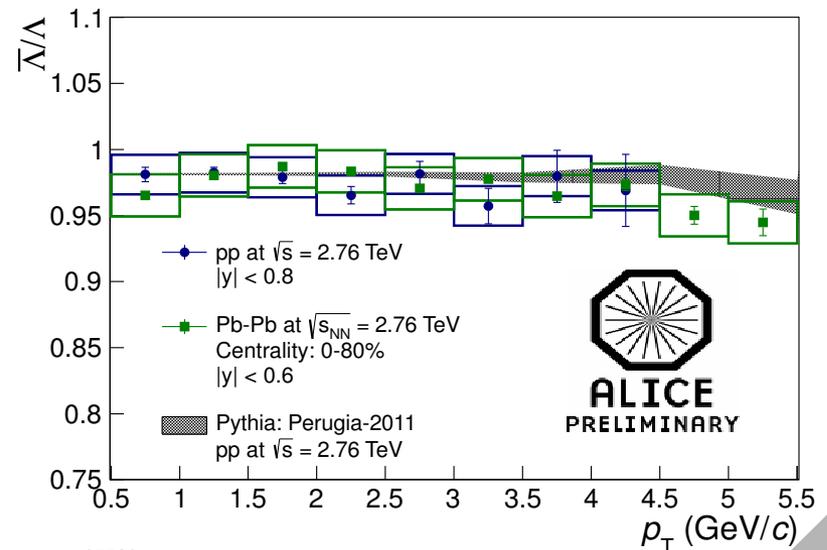
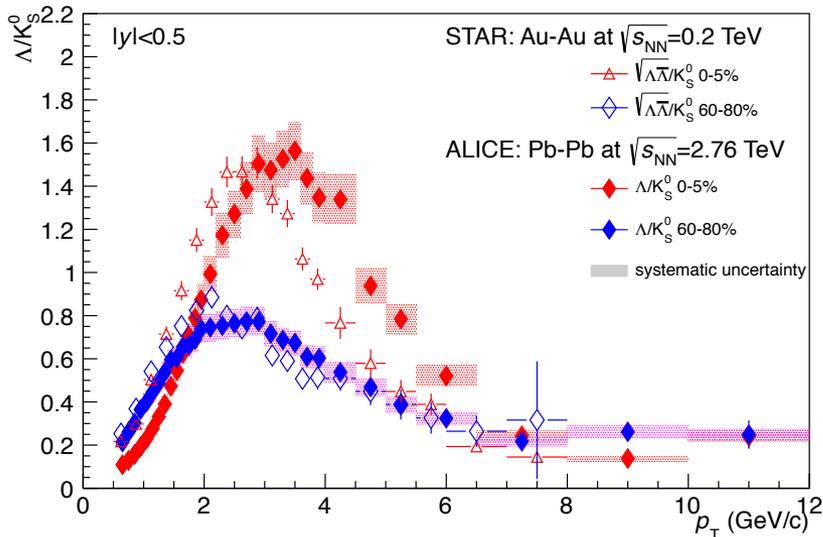


# $\Lambda/K_S^0$ Ratio - STAR

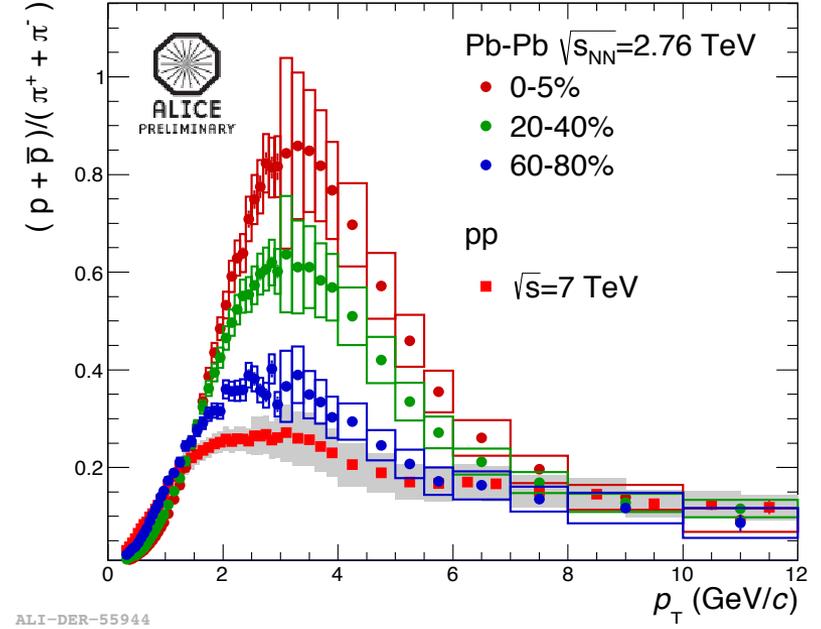
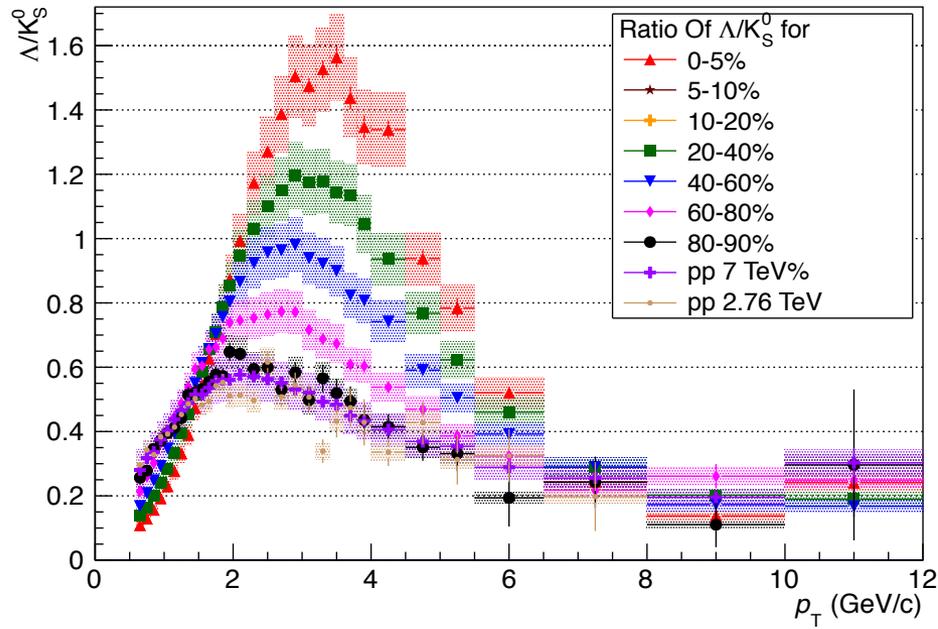


- Anti- $\Lambda/\Lambda$  ratio:
  - at STAR  $\sim 0.8$
  - at ALICE  $\sim 1$
- Magnitude of enhancement unchanged
- Persists to slightly higher  $p_T$  at LHC energy

STAR points from G. Agakishiev et al. (STAR Collaboration), *Phys. Rev. Lett.* 108, 072301 (2012), [arXiv:nucl-ex/1107.2955](https://arxiv.org/abs/1107.2955)

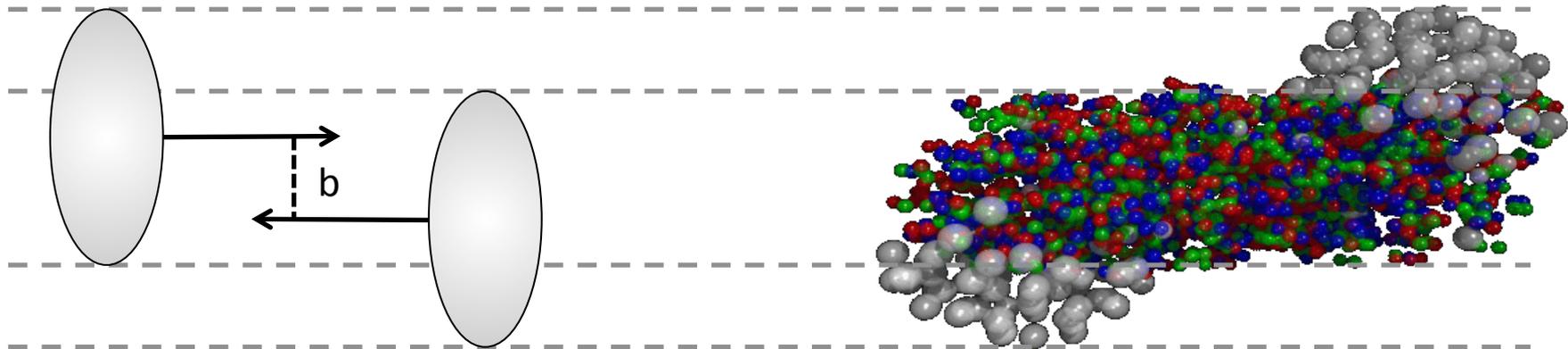


# $\Lambda/K_S^0$ Ratio – $p/\pi$



- Behaviour of the  $p/\pi$  ratio very similar to  $\Lambda/K_S^0$
- Enhancement of  $\sim x3$  in range  $1.5 < p_T < 6.5$  GeV/c
- $p/\pi$  ratio consistently  $\sim$ half of  $\Lambda/K_S^0$

# The Glauber Model

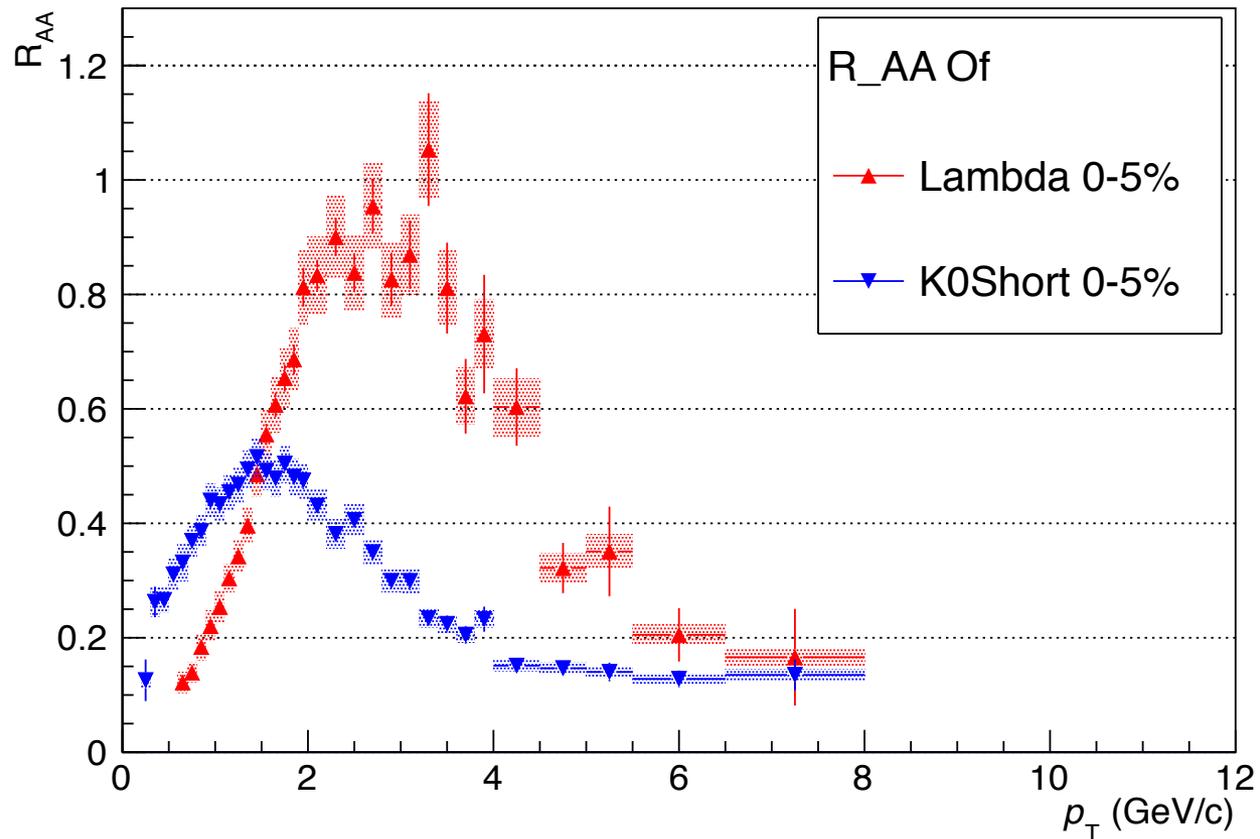


- The Glauber model is used to relate heavy ion collisions to p-p collisions
- Nucleons are considered to be distributed within the nucleus according to the Saxon-Woods nuclear density
$$\rho(r) = \rho_0 \frac{1 + \omega(r/R)^2}{1 + \exp\left(\frac{r-R}{a}\right)}$$
- It is assumed that all collisions are binary with a constant cross section, and that there is no deflection of nucleons
- This allows the calculation of the equivalent number of p-p collisions as the impact parameter,  $b$ , varies

# $R_{AA}$ of $\Lambda/K^0_S$

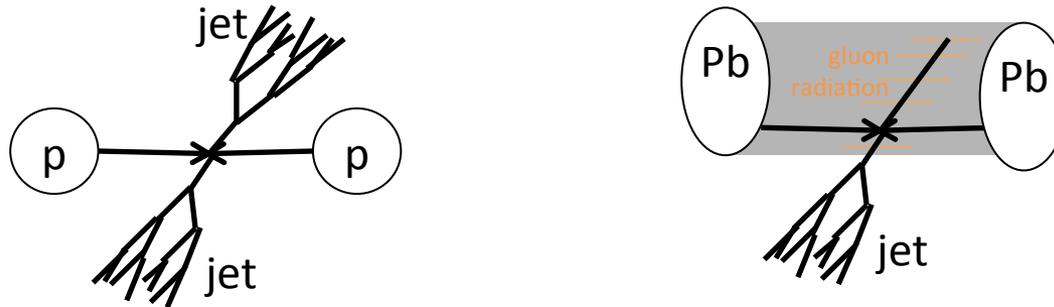
$$R_{AA}(b) = \frac{\frac{d^2 N^{AA}}{dp_T dy}}{N_{coll}^{AA}(b) \frac{d^2 N^{pp}}{dp_T dy}}$$

Centrality	$\langle N_{coll} \rangle$
0-5%	1685
10-20%	921.2
40-60%	127.7
80-100%	4.441



# Jet Quenching

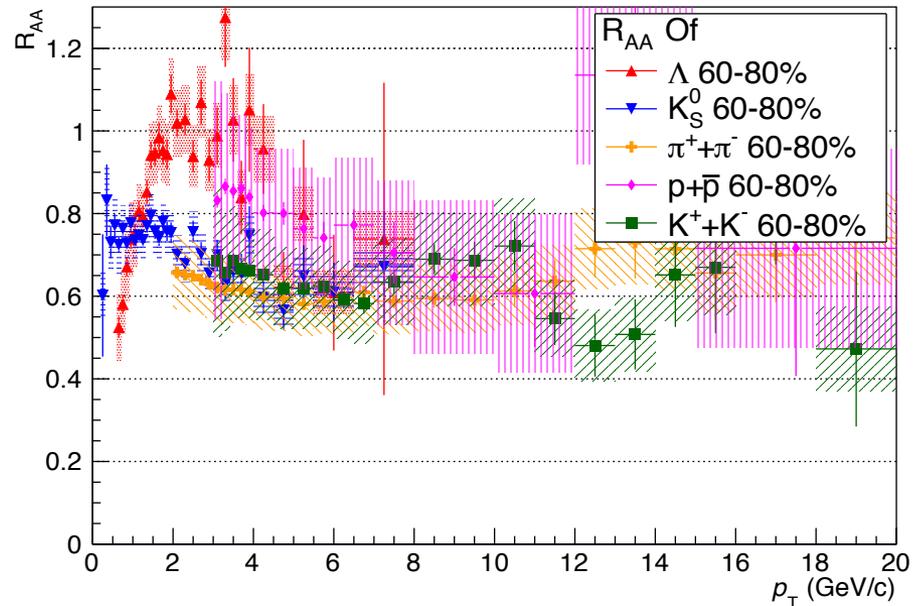
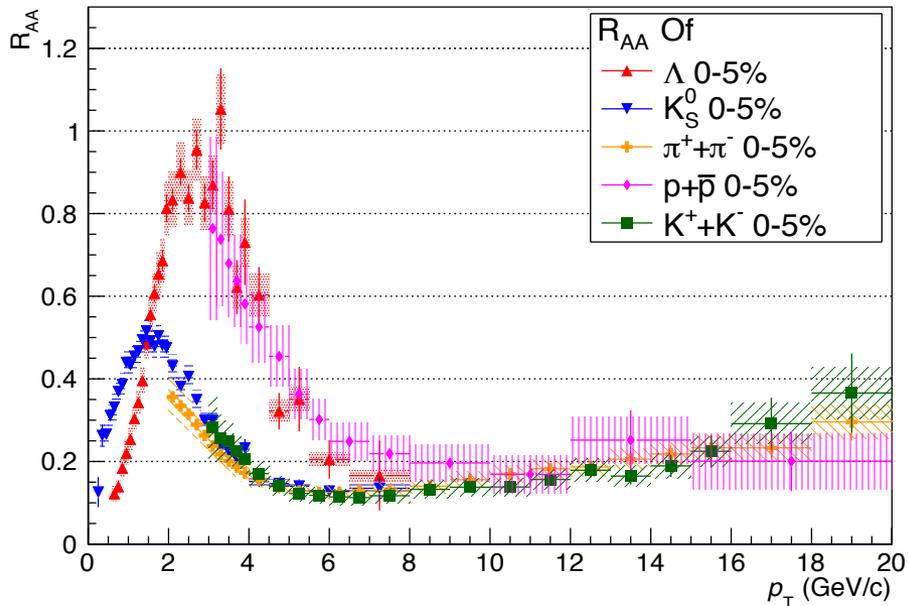
- Jets are formed as a high energy parton from the initial collision fragments



- If a QGP is formed, the parton would be expected to lose energy to strong interactions within the plasma
- This leads to jets being emitted at lower energy, or not escaping the medium at all
- Back-to-back jets may be asymmetric in energy, as one travels further through the medium than the other
- Some models suggest a variation in the suppression dependent on jet hadrochemistry

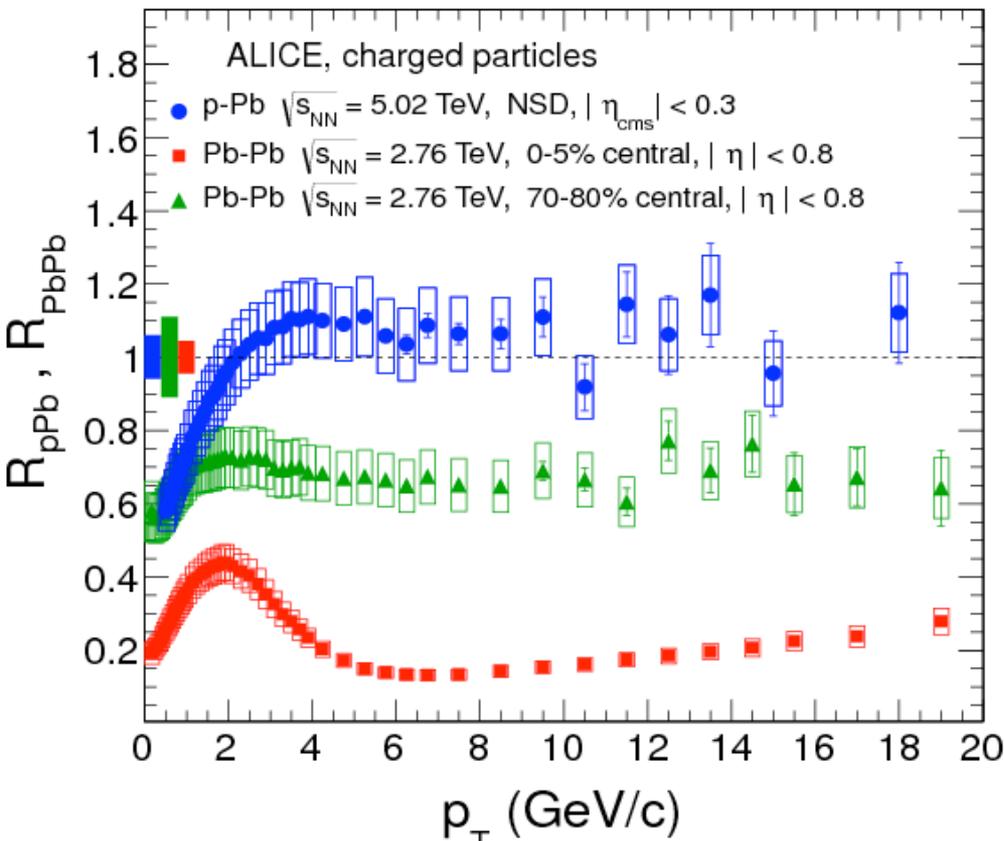
*S. Sapeta and U. A. Wiedemann, Eur.Phys.J. C55, 293 (2008), arXiv:0707.3494 [hep-ph] .*  
*P. Aurenche and B. Zakharov, Eur.Phys.J. C71, 1829 (2011), arXiv:1109.6819 [hep-ph] .*  
*R. Bellwied and C. Markert, Phys.Lett. B691, 208 (2010), arXiv:1005.5416 [nucl-th] .*

# $R_{AA}$ with other particles

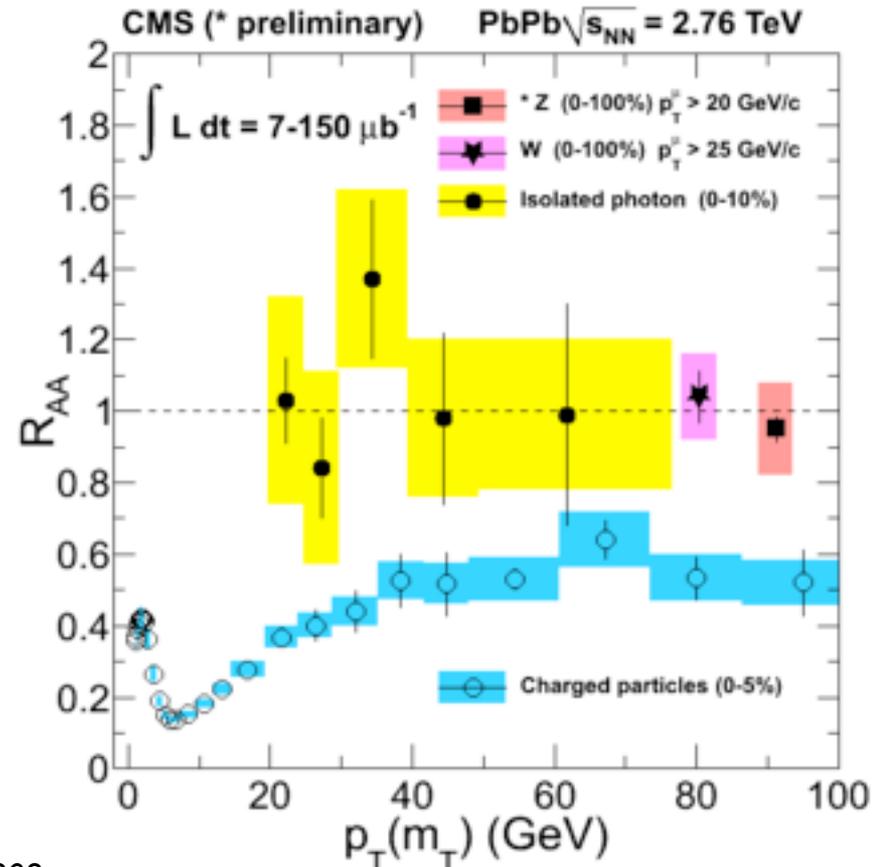


- Separation of baryons and mesons at lower  $p_T$ , as expected from  $\Lambda/K_S^0$  ratio behaviour
- Above 6 GeV/c, all hadrons are consistent
- No sign of hadrochemistry dependent jet suppression

# $R_{AA}$ for p-Pb & Electroweak probes



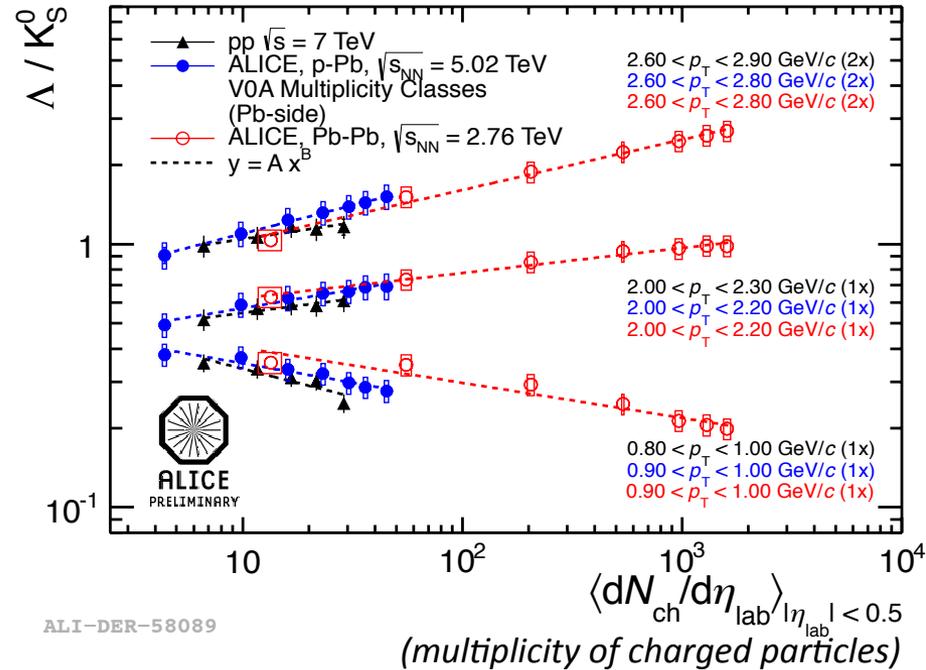
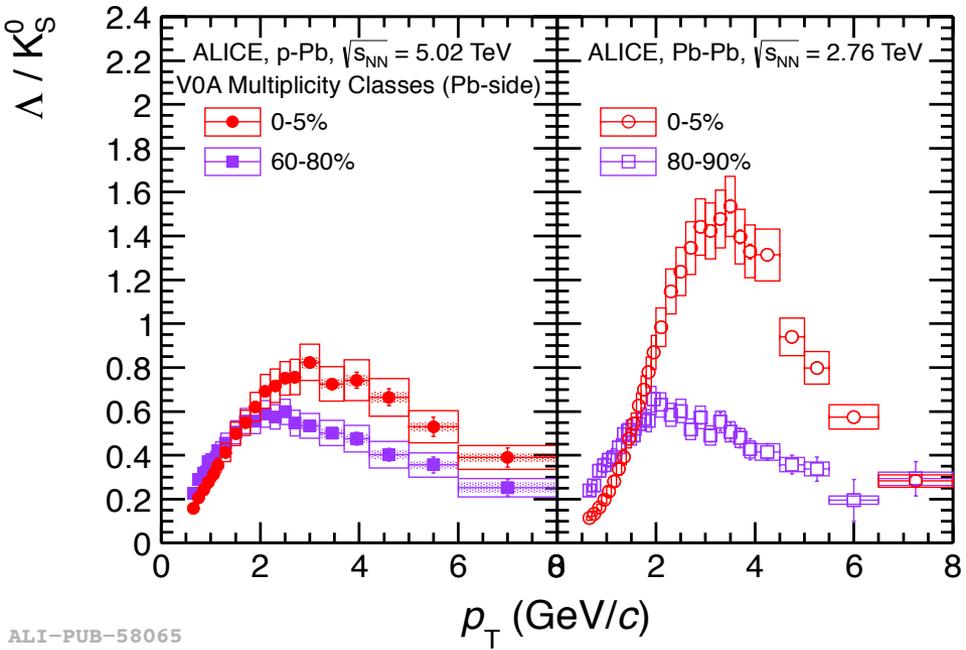
B. Abelev et al. (ALICE Collaboration), *Phys.Rev.Lett.* 110 (2013) 082302



S. Chatrchyan. et al., (CMS Collaboration), *Eur. Phys. J. C* 72, 1945 (2012).

- $R_{AA}$  is not consistent with 1 even for more peripheral events or  $p_T \sim 100$  GeV/c
- $R_{pPb}$  is consistent with 1 above  $p_T \sim 2$  GeV/c
- $R_{AA}$  of electroweak probes is consistent with 1

# $\Lambda/K^0_S$ Ratio – p-Pb



- In p-Pb collisions, an enhancement of the  $\Lambda/K^0_S$  ratio with multiplicity can be seen
- The enhancement is less pronounced when compared with Pb-Pb collisions
- The value of the  $\Lambda/K^0_S$  ratio appears to be a function of the multiplicity.
- Further work is needed to understand the origin of this behaviour

# Conclusions

- $\Lambda$  &  $K_S^0$  spectra measured for transverse momentum range 0.4-12 GeV/c in Pb-Pb and pp collisions
- $\Lambda/K_S^0$  ratio shows significant enhancement in central PbPb collisions when compared to pp collisions
- EPOS reproduces behaviour well, while pure Hydrodynamic and Recombination models are less successful
- Studies of the  $R_{AA}$  of identified particles shows no hadrochemical dependence at high  $p_T$ .
- Behaviour of the  $\Lambda/K_S^0$  ratio in p-Pb collisions is similar to that in PbPb



# Backup

# QGP

- Quark-gluon plasma is a **thermalised** system of **deconfined** quarks and gluons
- As a first approximation, we can treat Hadron Gas / QGP as an ideal gas:

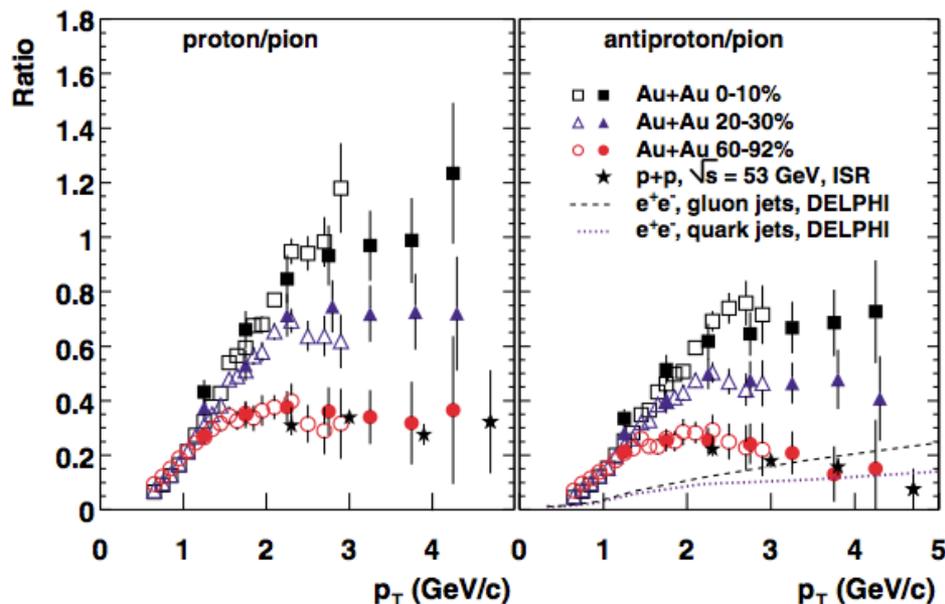
$$p = \frac{\varepsilon}{3} = g \frac{\pi^2}{90} T^4$$

$$\frac{\varepsilon}{T^4} = g \frac{\pi^2}{30}$$

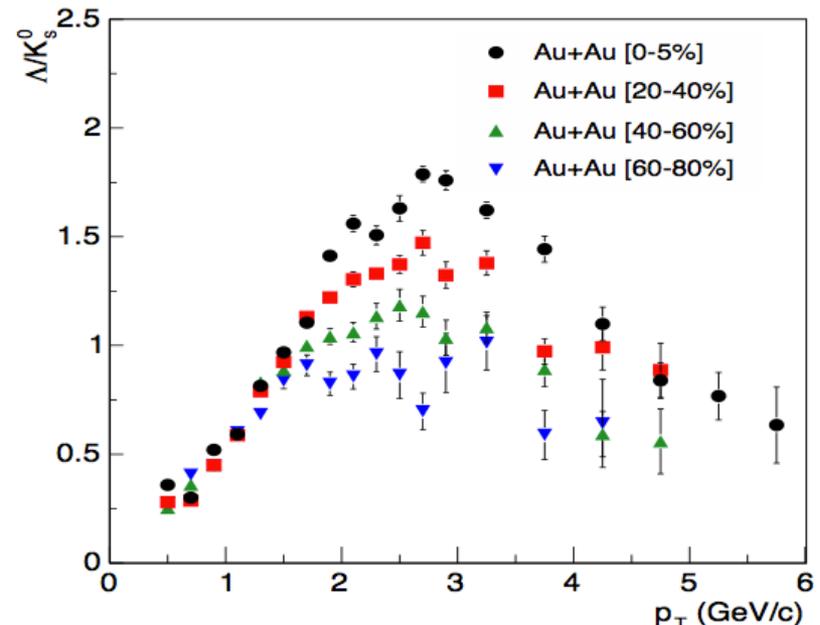
- For Hadron Gas, pions dominate so  $g \sim 3$
- For QGP  $g = g_{\text{gluons}} + g_{\text{quarks}} = 2_{\text{spin}} * 8_{\text{color}} + 7/8 * 3_{\text{flavor}} * 2_{\text{q/anti-q}} * 2_{\text{spin}} * 3_{\text{color}} \sim 48$
- By modelling hadrons as a 'bag', one can estimate the pressure holding them together
- A deconfined state can be attained if the QGP pressure is greater than the Bag pressure
  - Transition temperature  $\sim 140\text{MeV}$

# Motivation

- Observed at SPS & RHIC that  $p/\pi$  &  $\Lambda/K^0_S$  ratios are enhanced at intermediate momentum in heavy ion collisions, when compared to pp.
- Examining how this effect evolves with increased energy gives insight into the interplay between fragmentation and potential baryon-enhancing effects.
- $\Lambda$  &  $K^0_S$  can be identified with a single technique over a wide momentum range



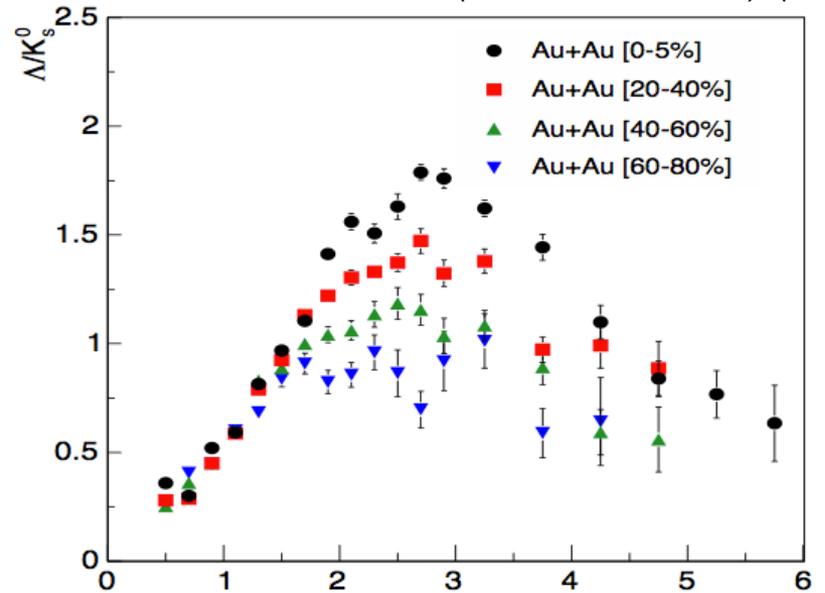
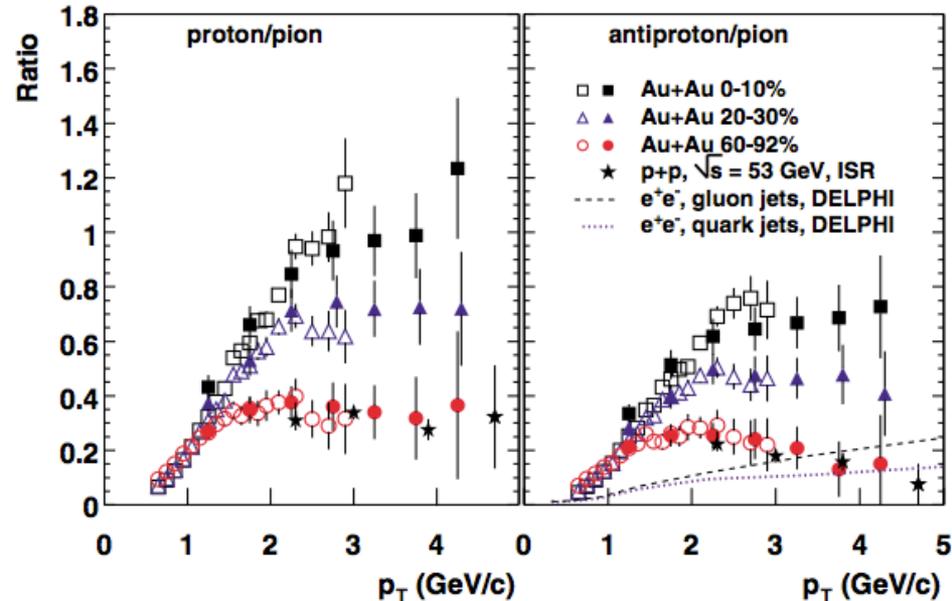
S. S. Adler et al. (PHENIX Collaboration), *Phys. Rev. Lett.* 91, 172301 (2003)



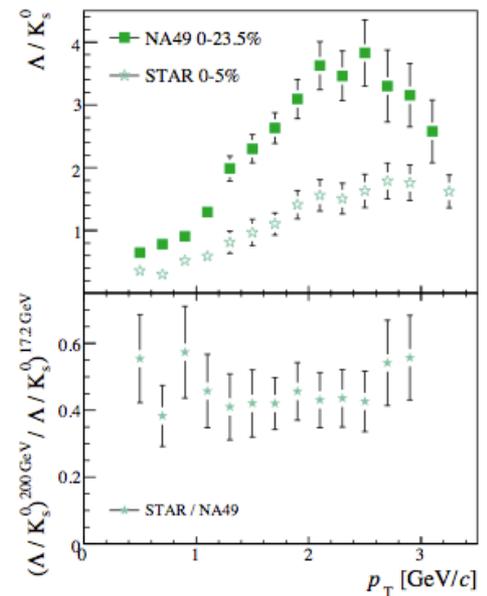
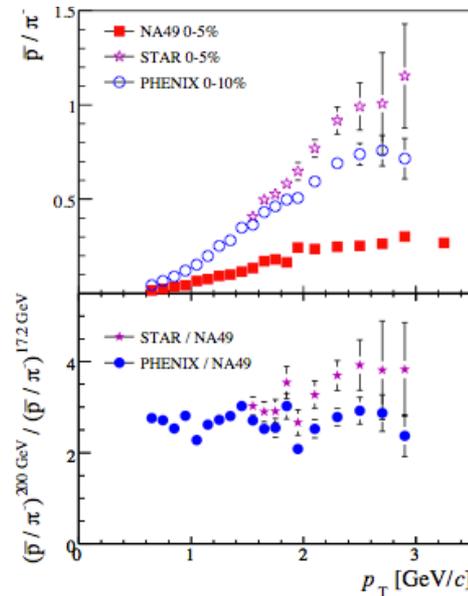
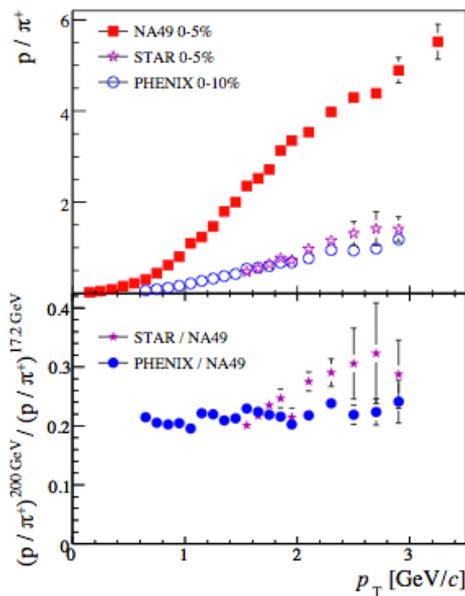
J. Adams et al. (STAR Collaboration), (2006),

# Motivation

*J. Adams et al. (STAR Collaboration), (2006),*



*S. S. Adler et al. (PHENIX Collaboration), Phys. Rev. Lett. 91, 172301 (2003)*



*T. Schuster and A. Laszlo (for the NA49 Collaboration). J.Phys. G32 (2006) S479-S482.*

# Theory Models used

**VISH2+1** - a viscous hydro code in (2+1) dimensions assuming longitudinal boost invariance. Sharp transition from fluid to non-interacting particles at freeze-out temperature.

*H. Song and U. W. Heinz, Phys. Lett. B658, 279 (2008), arXiv:nucl-th/0709.0742 .*  
*H. Song and U. W. Heinz, Phys. Rev. C77, 064901 (2008), arXiv:nucl-th/0712.3715 .*  
*H. Song and U. W. Heinz, Phys. Rev. C78, 024902 (2008), arXiv:nucl-th/0805.1756 .*

**EPOS 2.17v3** – String breaking model of particle creation, where low  $p_T$  particles are allowed to reinteract, simulating hydrodynamic behaviour. Strings breaking in-medium take their quarks from the fluid, rather than Schwinger Mechanism. Also uses Core-Corona model to determine fluid region.

*K. Werner, Phys. Rev. C 85, 064907 (2012) arXiv:1203.5704 .*  
*K. Werner, Phys. Rev. Lett. 109, 102301 (2012) arXiv:1204.1394 .*

**Recombination** – a model allowing quarks from the medium to coalesce into hadrons. Radial flow and jet quenching have been set to match that observed at the LHC.

*B. Muller, J. Schukraft, B. Wyslouch, arXiv:1202.3233 .*

# Reconstruction - Feeddown

Decay	Branching Ratio
$\Xi^- \Rightarrow \Lambda \pi^-$	$99.887 \pm 0.035\%$
$\Xi^0 \Rightarrow \Lambda \pi^0$	$99.525 \pm 0.012\%$

- $\Lambda$  coming from weak decays of  $\Xi$  are removed
- $\Omega$  decay also considered, but found to be negligible

- With MC, we create a Feed-down matrix relating the  $p_T$  distribution of  $\Lambda$  to that of  $\Xi$
- This can then be scaled to the measured  $\Xi$   $p_T$  spectrum
- The raw yield of primary  $\Lambda$  is then obtained as:

$$\Lambda_{primary}^{raw} = \Lambda_{measured}^{raw} - \sum_j F_{ij} \int_{p_T(bin)} \frac{dN}{dp_T}(\Xi^-)$$

$$F_{ij} = \frac{N_{recon}(\Lambda)_{from \Xi bin j}^{in bin i}}{N_{generated}(\Xi)_{in bin j}}$$

- Feed-down varies from  $\sim 25\%$  at low  $p_T$  to negligible levels at high  $p_T$

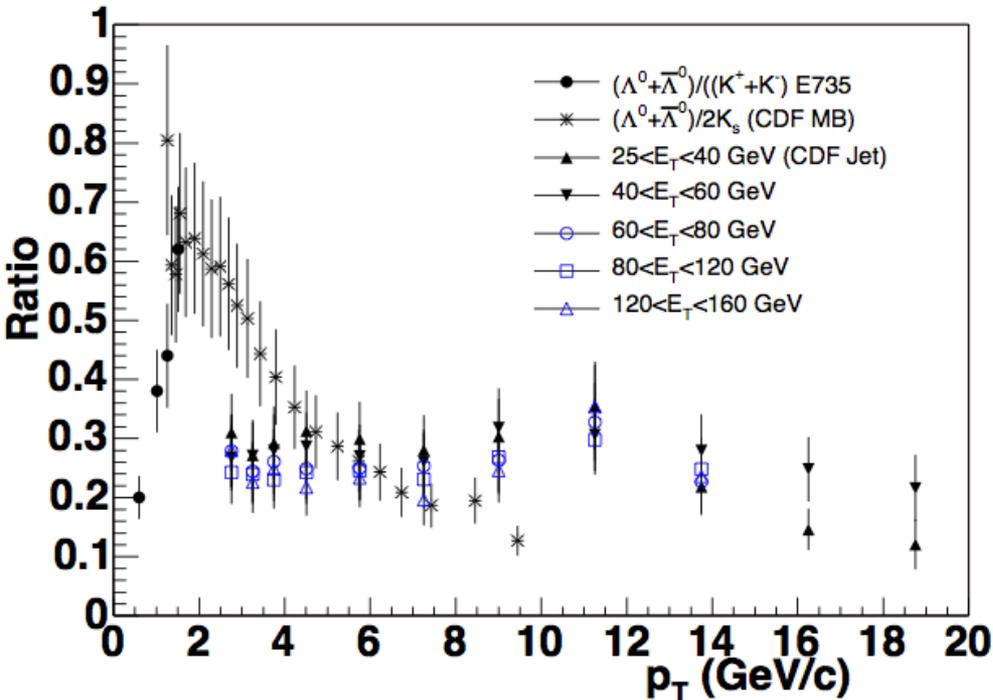
# Backup

Centrality	dN/dy		
	$\Lambda$	$K^0_S$	$\Lambda/K^0_S$
0-5%	26+3	110+10	0.24+0.02
5-10%	22+2	90+6	0.24+0.02
10-20%	17+2	68+5	0.25+0.02
20-40%	10+1	39+3	0.25+0.02
40-60%	3.8+0.4	14+1	0.26+0.03
60-80%	1.0+0.1	3.9+0.2	0.25+0.02
80-90%	0.21+0.03	0.85+0.09	0.25+0.02

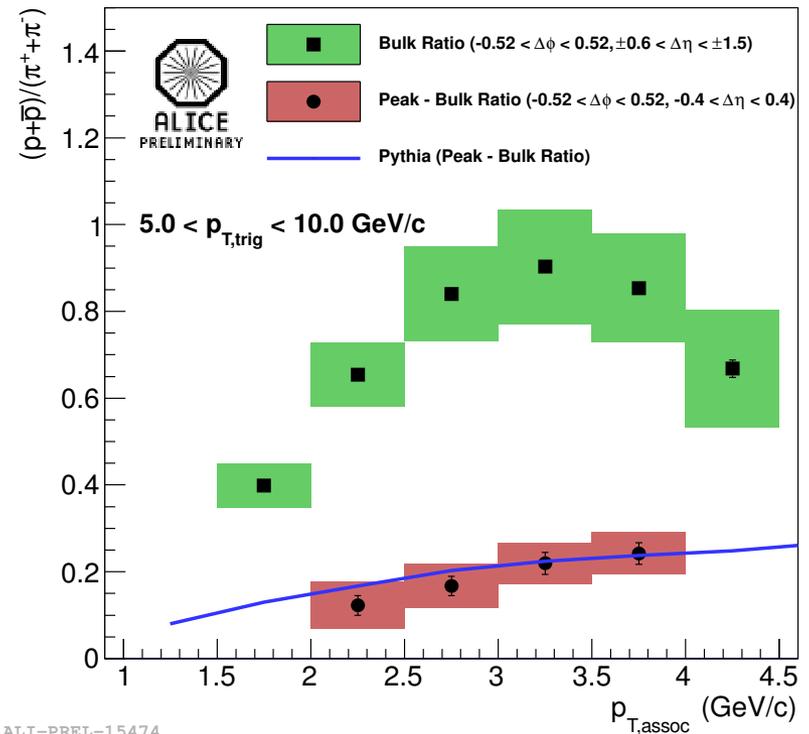
- Integrated ratio is constant within errors
- Suggests that baryons / mesons are redistributed in  $p_T$  rather than enhanced / suppressed

# L/K Ratio – Jets & Bulk

Pb-Pb,  $\sqrt{s_{NN}} = 2.76\text{TeV}$ , 0-10% central



T. Altonen et al. (CDF Collaboration), Phys. Rev. D 88, 092005 (2013)

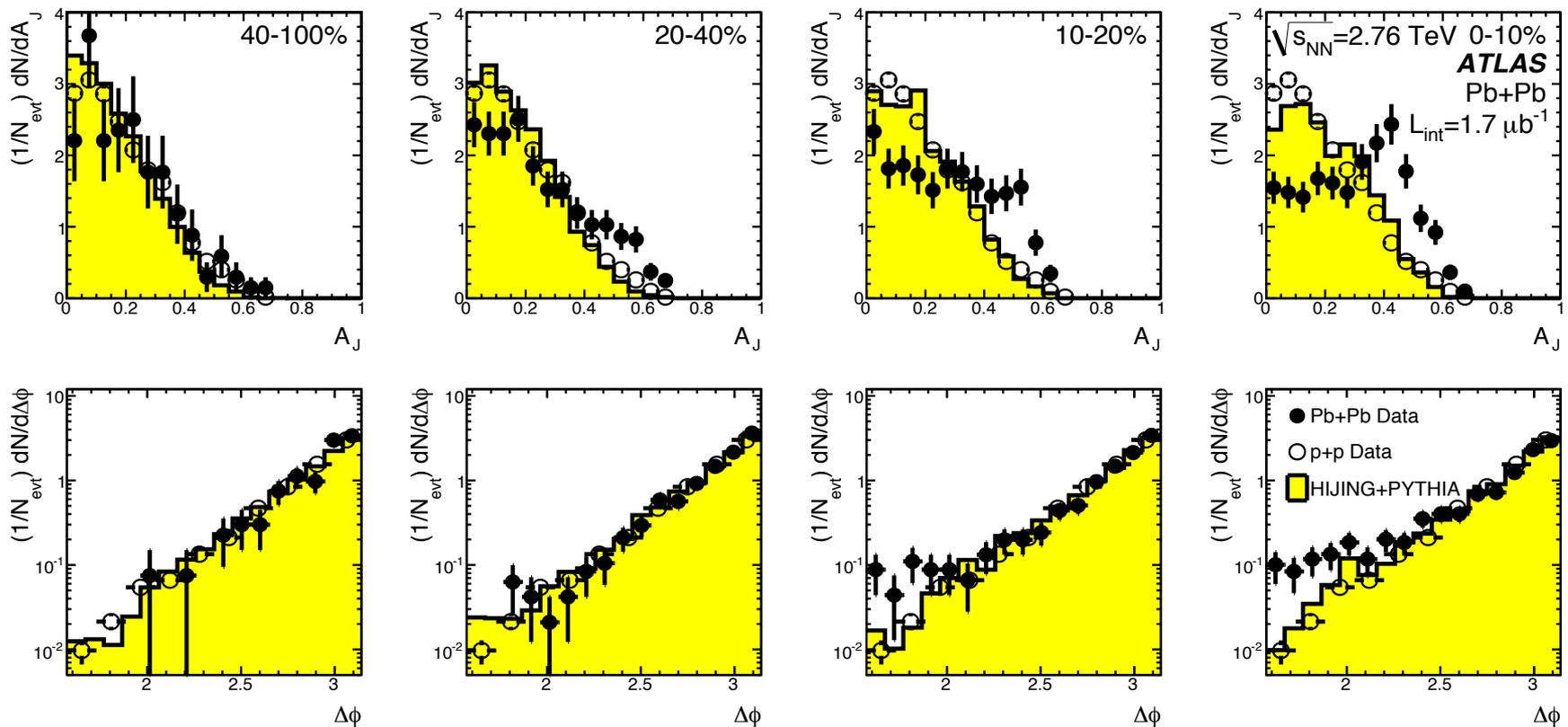


ALI-PREL-15474

Misha Veldhoen for the ALICE collaboration, arXiv:1207.7195

- In  $pp\bar{p}$  collisions at Fermilab the Ratio of  $\Lambda/K$  is suppressed in jets, when compared to inclusive studies
- This behaviour can be seen in Pb-Pb collisions at the LHC, where the  $p/\pi$  ratio in jets is suppressed relative to that in the underlying event

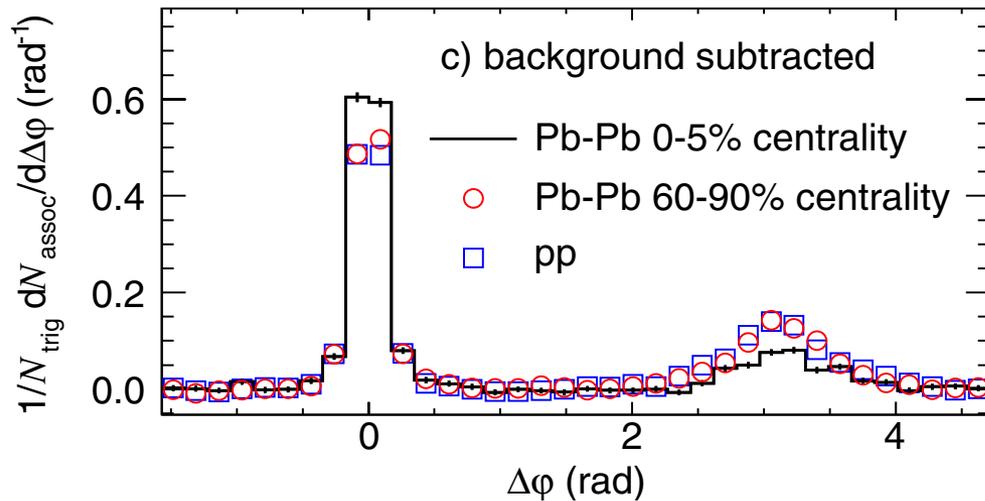
# Jet Quenching



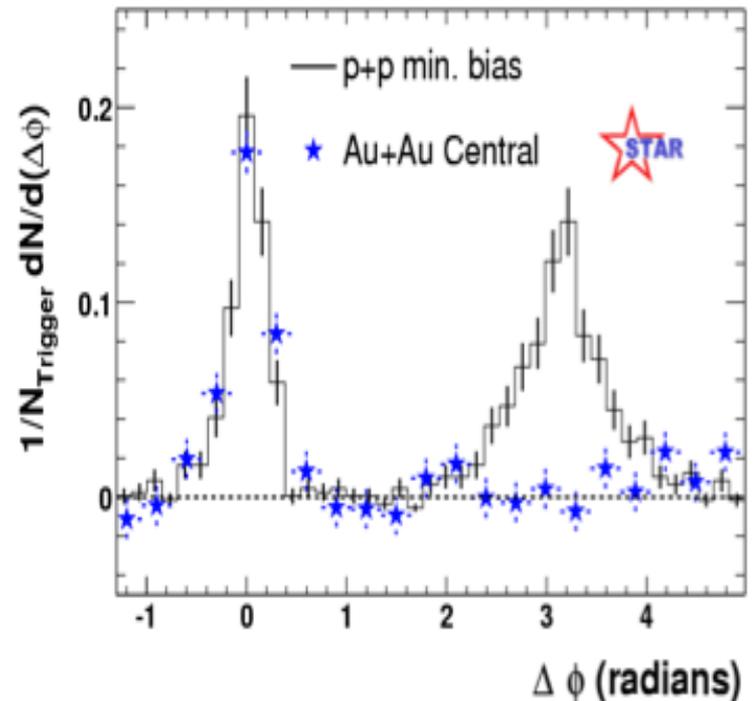
G. Aad et al (ATLAS Collaboration) PRL 105, 252303 (2010)

$$A_j = \frac{E_{T1} - E_{T2}}{E_{T1} + E_{T2}}$$

# Jet Quenching

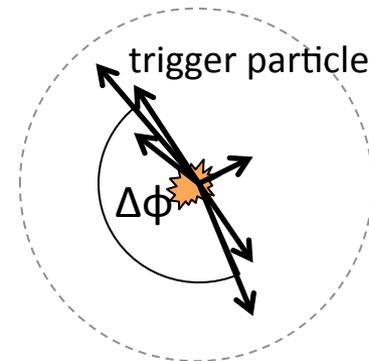


K. Aamodt et al., ALICE Collaboration, PRL 108, 092301 (2012)

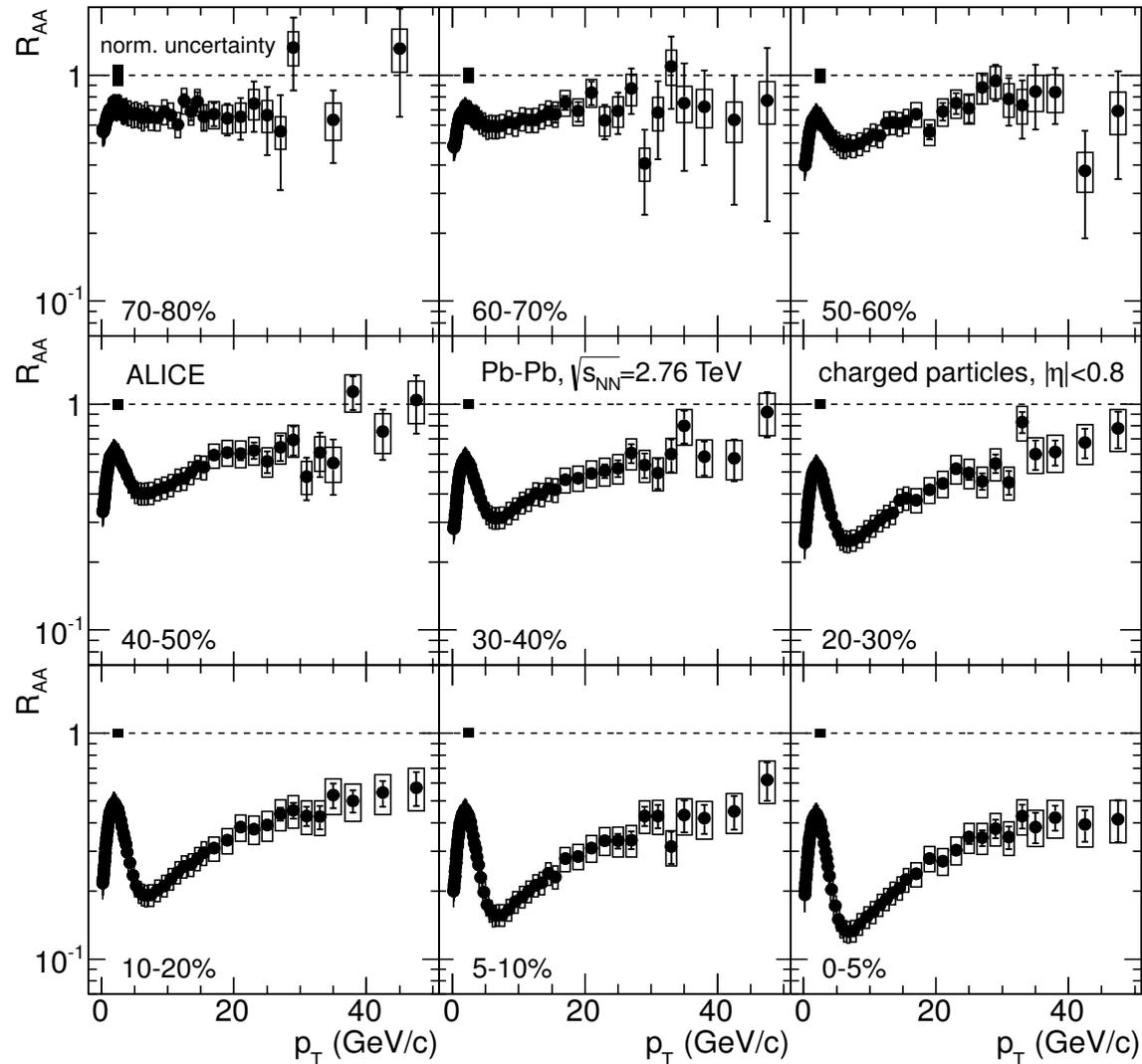


J. Adams, et al., STAR Collaboration, Phys. Rev. Lett. 91 (2003) 072304.

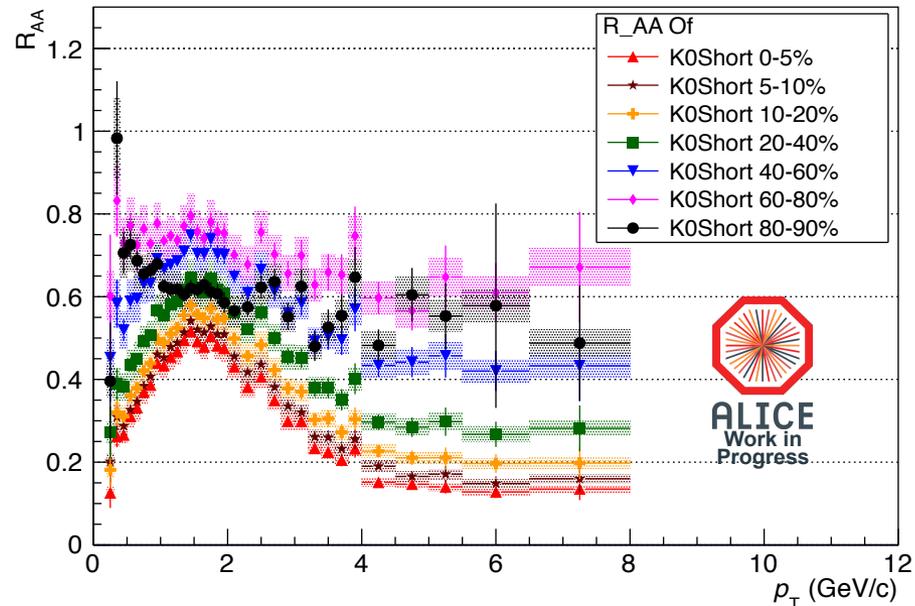
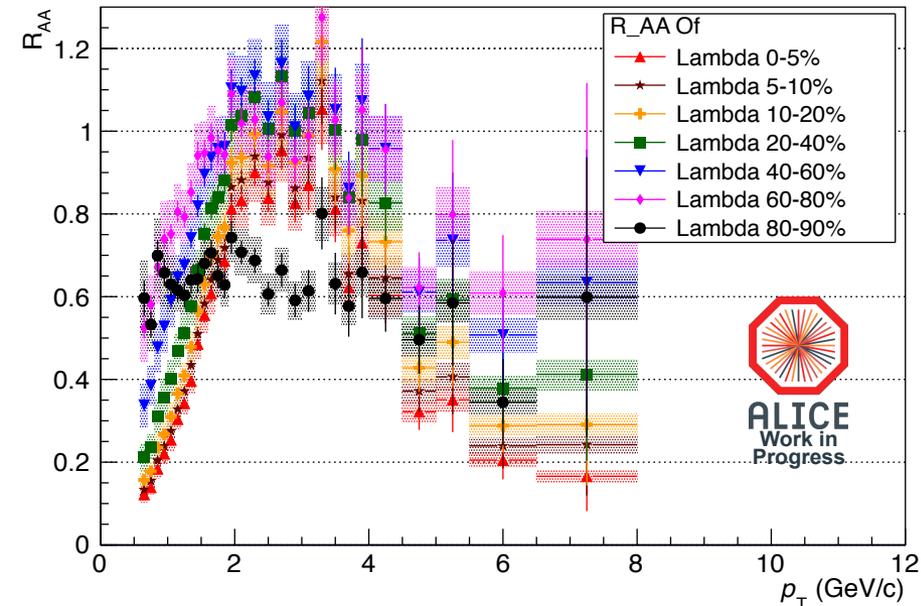
- Plots show distributions of particles relative to a triggered leading particle
- Underlying event subtracted to leave dijet structure



# $R_{AA}$ of Charged Particles



# R\_AA with centrality



- Ratio flattens out as we approach more peripheral events
- Ratio is not consistent with 1 even for most peripheral collisions